#### UNITARITY CONDITIONS FOR STOCHASTIC DIFFERENTIAL EQUATIONS DRIVEN BY NONLINEAR QUANTUM NOISE

Luigi Accardi

Centro Vito Volterra, Università di Roma Tor Vergata 00133 Roma, Italy

> Andreas Boukas American College of Greece Athens 153 42, Greece

# Contents

1	Introduction	3
<b>2</b>	Necessary and sufficient unitarity conditions	6
3	Sufficient unitarity conditions	8
4	Linear independence of the stochastic differentials	10

#### Abstract.

We prove the stochastic independence of the basic integrators of the renormalized square of white noise (SWN). We use this result to deduce the unitarity conditions for stochastic differential equations driven by the SWN.

## 1 Introduction

Linear quantum stochastic calculus on the Boson Fock space, as developed in [7], is associated with the stochastic differentials

$$dB(t) = b(t)dt$$
 ,  $dB^{+}(t) = b^{+}(t)dt$  ,  $dN(t) = b^{+}(t)b(t)dt$ 

corresponding to functionals of the Boson Fock white noise  $b, b^+$  satisfying the commutation relation  $[b(t), b^+(s)] = \gamma \cdot \delta(t-s)$  where  $\gamma > 0$  is the variance of the quantum Brownian motion defined by B and  $B^+$ , and  $\delta$  is the delta function (cf. [2], [5]). A general, representation free, quantum stochastic calculus which included [7] and all other known examples of linear quantum noise was developed in [1] (see also [4], [5], [6]).

The theory has recently been extended (cf. [2], [3]) to include *normally* ordered nonlinear stochastic differentials of the form

$$dB_{(m,n)} = b^+(t)^m b(t)^n dt$$

where  $m, n \in \{0, 1, ...\}$ . This extension required the introduction, in classical probability theory, of renormalization techniques, widely used in quantum field theory.

The white noise  $b^+$ , b is defined as follows: let  $L^2_{\text{sym}}(\mathbf{R}^n)$  denote the space of square integrable functions on  $\mathbf{R}^n$  symmetric under permutation of their arguments, and let

$$F = \bigoplus_{n=0}^{\infty} L_{\text{sym}}^2(\mathbf{R}^n)$$

where: if  $\psi = \{\psi^{(n)}\}_{n=1}^{\infty} \in F$ , then  $\psi^{(0)} \in \mathbf{C}, \ \psi^{(n)} \in L^2_{\text{sym}}(\mathbf{R}^n)$ , and

$$\|\psi\|^2 = |\psi(0)|^2 + \sum_{n=1}^{\infty} \int_{\mathbb{R}^n} |\psi^{(n)}(s_1, \dots, s_n)|^2 ds_1 \dots ds_n$$

Denote by  $S \subset L^2(\mathbf{R}^n)$  the Schwartz space of smooth functions decreasing at infinity faster than any polynomial and let

$$D = \left\{ \psi \in F | \psi^{(n)} \in S, \ \sum_{n=1}^{\infty} n \| \psi^{(n)} \|^2 < \infty \right\}$$

For each  $t \in \mathbf{R}$  define the linear operator  $b(t): D \to F$  by

$$(b(t)\psi)^{(n)}(s_1,\ldots,s_n) = \sqrt{n+1}\psi^{(n+1)}(t,s_1,\ldots,s_n)$$

and the operator valued distribution  $b^+(t)$  by

$$(b^{+}(t)\psi)^{(n)}(s_1,\ldots,s_n) = \frac{1}{\sqrt{n}} \sum_{i=1}^{\infty} \delta(t-s_i)\psi^{(n-1)}(s_1,\ldots,\hat{s}_i,\ldots,s_n)$$

Then

$$B(t) = \int_0^t b(s)ds \ , B^+(t) = \int_0^t b^+(s)ds \ , N(t) = \int_0^t b^+(s)b(s)ds$$

are, for each t, operators acting on D. The renormalized Itô table, proposed in [3], for the stochastic differentials dt, dB,  $dB^+$ ,  $dB_2 = dB_{(0,2)}$ ,  $dB_2^+ = dB_{(2,0)}$ , and dN is

