# DECOMPOSITIONS OF THE FREE PRODUCT OF GRAPHS 

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#### Abstract

We study the free product of rooted graphs and its various decompositions using quantum probabilistic methods. We show that the free product of rooted graphs is canonically associated with free independence, which completes the proof of the conjecture that there exists a product of rooted graphs canonically associated with each notion of noncommutative independence which arises in the axiomatic theory. Using the orthogonal product of rooted graphs, we decompose the branches of the free product of rooted graphs as "alternating orthogonal products". This leads to alternating decompositions of the free product itself, with the star product or the comb product followed by orthogonal products. These decompositions correspond to the recently studied decompositions of the free additive convolution of probability measures in terms of boolean and orthogonal convolutions, or monotone and orthogonal convolutions. We also introduce a new type of quantum decomposition of the free product of graphs, where the distance partition of the set of vertices is taken with respect to a set of vertices instead of a single vertex. We show that even in the case of widely studied graphs this yields new and more complete information on their spectral properties, like spectral measures of a (usually infinite) set of cyclic vectors under the action of the adjacency matrix.


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## 1. Introduction

Graph theory plays an important role in many branches of mathematics and its applications. In particular, in solid state physics the idea that the study of a free dynamics on a nonhomogeneous graph might be equivalent to the study of an
interacting dynamics on a homogeneous graph (a discrete version of the basic idea of general relativity) found interesting applications in the study of Bose-Einstein condensation. ${ }^{7,8,23}$

In many of these applications, the main focus was not so much on the combinatorial properties of single graphs, as on the analytical aspects of the asymptotics of large graphs (when the number of vertices tends to infinity). In that connection, typical objects of interest are spectral distributions of the adjacency matrix with respect to specially chosen states, or the full spectrum.

From these investigations an interesting class of graphs has emerged, namely those which are built from subgraphs expressible as products of simpler graphs. Intuitively, a product of graphs is a rule to construct a new graph by glueing together two given graphs subject to additional conditions (like associativity). It is known that there exist several different products among graphs, including those studied extensively in discrete mathematics like the lexicographic product, the Cartesian product or the strong product.

On the other hand, the experience of quantum probability teaches us that, at an algebraic level, certain types of products of quantum probability spaces correspond to different notions of stochastic independence. ${ }^{28,31}$ Recall that in the concept of a quantum probability space one has to distinguish a state, which, in the simplest graph setting, corresponds to a distinguished vertex called root. It is therefore natural to conjecture that certain types of products of rooted graphs could be canonically associated with the main notions of stochastic independence Moreover, one might expect that such graph products are important types of products, from which one could not only construct more complicated graphs, but also obtain information about their spectra, using (well-established or entirely new) quantum probabilistic techniques.

The well-known case of the Cayley graph of a free product of groups and its relation to free independence of Voiculescu ${ }^{33}$ and to the free product of states ${ }^{3,33}$ can be viewed as the first evidence that such a conjecture is true. That such a relation holds also in the general case of the free product of rooted graphs can be shown using the free probability techniques. Thus, we explicitly state and prove the fact that the free product of graphs, introduced by $\mathrm{Znojko}^{39}$ for symmetric graphs and generalized by Quenell ${ }^{30}$ and Gutkin ${ }^{11}$ to rooted graphs, is canonically associated with the notion of free independence. More generally, we show that the hierarchy of $m$-freeness introduced by one of the authors, ${ }^{18}$ corresponds to a natural hierarchy of $m$-free products of graphs, as well as to the natural inductive definition of the free product. ${ }^{39,11}$

The next important evidence supporting this conjecture was the discovery, by Accardi, Ben Ghorbal and Obata, ${ }^{1}$ that the comb product of rooted graphs is canonically related to the monotone independence. ${ }^{22,26}$ A similar connection between the star product of rooted graphs and boolean independence ${ }^{32}$ was established by one of the authors ${ }^{17}$ and Obata. ${ }^{29}$ Since it is well known that the Cartesian product of graphs is naturally related to tensor (or boson) independence, the correspondence
between the main notions of stochastic independence (which arise in the axiomatic theory) and certain types of products of rooted graphs is completed.