$\begin{array}{c ccccccccccccccccccccccccccccccccccc$		dt	dB	$dB^+$	$dB_2$	$dB_2^+$	dN
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$		0	0	0	0	0	0
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	dB	0	0	$\gamma dt$	0	$2\gamma dB^+$	$\gamma dB$
$dB_2^+$ 0 0 0 0 0 0 0	$dB^+$	0	0	0	0	0	0
$dB_2^+$ 0 0 0 0 0 0 0	$dB_2$	0	0	$2\gamma dB$	0	$4\gamma dN$	$2\gamma dB_2$
		0	0	0	0	0	0
		0	0	$\gamma dB^+$	0	$2\gamma dB_2^+$	$\gamma dN$

Since  $L^2_{\text{Sym}}(\mathbf{R}^n) = L^2_{\text{Sym}}(\mathbf{R}^{\otimes n})$  we can identify F with the symmetric (Boson) Fock space over S. In the case when the elements of S are defined on  $[0, +\infty)$  we denote the Fock space by  $\Gamma(S_+)$ . If  $\psi = \{(n!)^{-1/2} f^{\otimes n}\}$  we denote  $\psi$  by  $\psi(f)$ . With these notations:

$$b_t \psi(f) = f(t)\psi(f), \ b_t^2 \psi(f) = f(t)^2 \psi(f), \ \langle \psi(g), b_t^+ b_t \psi(f) \rangle = \overline{g(t)} f(t) \langle \psi(g), \psi(f) \rangle$$
(2)

Throughout this paper  $b_t^{+2}$  will be interpreted as a quadratic form defined on the linear span of the exponential vectors.

We couple  $\Gamma(S_+)$  with an *initial* Hilbert space  $H_0$  and we define an *adapted process*  $A = \{A(t) : t \geq 0\}$  to be a family of operators on  $H_0 \otimes \Gamma(S_+)$  such that for each  $t, A(t) = A_t \otimes 1$  where  $A_t$  acts on  $H_0 \otimes \Gamma(S_+^{t})$  and 1 is the identity operator on  $\Gamma(S_+^{(t)})$ , where  $S_+^{t]} = \{f \cdot \chi_{[0,t]} / f \in S\}$  and  $S_+^{(t)} = \{f \cdot \chi_{(t,+\infty)} / f \in S\}$ . If, for each  $t, A(t) = A \otimes 1$ , where A is on operator on  $H_0$  and 1 is the identity on  $\Gamma(S_+)$ , then A is a *constant process*. If, for each t, A(t) is a (locally) bounded operator then A is a *(locally) bounded process*.

In what follows we identify B(t),  $B^+(t)$ ,  $B_2(t)$ ,  $B_2^+(t)$ , and N(t) with  $1 \otimes B(t)$ ,  $1 \otimes B^+(t)$ , ...,  $1 \otimes N(t)$  where 1 is the identity on  $H_0$ . For a constant adapted process  $A = \{A(t)/t \ge 0\}$  we denote A(t) simply by A.

Once a quantum stochastic calculus has been constructed, one usually considers the problem of finding conditions under which stochastic differential equations driven by quantum noise admit unitary solutions. It is well known (cf. [7]) that the unique solution  $U = \{U(t) : t \ge 0\}$  of the initial value problem

$$dU(t) = \left[ (iH - \frac{1}{2}L^*L) dt - L^*W dB(t) + LdB^+(t) + (W - 1)dN(t) \right] U(t) \quad (3)$$
$$U(0) = 1 , \quad 0 \le t \le T < +\infty$$

where L, H, W are bounded, constant adapted processes with H self-adjoint and W is unitary i.e.  $U(t)U^*(t) = U^*(t)U(t) = 1$  for each t.

In this note we discuss, in Sections 2 and 3, the unitarity of the solution U of the initial value problem

$$dU(t) = [A_1dt + A_2dB(t) + A_3dB^+(t) + A_4dB_2(t) + A_5dB_2^+(t) + A_6dN(t)]U(t)$$
(4)
$$U(0) = 1 , \qquad 0 < t < T < +\infty$$

where the coefficients  $A_1, A_2, \ldots, A_6$  are bounded, constant adapted processes.

The derivation of the unitarity conditions depends on the linear independence of the stochastic differentials which is established in Section 4.

For an operator K we denote its adjoint by  $K^*$  while its real part is Re  $K = \frac{K+K^*}{2}$ .

## 2 Necessary and sufficient unitarity conditions

In this note we suppose that equation (4) has a solution, defined as a quadratic form on the exponential vectors and we also assume that the expression

 $\langle U(t)\psi(f), U(t)\psi(g)\rangle$ 

has a meaning as a quadratic form on the exponential vectors. Under these assumptions we study under which conditions on the coefficients of equation (1.4), the solution of this equation is unitary, in the sense that the identity between quadratic forms on the exponential vectors

$$\langle U(t)\psi(f), U(t)\psi(g)\rangle = \langle \psi(f), \psi(g)\rangle$$

takes place.

**Proposition 1** The solution U of (4) is unitary if and only if

$$A_1 + A_1^* + A_2 A_2^* \gamma = 0 \tag{5}$$

$$A_2 + A_3^* + A_4 A_2^* 2\gamma + A_2 A_6^* \gamma = 0 \tag{6}$$

$$A_4 + A_5^* + A_4 A_6^* 2\gamma = 0 \tag{7}$$

$$A_6 + A_6^* + A_4 A_4^* 4\gamma + A_6 A_6^* \gamma = 0 \tag{8}$$

and

$$A_1^* + A_1 + A_3^* A_3 \gamma = 0 \tag{9}$$

$$A_3^* + A_2 + A_3^* A_6 \gamma + A_5^* A_3 2 \gamma = 0 \tag{10}$$

$$A_5^* + A_4 + A_5^* A_6 2\gamma = 0 \tag{11}$$

$$A_6^* + A_6 + A_5^* A_5 4\gamma + A_6^* A_6 \gamma = 0$$
(12)

*Proof*: U is unitary if and only if for each  $t \in [0,T]$ ,  $U(t)U^*(t) = U^*(t)U(t) = 1$ . Since  $U(0) = U^*(0) = 1$ ,  $U(t)U^*(t) = 1 \Leftrightarrow d(U(t)U^*(t)) = 0 \Leftrightarrow dU(t) \cdot U^*(t) + U(t) \cdot dU^*(t) + dU(t) \cdot dU^*(t) = 0$  which by Itô's table (1.1) is equivalent to

$$(A_{1}+A_{1}^{*}+A_{2}A_{2}^{*}\gamma)dt + (A_{2}+A_{3}^{*}+A_{4}A_{2}^{*}2\gamma + A_{2}A_{6}^{*}\gamma)dB(t) + (A_{3}+A_{2}^{*}+A_{2}A_{4}^{*}2\gamma + A_{6}A_{2}^{*}\gamma)dB^{+}(t) + (A_{4}+A_{5}^{*}+A_{4}A_{6}^{*}2\gamma)dB_{2}(t) + (A_{5}+A_{4}^{*}+A_{6}A_{4}^{*}2\gamma)dB_{2}^{+}(t) + (A_{6}+A_{6}^{*}+A_{4}A_{6}^{*}4\gamma + A_{6}A_{6}^{*}\gamma)dN(t) = 0$$
(13)

By the linear independence of the stochastic differentials (see Proposition 38), (13) is equivalent to (5) - (8).

Similarly  $U^*(t)U(t) = 1$  is equivalent to (9) – (12).

*Remark*: By (6) and (10),  $A_2$  and  $A_3$  are either both zero or both nonzero. The same is true, by (7) and (11), for  $A_4$  and  $A_5$ .

Corollary 2.1: The solution U of

$$dU(t) = [A_1dt + A_4dB_2(t) + A_5dB_2^+(t) + A_6dN(t)]U(t)$$
(14)  
$$U(0) = 1 , \qquad 0 \le t \le T < +\infty$$

is unitary if and only if

$$A_1 + A_1^* = 0 \tag{15}$$

$$A_4 + A_5^* + A_4 A_6^* 2\gamma = 0 \tag{16}$$

$$A_6 + A_6^* + A_4 A_4^* 4\gamma + A_6 A_6^* \gamma = 0 \tag{17}$$

$$A_5^* + A_4 + A_5^* A_6 2\gamma = 0 \tag{18}$$

$$A_6^* + A_6 + A_5^* A_5 4\gamma + A_6^* A_6 \gamma = 0$$
<sup>(19)</sup>

*Proof*: The proof follows from Proposition 5 by letting  $A_2 = A_3 = 0$ .  $\Box$ 