In all the above-mentioned cases, the canonical relation between a notion of product of two graphs, $\mathcal{G}_{1}=\left(V_{1}, E_{1}\right)$ and $\mathcal{G}_{2}=\left(V_{2}, E_{2}\right)$, and a notion of independence $\mathcal{I}$ is realized by showing that the adjacency matrix of their product is naturally split into a sum of operator random variables which are $\mathcal{I}$-independent with respect to a given state. In all cases this decomposition can be obtained by embedding the algebra of operators on the $l^{2}$-space of the product graph into an appropriate tensor product, paralleling the construction of Refs. 18 and 19.

Let us remark in this context that the Cartesian, comb and star products appear to be "basic" graph products since the corresponding vertex sets are subsets of $V_{1} \times V_{2}$, whereas the vertex set of the free product of graphs is the free product $V_{1} * V_{2}$ of rooted sets and the construction of $\mathcal{G}_{1} * \mathcal{G}_{2}$ involves infinitely many copies of $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$. This is the main reason why the free product of graphs has this peculiar feature that it admits a variety of natural decompositions. In particular, the decomposition into the sum of "freely independent subgraphs", although the most natural from the point of view of free independence, is not always the most intuitive or the most convenient.

In particular, we find new decompositions related to the "growth" of $\mathcal{G}_{1} * \mathcal{G}_{2}$ exhibited by its inductive definitions. Motivated by the recent work of one of the authors on the decompositions of the free additive convolution of probability measures ${ }^{20}$ (see also Ref. 21), we study a new type of "basic" product of rooted graphs called the orthogonal product of graphs, related to the orthogonal convolution of probability measures introduced in Ref. 20. We show that this is the orthogonal product which is the main building block of $\mathcal{G}_{1} * \mathcal{G}_{2}$ since it is responsible for its "growth". In fact, it allows us to decompose its "branches" ${ }^{30}$ into products of alternating $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$. Since one obtains the free product by taking the comb product or the star product of "branches", as shown in Ref. 30 (we use quantum probabilistic techniques to simplify the proofs presented there), we arrive at two alternating decompositions of the free product of graphs - the comb-orthogonal decomposition and the star-orthogonal decomposition. More importantly, using the orthogonal convolution, one can study spectral distributions of free products of uniformly locally finite graphs in a very intuitive manner and see their direct relation to continued fractions, especially periodic and mixed periodic Jacobi continued fractions ${ }^{14}$ (without using the Rtransforms).

Finally, in order to get a more detailed information about the structure of the spectrum of $\mathcal{G}_{1} * \mathcal{G}_{2}$, we introduce a new type of quantum decomposition of the adjacency matrix of a given graph $\mathcal{G}$. This decomposition is based on a new type of distance partition $V=\bigcup_{n=0}^{\infty} \mathcal{V}_{n}$ of the set of vertices, where $\mathcal{V}_{n}$ is the set of vertices of $\mathcal{G}$, whose distance from a set of vertices (instead of a single vertex) is equal to $n$. This leads to a different quantum decomposition of the adjacency matrix $A(\mathcal{G})$ into the sum of a creation, annihilation and diagonal operators. It allows us to
derive a cyclic direct sum decomposition of the Hilbert space $l^{2}(V)$ together with the spectral distributions associated with different cyclic (vacuum) vectors.

For classical (random walk) methods applied to the free products of Cayley graphs and other infinite graphs, we refer the reader to Refs. 2, 6, 9, 10, 15, 16, 37, 38 and references therein.

## 2. Notation

By a rooted set we understand a pair $(X, e)$, where $X$ is a countable set and $e$ is a distinguished element of $X$ called root. By a rooted graph we understand a pair $(\mathcal{G}, e)$, where $\mathcal{G}=(V, E)$ is a non-oriented graph with the set of vertices $V=V(\mathcal{G})$, and the set of edges $E=E(\mathcal{G}) \subseteq\left\{\left\{x, x^{\prime}\right\}: x, x^{\prime} \in V, x \neq x^{\prime}\right\}$ and $e \in V$ is a distinguished vertex called the root. We will also denote by $\mathcal{G}$ the rooted graph $(\mathcal{G}, e)$ if no confusion arises, especially if the graph is symmetric, i.e. for any $x \neq x^{\prime}$ there exists an automorphism $\tau$ of $\mathcal{G}$ for which $\tau(x)=x^{\prime}$ (in other words, all vertices are equivalent).

For rooted graphs we will use the notation

$$
\begin{equation*}
V^{0}=V \backslash\{e\} \tag{2.1}
\end{equation*}
$$

Two vertices $x, x^{\prime} \in V$ are called adjacent if $\left\{x, x^{\prime}\right\} \in E$, i.e. vertices $x, x^{\prime}$ are connected with an edge. Then we write $x \sim x^{\prime}$. Simple graphs have no loops, i.e. $\{x, x\} \notin E$ for all $x \in V$. The degree of $x \in V$ is defined by $\kappa(x)=\mid\left\{x^{\prime} \in V\right.$ : $\left.x^{\prime} \sim x\right\} \mid$, where $|I|$ stands for the cardinality of $I$. A graph is called locally finite if $\kappa(x)<\infty$ for every $x \in V$. It is called uniformly locally finite if $\sup \{\kappa(x): x \in$ $V\}<\infty$.

For $x \in V$, let $\delta(x)$ be the indicator function of the one-element set $\{x\}$. Then $\{\delta(x), x \in V\}$ is an orthonormal basis of the Hilbert space $l^{2}(V)$ of square integrable functions on the set $V$, with the usual inner product.

The adjacency matrix $A=A(\mathcal{G})$ of $\mathcal{G}$ is a $0-1$ matrix defined by

$$
A_{x, x^{\prime}}= \begin{cases}1 & \text { if } x \sim x^{\prime}  \tag{2.2}\\ 0 & \text { otherwise }\end{cases}
$$

We identify it with the densely defined symmetric operator on $l^{2}(V)$, denoted by the same symbol, with the action on the basis of $l^{2}(V)$ given by

$$
\begin{equation*}
A \delta(x)=\sum_{x \sim x^{\prime}} \delta\left(x^{\prime}\right) \tag{2.3}
\end{equation*}
$$

for $x \in V$. Notice that the sum on the right-hand side is finite since our graph is assumed to be locally finite. It is known that $A(\mathcal{G})$ is bounded iff $\mathcal{G}$ is uniformly locally finite. The closure of $A(\mathcal{G})$ is called the adjacency operator of $\mathcal{G}$ and its spectrum - the spectrum of $\mathcal{G}$.

The unital algebra generated by $A$, i.e. the algebra of polynomials in $A$, is called the adjacency algebra of $\mathcal{G}$ and is denoted by $\mathcal{A}(\mathcal{G})$, or simply $\mathcal{A}$.

For simplicity of presentation, by a graph we shall understand a non-oriented, connected and locally finite simple graph with a non-empty set of edges. Any rooted graph of type $(\mathcal{G}, e)$, where $\mathcal{G}$ is a graph in this sense, will also be called a graph if no confusion arises. However, it is not hard to observe that the theorems remain true if we allow graphs to have loops, or even if we take multigraphs, i.e. graphs with multiple edges.

## 3. Convolutions, Transforms and Graph Products

By the spectral distribution of $A=A(\mathcal{G})$ in a state $\psi$ on $l^{2}(V)$ we understand the sequence $\left(\psi\left(A^{n}\right)\right)_{n \geq 0}$, or the associated measure $\mu$ (in case the moment problem has the unique solution), for which

$$
\begin{equation*}
\psi\left(A^{n}\right)=\int_{\mathbb{R}} x^{n} \mu(d x), \quad n \in \mathbb{N} \cup\{0\} \tag{3.1}
\end{equation*}
$$

and by the spectral distribution of the rooted $\operatorname{graph}(\mathcal{G}, e)$ we understand the spectral distribution of $A(\mathcal{G})$ in the state $\varphi_{e}(\cdot)=\langle\cdot \delta(e), \delta(e)\rangle$. The spectral distribution of $(\mathcal{G}, e)$ is important in the evaluation of the $\operatorname{spectrum} \operatorname{spec}(\mathcal{G})$ of the graph $\mathcal{G}$. In some cases (homogenous trees and $n$-ary trees are the easiest examples) it is even so that $\operatorname{spec}(\mathcal{G})$ agrees with the support of the spectral distribution of $(\mathcal{G}, e)$.