**Corollary**: The solution U of

$$dU(t) = [A_1dt + A_2dB(t) + A_3dB^+(t) + A_6dN(t)]U(t)$$
(20)  
$$U(0) = 1, \quad 0 \le t \le T < +\infty$$

is unitary if and only if

$$A_1 + A_1^* + A_2 A_2^* \gamma = 0 \tag{21}$$

$$A_2 + A_3^* + A_2 A_6^* \gamma = 0 \tag{22}$$

$$A_6 + A_6^* + A_6 A_6^* \gamma = 0 \tag{23}$$

$$A_1^* + A_1 + A_3^* A_3 \gamma = 0 \tag{24}$$

$$A_3^* A_2 + A_3^* A_6 \gamma = 0 \tag{25}$$

*Proof*: The proof follows from Proposition 5 by letting  $A_4 = A_5 = 0$ .  $\Box$ 

*Remark*: Conditions (21) – (25), for  $\gamma = 1$ , are well known (cf. [7]) and they are certainly satisfied if  $A_1 = \frac{W-1}{\gamma}$ ,  $A_2 = -L^*W$ ,  $A_3 = L$ , and  $A_1 = iH - \frac{\gamma}{2}L^*L$  where L, W, H are bounded operators with W unitary and H self-adjoint. For  $\gamma = 1$  we obtain (3).

## 3 Sufficient unitarity conditions

**Proposition 2** Let L, H, W, M be bounded operators with H self-adjoint and W, M unitary. Suppose also that

$$L^*(1-W) + \sqrt{2}(1 - \operatorname{Re} W)^{1/2}ML = 0$$
(26)

Then  $A_1 = iH - \frac{\gamma}{2}L^*L$ ,  $A_2 = -L^*W$ ,  $A_3 = L$ 

$$A_4 = -\left(\frac{1 - \text{Re }W}{8\gamma^2}\right)^{1/2} MW, \quad A_5 = M^* \left(\frac{1 - \text{Re }W}{8\gamma^2}\right)^{1/2} \text{ and } A_6 = \frac{W - 1}{2\gamma}$$

satisfy (5) - (12). Therefore, the solution U of

$$dU(t) = \left[ \left( iH - \frac{\gamma}{2} L^*L \right) dt - L^*W dB(t) + L dB^+(t) - \left( \frac{1 - \operatorname{Re} W}{8\gamma^2} \right)^{1/2} MW dB_2(t) + M^* \left( \frac{1 - \operatorname{Re} W}{8\gamma^2} \right)^{1/2} dB_2^+ + \frac{W - 1}{2\gamma} dN(t) \right] U(t)$$

$$U(0) = 1 , \qquad 0 \le t \le T < +\infty$$
(27)

is unitary.

Proof:

$$A_{1} + A_{1}^{*} + A_{2}A_{2}^{*}\gamma = iH - \frac{\gamma}{2}L^{*}L - iH - \frac{\gamma}{2}L^{*}L + (-L^{*}W)(-W^{*}L)\gamma = 0$$

$$A_{2} + A_{3}^{*} + 2\gamma A_{4}A_{2}^{*} + \gamma A_{2}A_{6}^{*} = -L^{*}W + L^{*} + 2\gamma \left(-\left(\frac{1 - \operatorname{Re}W}{8\gamma^{2}}\right)^{1/2}\right)MW(-W^{*}L) + \gamma (-L^{*}W)\frac{W^{*} - 1}{2\gamma} = -L^{*}W + L^{*} + \frac{1}{\sqrt{2}}(1 - \operatorname{Re}W)^{1/2}ML + \frac{L^{*}W - L^{*}}{2} = \frac{1}{2}[L^{*}(1 - W) - \sqrt{2}(1 - \operatorname{Re}W)^{1/2}ML] = 0 \text{ by } (28).$$
(28)

Finally,

$$A_{6} + A_{6}^{*} + A_{4}A_{4}^{*}4\gamma + A_{6}A_{6}^{*}\gamma = \frac{W-1}{2\gamma} + \frac{W^{*}-1}{2\gamma} + \left(\frac{1-\operatorname{Re}W}{8\gamma^{2}}\right)^{1/2} MWW^{*}M\left(\frac{1-\operatorname{Re}W}{8\gamma^{2}}\right)^{1/2} 4\gamma$$

$$+\frac{W-1}{2\gamma}\frac{W^{*}-1}{2\gamma}\gamma = \frac{\text{Re }W-1}{\gamma} + \frac{1-\text{ Re }W}{2\gamma} + \frac{1-\text{ Re }W}{2\gamma} = 0$$

thus proving (5) - (8).

The proof of (9) - (12) is similar.  $\Box$ 

*Remark*: Equation (28) connects the linear case (20) with the non–linear case (14). Several examples of L, M, W satisfying (28) are given in the following corollary to Proposition 28.

**Corollary 3.1**: Let L, M, W, H be bounded operators with H selfadjoint and M, W unitary. The solution  $U = \{U(t) : 0 \le t \le T < +\infty\}$ of each of the following initial value problems, with initial value U(0) = 1, is unitary.

$$dU(t) = [iHdt]U(t) \tag{29}$$

$$dU(t) = \left[ \left( iH - \frac{\gamma}{2} L^*L \right) dt - L^* dB(t) + L dB^+(t) \right] U(t)$$
(30)

$$dU(t) = \left[iHdt + \frac{W-1}{2\gamma}dN(t)\right]U(t)$$
(31)

$$dU(t) = \left[iHdt - \left(\frac{1-\text{Re }W}{8\gamma^2}\right)^{1/2} MW dB_2(t) + M^* \left(\frac{1-\text{Re }W}{8\gamma^2}\right)^{1/2} dB_2^+(tB2) + \frac{W-1}{2\gamma} dN(t)\right] U(t)$$
(33)

$$dU(t) = \left[ \left( iH - \frac{\gamma}{2} |\gamma|^2 \right) dt - \overline{\gamma} e^{i\theta} dB(t) + \lambda dB^+(t) \right]$$
(34)

$$-\frac{(e^{i\theta}-1)e^{i\theta}}{4\gamma}\frac{\overline{\lambda}}{\lambda}dB_2(t) + \frac{(e^{i\theta}-1)}{4\gamma}\frac{\overline{\lambda}}{\overline{\lambda}}dB_2^+(t) + \frac{e^{i\theta}-1}{2\gamma}dN(t)\Big]U(t) \quad (35)$$

where  $\theta \in \mathbf{R} - \{2k\pi/k = 0, \pm 1, \pm 2, ...\}, \lambda \in \mathbf{C} - \{0\}.$ 

$$dU(t) = \left[ \left( iH - \frac{\gamma}{2} \right) dt - e^{i\theta} L^* dB(t) + L dB^+(t) - \frac{(e^{i\theta} - 1)}{4\gamma} (L^2)^* dB_2(t) (36) + \frac{e^{-i\theta} - 1}{4\gamma} L^2 dB_2^+(t) + \frac{e^{i\theta} - 1}{2\gamma} dN(t) \right] U(t)$$
(37)

where  $\theta \in \mathbf{R} - \{2k\pi/k = 0, \pm 1, \pm 2, ...\}$  and  $LL^* = L^*L = 1$ .

*Proof*: The proof follows from Proposition 28 by taking

(a) 
$$L = 0, \quad W = 1$$
  
(b)  $W = 1$   
(c)  $L = M = 0$   
(d)  $L = 0$   
(e)  $L = \lambda 1, \quad W = e^{i\theta}, \quad M = \frac{e^{i\theta} - 1}{\sqrt{2}(1 - \cos\theta)^{1/2}\lambda} \frac{\overline{\lambda}}{\lambda} 1$   
(f)  $W = e^{i\theta} 1, \quad M = \frac{e^{i\theta} - 1}{\sqrt{2}(1 - \cos\theta)^{1/2}} (L^2)^* \square$ 

## 4 Linear independence of the stochastic differentials

**Proposition 3** For each i = 1, 2, ..., 6 let  $A_i = \hat{A}_i \otimes 1$  be a constant adapted process acting on  $H_0 \otimes \Gamma(S_+)$  and suppose that for all  $t \ge 0$ 

$$A_1dt + A_2dB(t) + A_3dB^+(t) + A_4dB_2(t) + A_5dB_2^+(t) + A_6dN(t) = 0 \quad (38)$$

Then

$$A_1 = A_2 = \dots = A_6 = 0 \tag{39}$$

on the exponential domain.