For any probability distribution $\mu$ with the sequence of moments $\left(M_{n}\right)_{n \geq 0}$, we define its moment generating function as the formal power series $M_{\mu}(z)=$ $\sum_{n=0}^{\infty} M_{n} z^{n}$. The corresponding Cauchy transform, K-transform, reciprocal Cauchy transform and R-transform are defined, respectively, by the formal power series

$$
\begin{align*}
& G_{\mu}(z)=\frac{1}{z} M_{\mu}\left(\frac{1}{z}\right)  \tag{3.2}\\
& K_{\mu}(z)=z-\frac{1}{G_{\mu}(z)}  \tag{3.3}\\
& F_{\mu}(z)=\frac{1}{G_{\mu}(z)}  \tag{3.4}\\
& R_{\mu}(z)=-\frac{1}{z}+G_{\mu}^{-1}(z) . \tag{3.5}
\end{align*}
$$

Let $\left(\mathcal{G}_{1}, e_{1}\right)$ and $\left(\mathcal{G}_{2}, e_{2}\right)$ be two graphs with adjacency matrices $A_{1}=A\left(\mathcal{G}_{1}\right)$, $A_{2}=A\left(\mathcal{G}_{2}\right)$ and spectral distributions $\mu, \nu$, respectively. By $\mu \uplus \nu$ we denote the boolean convolution of $\mu$ and $\nu$ associated with boolean independence. ${ }^{32}$ By $\mu \triangleright \nu$ we denote the monotone convolution associated with monotone independence. ${ }^{26}$ By $\mu \boxplus \nu$ we denote the free additive convolution associated with free independence. ${ }^{33}$ Finally, by $\mu \vdash \nu$ we denote the orthogonal convolution associated with orthogonal independence, recently introduced in Ref. 20. The following identities hold:

$$
\begin{equation*}
K_{\mu \uplus \nu}(z)=K_{\mu}(z)+K_{\nu}(z), \tag{3.6}
\end{equation*}
$$

$$
\begin{align*}
& R_{\mu \boxplus \nu}(z)=R_{\mu}(z)+R_{\nu}(z)  \tag{3.7}\\
& F_{\mu \triangleright \nu}(z)=F_{\mu}\left(F_{\nu}(z)\right)  \tag{3.8}\\
& K_{\mu \vdash \nu}(z)=K_{\mu}\left(F_{\nu}(z)\right) \tag{3.9}
\end{align*}
$$

The above relations can be treated as definitions of the associated convolutions, although these are usually introduced by using some notion of noncommutative "independence" which parallels the connection between the usual (classical) convolution of distributions (measures) and the notion of classical independence. Note that the K - and R-transforms are additive under the considered convolutions (see Refs. 32, 33 and 37 ), thus they play the role of the logarithm of the Fourier transform. In the case of the monotone and orthogonal convolutions, addition of transforms is replaced by composition (see Refs. 26 and 20).

Proposition 3.1. The following relations hold:

$$
\begin{align*}
& F_{\mu \uplus \nu}(z)=F_{\mu}(z)+F_{\nu}(z)-z,  \tag{3.10}\\
& F_{\mu \vdash \nu}(z)=F_{\mu}\left(F_{\nu}(z)\right)-F_{\nu}(z)+z . \tag{3.11}
\end{align*}
$$

Proof. These are straightforward consequences of (3.3), (3.4) and (3.6), (3.9).
It is natural that the free additive convolution and the R-transform appear in the context of spectral theory of free product graphs. It seems less so in the case of the other three types of convolutions. However, we will show that one can decompose the free product of graphs using the products of graphs associated with these convolutions. Let us give the definitions of these products.

Definition 3.1. The comb product of rooted graphs $\left(\mathcal{G}_{1}, e_{1}\right)$ and $\left(\mathcal{G}_{2}, e_{2}\right)$ is the rooted graph $\left(\mathcal{G}_{1} \triangleright \mathcal{G}_{2}, e\right)$ obtained by attaching a copy of $\mathcal{G}_{2}$ by its root $e_{2}$ to each vertex of $\mathcal{G}_{1}$, where we denote by $e$ the vertex obtained by identifying $e_{1}$ and $e_{2}$. If no confusion arises, we denote the comb product by $\mathcal{G}_{1} \triangleright \mathcal{G}_{2}$. If we identify its set of vertices with $V_{1} \times V_{2}$, then its root is identified with $e_{1} \times e_{2}$.

Note that the comb product of rooted graphs is not commutative and it depends on the choice of the root. Let us also remark that the definition given in Ref. 1 is equivalent to the one given above, except that in our definition the information about the role of the root $e_{2}$ in the glueing is encoded in the definition of the rooted graph ( $\mathcal{G}_{2}, e_{2}$ ). Moreover, as our product is taken in the (natural) category of rooted graphs, we define the root of the comb product to be $e$, which makes the product associative.

Theorem 3.1. ${ }^{1}$ Let $\left(\mathcal{G}_{1}, e_{1}\right)$ and $\left(\mathcal{G}_{2}, e_{2}\right)$ be rooted graphs with spectral distributions $\mu$ and $\nu$, respectively. Then, the adjacency matrix of their comb product can be decomposed as

$$
\begin{equation*}
A\left(\mathcal{G}_{1} \triangleright \mathcal{G}_{2}\right)=A^{(1)}+A^{(2)}, \tag{3.12}
\end{equation*}
$$

where $A^{(1)}$ and $A^{(2)}$ are monotone independent with respect to $\varphi(\cdot)=\langle\cdot \delta(e), \delta(e)\rangle$. Moreover, the spectral distribution of $\left(\mathcal{G}_{1} \triangleright \mathcal{G}_{2}, e\right)$ is given by $\mu \triangleright \nu$.

Definition 3.2. The star product of $\left(\mathcal{G}_{1}, e_{1}\right)$ and $\left(\mathcal{G}_{2}, e_{2}\right)$ is the graph $\left(\mathcal{G}_{1} \star \mathcal{G}_{2}, e\right)$ obtained by attaching a copy of $\mathcal{G}_{2}$ by its root $e_{2}$ to the root $e_{1}$ of $\mathcal{G}_{1}$, where we denote by $e$ the vertex obtained by identifying $e_{1}$ and $e_{2}$. If no confusion arises, we also denote the star product by $\mathcal{G}_{1} \star \mathcal{G}_{2}$. If we identify its set of vertices with $V_{1} \star V_{2}:=\left(V_{1} \times\left\{e_{2}\right\}\right) \cup\left(\left\{e_{1}\right\} \times V_{2}\right)$, then its root is identified with $e_{1} \times e_{2}$.

Theorem 3.2. ${ }^{16,29}$ Let $\left(\mathcal{G}_{1}, e_{1}\right)$ and $\left(\mathcal{G}_{2}, e_{2}\right)$ be rooted graphs with spectral distributions $\mu$ and $\nu$, respectively. Then, the adjacency matrix of their star product can be decomposed as

$$
\begin{equation*}
A\left(\mathcal{G}_{1} \star \mathcal{G}_{2}\right)=A^{(1)}+A^{(2)}, \tag{3.13}
\end{equation*}
$$

where $A^{(1)}$ and $A^{(2)}$ are boolean independent with respect to $\varphi$, where $\varphi(\cdot)=$ $\langle\cdot \delta(e), \delta(e)\rangle$. Moreover, the spectral distribution of $\left(\mathcal{G}_{1} \star \mathcal{G}_{2}, e\right)$ is given by $\mu \uplus \nu$.

## 4. Orthogonal Product of Graphs

Let us introduce now a new basic product of rooted graphs called "orthogonal", which is related to the orthogonal convolution introduced in Ref. 20. Using this new product, together with the comb product of graphs (or, the star product of graphs), one can construct their free product $\mathcal{G}_{1} * \mathcal{G}_{2}$, using copies of $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$. This is an application of the more general theory of constructing the free additive convolution from the orthogonal and monotone (or, orthogonal and boolean) convolutions given in Ref. 20 and we use the results contained there.

Definition 4.1. The orthogonal product of two rooted graphs $\left(\mathcal{G}_{1}, e_{1}\right)$ and $\left(\mathcal{G}_{2}, e_{2}\right)$ is the rooted graph $\left(\mathcal{G}_{1} \vdash \mathcal{G}_{2}, e\right)$ obtained by attaching a copy of $\mathcal{G}_{2}$ by its root $e_{2}$ to each vertex of $\mathcal{G}_{1}$ but the root $e_{1}$, where $e$ is taken to be equal to $e_{1}$. If its set of vertices is identified with $V_{1} \vdash V_{2}:=\left(V_{1}^{0} \times V_{2}\right) \cup\left\{e_{1} \times e_{2}\right\}$, then $e$ is identified with $e_{1} \times e_{2}$.

It is worth noting that the orthogonal product of graphs resembles their comb product. The difference is that in the comb product the second graph is glued by its root to all vertices of the first graph, whereas in the orthogonal product the second graph is glued to all vertices but the root of the first graph. An example of the orthogonal product of graphs is given in Fig. 1.


Fig. 1. Orthogonal product $\mathcal{G}_{1} \vdash \mathcal{G}_{2}$.

The notion of the orthogonal product of graphs is related to the concept of orthogonal subalgebras introduced in Ref. 20. By a non-unital subalgebra of a unital algebra $\mathcal{A}$ we understand a subalgebra which does not contain the unit of $\mathcal{A}$.