Proof: By (38), for all 
$$f, g \in S_+, u, v \in H_0$$
, and  $t \ge 0$   
 $\langle u \otimes \psi(f), [A_1 dt + \ldots + A_6 dN(t)]v \otimes \psi(g) \rangle = 0$  (40)

By (2), (5) implies

$$\langle u \otimes \psi(f), A_1 v \otimes \psi(g) \rangle + g(t) \langle u \otimes \psi(f), A_2 v \otimes \psi(g) \rangle +$$
  
$$\overline{f(t)} \langle u \otimes \psi(f), A_3 v \otimes \psi(g) \rangle + g(t)^2 \langle u \otimes \psi(f), A_4 v \otimes \psi(g) \rangle +$$
  
$$\overline{f(t)}^2 \langle u \otimes \psi(f), A_5 v \otimes \psi(g) \rangle + \overline{f}(t) g(t) \langle u \otimes \psi(f), A_6 v \otimes \psi(f) \rangle = 0 \quad (41)$$

Since, for each i = 1, 2, ..., 6,  $\langle u \otimes \psi(f), A_i v \otimes \psi(g) \rangle = \langle u, \hat{A}_i v \rangle \langle \psi(f), \psi(g) \rangle$ and  $\langle \psi(f), \psi(g) \rangle = \exp(\langle f, g \rangle) \neq 0$ , (41) implies

$$\langle u, \hat{A}_1 v \rangle + g(t) \langle u, \hat{A}_2 v \rangle + \overline{f(t)} \langle u, \hat{A}_3 v \rangle + g(t)^2 \langle u, \hat{A}_4 v \rangle + \overline{f(t)}^2 \langle u, \hat{A}_5 v \rangle + \overline{f(t)} g(t) \langle u, A_6 v \rangle = 0$$

$$(42)$$

Taking f and g to be such that g(t) = 0 and  $f(t) \neq 0$  we obtain

$$\langle u, \hat{A}_3 v \rangle + \overline{f(t)} \langle u, \hat{A}_5 v \rangle = 0$$
(43)

for all  $u, v \in H_0$ . Since f(t) can be fixed, while u, v are arbitrary, it follows that  $\langle u, \hat{A}_3 v \rangle = \langle u, \hat{A}_5 v \rangle = 0$  and so  $\hat{A}_3 = \hat{A}_5 = 0$ .

Similarly, if g is such that g(t) = 0 then

$$\langle u, \hat{A}_2 v \rangle + g(t) \langle u, \hat{A}_4 v \rangle = 0 \tag{44}$$

from which we obtain  $\hat{A}_2 = \hat{A}_4 = 0$ .

Thus  $\hat{A}_i = 0$  and so  $A_i = 0$  for all  $i = 1, 2, \ldots, 6$ .  $\Box$ 

**Proposition 4** For each i = 1, 2, ..., 6 and  $t \ge 0$ , let  $A_i = \{A_i(t) = \hat{A}_i(t) \otimes 1 : t \ge 0\}$  be an adapted process acting on  $H_0 \otimes \Gamma(S_+)$  and assume that: (a) The map  $t \in [0, +\infty) \to \langle u \otimes \psi(f), A_i(t)v \otimes \psi(g) \rangle$  is continuous for all  $u, v \in H_0$  and  $f, g \in S_+$ . (b) The map  $(v, g) \in H_0 \times S_+ \to A_i(t)v \otimes \psi(g)$  is continuous for all  $t \ge 0$ . (c) For all  $t \ge 0, u, v \in H_0$ , and  $f, g \in S_+$  with Re  $f \cdot$  Re g = 0, Im  $\langle u \otimes \psi(f), A_i(t)v \otimes \psi(g) \rangle = 0$  for all i = 2, 3, ..., 6. Then

$$A_{1}(t)dt + A_{2}(t)dB(t) + A_{3}(t)dB^{+}(t) + A_{4}dB_{2}(t) + A_{5}(t)dB_{2}^{+}(t) + A_{6}(t)dN(t) = 0$$
(45)

for all  $t \ge 0$ , implies

$$A_1 \equiv A_2 \equiv \ldots \equiv A_6 \equiv 0$$

on the exponential domain.