Definition 4.2. Let $(\mathcal{A}, \varphi, \psi)$ be a unital algebra with a pair of linear normalized functionals and let $\mathcal{A}_{1}$ and $\mathcal{A}_{2}$ be non-unital subalgebras of $\mathcal{A}$. We say that $\mathcal{A}_{2}$ is orthogonal to $\mathcal{A}_{1}$ with respect to $(\varphi, \psi)$ if
(i) $\varphi\left(b a_{2}\right)=\varphi\left(a_{1} b\right)=0$,
(ii) $\varphi\left(w_{1} a_{1} b a_{2} w_{2}\right)=\psi(b)\left(\varphi\left(w_{1} a_{1} a_{2} w_{2}\right)-\varphi\left(w_{1} a_{1}\right) \varphi\left(a_{2} w_{2}\right)\right)$
for any $a_{1}, a_{2} \in \mathcal{A}_{1}, b \in \mathcal{A}_{2}$ and any elements $w, v$ of the algebra $\operatorname{alg}\left(\mathcal{A}_{1}, \mathcal{A}_{2}\right)$ generated by $\mathcal{A}_{1}$ and $\mathcal{A}_{2}$. We say that the pair $(a, b)$ of elements of $\mathcal{A}$ is orthogonal with respect to $(\varphi, \psi)$ if the algebra generated by $a \in \mathcal{A}$ is orthogonal to the algebra generated by $b \in \mathcal{A}$.

In analogy to Theorems 3.1-3.2, one can decompose the adjacency matrix of the orthogonal product of graphs. The proof is based on the tensor product realization of orthogonal subalgebras. ${ }^{20}$ Note that tensor product realizations of noncommutative random variables, originated in Ref. 17 for boolean, $m$-free and free random variables, (see also Ref. 19), are especially useful in the context of graph products since the projections introduced in this scheme tell us how the graphs should be glued together. This technique was later used in a number of papers. ${ }^{111,18,29}$

Theorem 4.1. Let $\left(\mathcal{G}_{1}, e_{1}\right)$ and $\left(\mathcal{G}_{2}, e_{2}\right)$ be rooted graphs with spectral distributions $\mu$ and $\nu$, respectively. Then, the adjacency matrix of their orthogonal product can be decomposed as

$$
\begin{equation*}
A\left(\mathcal{G}_{1} \vdash \mathcal{G}_{2}\right)=A^{(1)}+A^{(2)} \tag{4.1}
\end{equation*}
$$

where the pair $\left(A^{(1)}, A^{(2)}\right)$ is orthogonal with respect to $(\varphi, \psi)$, where $\varphi$ and $\psi$ are states associated with vectors $\delta(e), \delta(v) \in l^{2}\left(V_{1} \vdash V_{2}\right)$ and $v \in V_{1}^{0}$. Moreover, the spectral distribution of $\left(\mathcal{G}_{1} \vdash \mathcal{G}_{2}, e\right)$ is given by $\mu \vdash \nu$.

Proof. In order to prove the decomposition, it is convenient to identify the adjacency matrix of $\mathcal{G}_{1} \vdash \mathcal{G}_{2}$ with the sum

$$
A\left(\mathcal{G}_{1} \vdash \mathcal{G}_{2}\right)=A_{1} \otimes P_{\xi_{2}}+P_{\xi_{1}}^{\perp} \otimes A_{2}
$$

on the Hilbert space $l^{2}\left(V_{1} \vdash V_{2}\right) \subset l^{2}\left(V_{1} \times V_{2}\right) \cong l^{2}\left(V_{1}\right) \otimes l^{2}\left(V_{2}\right)$, where $A_{i}$ is the adjacency matrix of $\mathcal{G}_{i}$ and $P_{\xi_{1}}^{\perp}=1-P_{\xi_{1}}$, with $P_{\xi_{i}}$ denoting the projection onto $\mathbb{C} \xi_{i}$, where $\xi_{i}=\delta\left(e_{i}\right)$ and $i=1$, 2. Projection $P_{\xi_{1}}^{\perp}$ indicates that graph $\mathcal{G}_{2}$ should be glued to all vertices of $\mathcal{G}_{1}$ but the root, whereas projection $P_{\xi_{2}}$ indicates that graph $\mathcal{G}_{1}$ should be glued only to vertex $e_{2}$ of $\mathcal{G}_{2}$, which reproduces Definition 4.1. It remains to take $\varphi$ and $\psi$ to be the states associated with unit vectors $\delta\left(e_{1}\right) \times \delta\left(e_{2}\right)$ and $\delta(v) \times \delta\left(e_{2}\right)$, respectively, where $v$ is an arbitrary vertex from $V_{1}^{0}$. Now, in view of Theorem 4.1 of Ref. 20, the above summands form a pair of orthogonal elements
(of the algebra they generate) with respect to the pair of states $(\varphi, \psi)$. The proof consists of checking (i)-(ii) of Definition 4.2 and in the case of graphs it is very similar to the general convolution case. ${ }^{20}$

It follows from Corollary 4.2 of Ref. 20 that the spectral distribution of such a sum in the state $\varphi$ is equal to the orthogonal convolution $\mu \vdash \nu$. Nevertheless, we choose to show this fact here since we can present a proof which nicely exhibits the relation between the comb product and the orthogonal product. Namely, recall that $\mathcal{G}_{1} \vdash \mathcal{G}_{2}$ differs from $\mathcal{G}_{1} \triangleright \mathcal{G}_{2}$ by the fact that no copy of $\mathcal{G}_{2}$ is glued to the vertex $e_{1}$. Therefore, one can obtain $\mathcal{G}_{1} \triangleright \mathcal{G}_{2}$ by glueing $\mathcal{G}_{1} \vdash \mathcal{G}_{2}$ and $\mathcal{G}_{2}$ at their roots, which corresponds to their star product. Therefore, $\mathcal{G}_{1} \triangleright \mathcal{G}_{2}=\left(\mathcal{G}_{1} \vdash \mathcal{G}_{2}\right) \star \mathcal{G}_{2}$, which leads to the formula $\mu \triangleright \nu=\sigma \uplus \nu$ for spectral distributions, where $\sigma$ is the spectral distribution of $\mathcal{G}_{1} \vdash \mathcal{G}_{2}$. Using transforms, we get $F_{\mu}\left(F_{\nu}(z)\right)=F_{\sigma}(z)+F_{\nu}(z)-z$, which, in view of (3.11), gives our assertion.

Example 4.1. Let us apply Theorem 4.1 to the orthogonal product in Fig. 1. We have

$$
G_{\mu}(z)=\frac{z^{2}-1}{z\left(z^{2}-2\right)}, \quad G_{\nu}(z)=\frac{z^{2}-2}{z\left(z^{2}-3\right)}
$$

and therefore,

$$
K_{\mu}(z)=\frac{1}{z-\frac{1}{z}}, \quad F_{\nu}(z)=z-\frac{1}{z-\frac{2}{z}} .
$$

In view of Theorem 4.1 and (3.9), we obtain the explicit formula for the Cauchy transform $G_{\mu \vdash \nu}(z)$. Algebraic calculations lead to the continued fraction representation of $G_{\mu \vdash \nu}(z)$ associated with the sequences of Jacobi coefficients $\omega=\left(\omega_{n}\right)=$ $(1,2,3 / 2,5 / 6,4 / 15,12 / 5,0, \ldots)$ and $\alpha=\left(\alpha_{n}\right)=(0,0, \ldots)$. The corresponding measure is a discrete measure consisting of seven atoms (since their explicit values and corresponding masses are rather complicated, we do not give them here).

## 5. Free Decomposition of the Free Product of Graphs

In this section we recall after Refs. 39 and 10 the definition of the free product of rooted graphs and we show that the adjacency matrix of the free product of a finite family of rooted graphs is the sum of freely independent copies of the adjacency matrices of the factors of the product. Essentially, this fact is a natural consequence of free probability and one just needs to adapt the proof of Voiculescu. We also define a corresponding approximating sequence of $m$-free products.

Consider rooted graphs $\left(\mathcal{G}_{i}, e_{i}\right)=\left(V_{i}, E_{i}, e_{i}\right)$, where $i \in I$ and $I$ is a finite index set, and denote $V_{i}^{0}=V_{i} \backslash\left\{e_{i}\right\}$. By the free product of rooted sets $\left(V_{i}, e_{i}\right), i \in I$, we understand the rooted set ( $\left.*_{i \in I} V_{i}, e\right)$, where

$$
*_{i \in I} V_{i}=\{e\} \cup\left\{v_{1} v_{2} \cdots v_{m} ; v_{k} \in V_{i_{k}}^{0} \text { and } i_{1} \neq i_{2} \neq \cdots \neq i_{n}, m \in \mathbb{N}\right\}
$$

and $e$ is the empty word. For notational convenience, we will sometimes use words containing roots $e_{k}$ but then we shall always understand that $w e_{k}=e_{k} w \equiv w$, where $w \in *_{i \in I} V_{i}$, thus any $e_{k}$ will be treated as the "unit", or the empty word. We are ready to give the definition of the free product of graphs.