*Proof*: As in the proof of Proposition 38, in view of (45), for all  $f, g \in S_+$ ,  $u, v \in H_0$ , and  $t \ge 0$ 

$$\langle u \otimes \psi(f), A_1(t)v \otimes \psi(g) \rangle + g(t) \langle u \otimes \psi(f), A_2(t)v \otimes \psi(g) \rangle$$

$$+\overline{f(t)}\langle u \otimes \psi(f), A_3(t)v \otimes \psi(g) \rangle + g(t)^2 \langle u \otimes \psi(f), A_4(t)v \otimes \psi(g) \rangle +\overline{f(t)}^2 \langle u \otimes \psi(f), A_5(t)v \otimes \psi(g) \rangle + \overline{f(t)}g(t) \langle u \otimes \psi(f), A_6(t)v \otimes \psi(g) \rangle = 0$$
(46)

Allowing f, g to vary over continuous functions with compact support contained in [0,t), (46) implies  $\langle u \otimes \psi(f), A_1(t)v \otimes \psi(g) \rangle = 0$ . By the adaptedness of  $A_1$  and the totality of  $\{u \otimes \psi(f)/f \text{ continuous, sup } pf \subset [0,t]\}$ in  $H_0 \otimes \Gamma(S_+)$  it follows that  $\hat{A}_i(t) = 0$  and so  $A_1 \equiv 0$ .

Allowing g to vary over continuous functions with compact support contained in [0, t), and taking f such that Re f = 0 and  $f(t) \neq 0$  we obtain from (46), with  $A_1(t) = 0$ ,

$$i\langle u \otimes \psi(f), A_3(t)v \otimes \psi(g) \rangle + \operatorname{Im} f(t)\langle u \otimes \psi(f), A_5(6)v \otimes \psi(g) \rangle = 0 \quad (47)$$

By assumption (c) and the fact that Im  $f(t) = f(t) \neq 0$ , (47) yields

$$\langle u \otimes \psi(f), A_3(t)v \otimes \psi(g) \rangle = \langle u \otimes \psi(f), A_5(t)v \otimes \psi(g) \rangle = 0$$

and so, as before,  $A_3 \equiv A_5 \equiv 0$ .

Repeating the above argument interchanging the roles of f and g we obtain  $A_2 \equiv A_4 \equiv 0$ .

Finally, if we vary f, g over continuous functions such that  $\overline{f(t)}g(t) \neq 0$ , we conclude that  $A_6 \equiv 0$ .  $\Box$ 

**Proposition 5** Let  $A_i = \{A_i(t)/t \ge 0\}$  be, for each i = 1, 2, ..., 6, an adapted process in  $H_0 \otimes \Gamma(S_+)$  such that

(a) For fixed  $f, g \in S_+$  and  $u, v \in H_0$  the map  $t \in [0, +\infty) \to \langle u \otimes \psi(f), A_i(t)v \otimes \psi(g) \rangle$  is continuous

(b) For fixed  $t \in [0, +\infty)$  the map  $(v, g) \in H_0 \times S_+ \to A_i(t)v \otimes \psi(g)$  is continuous.

Then the stochastic differentials dt, dB,  $dB^+$ ,  $dB_2$ ,  $dB_2^+$ , and dN are linearly independent i.e.

 $A_1(t)dt + \ldots + A_6(t)dN(t) = 0$  for all  $t \ge 0$  implies  $A_1 \equiv \ldots \equiv A_6 \equiv 0$ 

on the exponential domain.

*Proof*: For all  $u, v \in H_0$  and  $f, g \in S_+$ 

$$\langle u \otimes \psi(f), [A, (t)dt + A_2(t)dB(t) + A_3(t)dB^+(t) + A_4(t)dB_2(t) + A_5(t)dB_2^+(t) \rangle$$

$$+A_6(t)dN(t)]v\otimes\psi(g)\rangle=0$$

implies

$$\langle u \otimes \psi(f), [A_1(t) + g(t)A_2(t) + \overline{f(t)}A_3(t) + g(t)^2 + \overline{f(t)}^2 A_5(t) + f(t)\overline{g(t)}A_6(t)]v \otimes \psi(g) \rangle = 0$$

$$(48)$$

Allowing f, g to vary over continuous functions with compact support contained in [0, t) we obtain from (4.11)  $\langle u \otimes \psi(f), A_1(t)v \otimes \psi(g) \rangle = 0$ . By the adaptedness of  $A_1$  and the totality of  $\{u \otimes \psi(f)/u \in H_0, f \text{ continuous}, \sup pg \subset [0, t)\}$  in  $H_0 \otimes \Gamma(S_+)$  it follows that  $A_1 \equiv 0$ .

Letting f, g vary over continuous functions with  $f(t) \neq 0$  and  $\sup pg \subset [0, t)$  (48) implies

$$\langle u \otimes \psi(f), [\overline{f(t)}A_3(t) + \overline{f(t)}^2 A_5(t)]v \otimes \psi(g) \rangle = 0$$

and so

$$\langle u \otimes \psi(f), [A_3(t) + \overline{f(t)}A_5(t)]v \otimes \psi(g) \rangle = 0$$

Thus

$$\langle u \otimes \psi(f), [A_3(t)dt + \overline{f(t)}A_5(t)dt]v \otimes \psi(g) \rangle = 0$$

i.e.

$$\langle u \otimes \psi(f), [A_3(t)dt + A_5(t)dB^+(t)]v \otimes \psi(g) \rangle = 0$$
(49)

By a totality argument and by the arbitrariness of t, (49) implies

$$A_3(t)dt + A_5(t)dB^+(t) = 0 (50)$$

But dt, dB,  $dB^+$  and dN are known to be linearly independent (cf. [8]). Thus (50) implies  $A_3 \equiv A_5 \equiv 0$ .

Similarly  $A_2 \equiv A_4 \equiv 0$ .

Now letting f, g be such that  $f(t)\overline{g(t)} \neq 0$  we conclude that  $A_6 \equiv 0$ .  $\Box Remark$ : A formal proof of the linear independence of dt, dB,  $dB^+$ ,  $dB_2$ ,

 $dB_2^+,\,\mathrm{and}\;dN$  can be obtained with the use of Itô's table as follows: suppose that

$$A_1dt + A_2dt + A_3dB^+ + A_4dB_2 + A_5dB_2^+ + A_6dN = 0$$
(51)

Multiplying (51) from the left by dB we obtain by (1)

$$A_3dt + 2\gamma A_5dB^+ + \gamma A_6dN = 0 \tag{52}$$

By the linear independence of dt, dB,  $dB^+$ , and dN (52) yields

$$A_3 \equiv A_5 \equiv A_6 \equiv 0$$

Multiplying (4.14) from the right by  $dB^+$  we obtain

$$A_2dt + 2\gamma A_4dB = 0 \tag{53}$$

As before, (53) yields  $A_2 \equiv A_4 \equiv 0$ . Finally  $A_1 dt = 0$  implies  $A_1 \equiv 0$ .

#### References

- Accardi L., Fagnola F., Quagebeur J., Representation free quantum stochastic calculus, J. Func. Anal., 104 (1992), 140–197.
- [2] Accardi L., Lu Y.G., Obata N., Towards a non-linear extension of stochastic calculus, Publications of the Research Institute for Mathematical Sciences, Kyoto, RIMS Kokyuroki 957, Obata N. (ed.), (1996), 1–15.
- [3] Accardi L., Lu Y.G., Volovich I., *White noise approach to stochastic calculus and nonlinear Ito tables*, submitted to: Nagoya Journal of Mathematics.
- [4] Accardi L., Quantum Stochastic Calculus, In "Proceedings IV Vilnius Conference on Probability and Mathematical Statistics", VNU Science, 1986.
- [5] Accardi L., A mathematical theory of quantum noise, in "Proceedings 1<sup>st</sup> World Congress of the Bernoulli Society" (Y. Prohorov and V. Sazonov, Eds.), 1 (1987) pp. 427–444, VNU Science.
- [6] Accardi L., Parthasarathy K.R., Stochastic calculus on local algebras, in "Quantum Probability and Application II (Proceedings Heidelberg, 1984)", Lecture Notes in Math., **1136**, (1985) pp. 9–23, Springer–Verlag, New York.
- [7] Hudson R.L., Parthasarathy K.R., Quantum Ito's formula and stochastic evolutions, Comm. Math. Phys. 93 (1984), 301–323.
- [8] Parthasarathy K.R., Quantum Stochastic Calculus, in "Proceedings of Stochastic Processes and their applications (Nagoya 1985)" (K. Itô, T. Hida, Eds.), Lecture Notes in Math., **1203**, (1986) pp. 177–196, Springer–Verlag, New York.