Definition 5.1. By the free product of rooted graphs $\left(\mathcal{G}_{i}, e_{i}\right), i \in I$, we understand the rooted graph $\left(*_{i \in I} \mathcal{G}_{i}, e\right)$ with the set of vertices $*_{i \in I} V_{i}$ and the set of edges $*_{i \in I} E_{i}$ consisting of pairs of vertices from $*_{i \in I} V_{i}$ of the form

$$
*_{i \in I} E_{i}=\left\{\left\{v u, v^{\prime} u\right\}:\left\{v, v^{\prime}\right\} \in \bigcup_{i \in I} E_{i} \text { and } u, v u, v^{\prime} u \in *_{i \in I} V_{i}\right\} .
$$

We denote this product by $*_{i \in I}\left(\mathcal{G}_{i}, e_{i}\right)$ or simply $*_{i \in I} \mathcal{G}_{i}$ if no confusion arises.
The most intuitive construction of the free product of graphs is given by some inductive procedure which gives a sequence of growing graphs whose inductive limit is the free product of graphs. In fact, one natural procedure was given in Ref. 39, where it was one of the equivalent definitions of the free product of graphs. Interestingly enough, this procedure gives a sequence of iterates indexed by $m \in \mathbb{N}$ which corresponds to the $m$-free product of states introduced in Ref. 17. This leads us to the following formal definition.

Definition 5.2. By the $m$-free product of rooted graphs $\left(\mathcal{G}_{i}, e_{i}\right), i \in I$, we understand the subgraph $\left(*_{i \in I}^{(m)} \mathcal{G}_{i}, e\right)$ of $\left(*_{i \in I} \mathcal{G}_{i}, e\right)$ obtained by restricting the set of vertices to words $w$ of length $|w| \leq m$.

Example 5.1. Consider two "segments" $\mathcal{G}_{1} \cong \mathbb{Z}_{2}$ and $\mathcal{G}_{2} \cong \mathbb{Z}_{2}$. They are graphs consisting of one edge $x \sim e_{1}$ and $y \sim e_{2}$. Then

$$
\begin{gathered}
V_{1} * V_{2}=\{e, x, y, x y, y x, x y x, y x y, \ldots\} \\
E_{1} * E_{2}=\{\{e, x\},\{e, y\},\{x, y x\},\{y, x y\},\{y x, x y x\},\{x y, y x y\}, \ldots\}
\end{gathered}
$$

and it is easy to see that $\mathcal{G}_{1} * \mathcal{G}_{2} \cong \mathbb{Z}$, where $\mathbb{Z}$ denotes the two-way infinite path (or, one-dimensional integer lattice) with the root at 0 . Truncations of this product to words of length $\leq m$ give $\mathcal{G}_{1} *^{m} \mathcal{G}_{2}$.

In order to give the explicit form of the adjacency matrix of the free product of graphs, consider the Hilbert space $\mathcal{H}_{I}=l^{2}\left(*_{i \in I} V_{i}\right)$ spanned by $\delta(e)$ and vectors of the form

$$
\delta(w) \text { for } w=v_{1} v_{2} \cdots v_{n} \in *_{i \in I} V_{i}
$$

and let $\varphi$ be the vacuum expectation on $B(\mathcal{H})$ given by $\varphi(T)=\langle T \delta(e), \delta(e)\rangle$. We have

$$
\left(\mathcal{H}_{I}, \delta(e)\right) \cong *_{i \in I}\left(l^{2}\left(V_{i}\right), \delta\left(e_{i}\right)\right)
$$

where the RHS is understood as the free product of Hilbert spaces with distinguished unit vectors. ${ }^{33}$

In the sequel we will need the following subsets of $*_{i \in I} V_{i}$ :

$$
W_{j}(n)=\left\{v_{1} v_{2} \cdots v_{n} \in *_{i \in I} V_{i}: v_{1} \notin V_{j}^{0}\right\},
$$

where $n \in \mathbb{N}$. Thus, $W_{j}(n)$ is the subset consisting of words of length $n$ which do not begin with a letter from $V_{j}^{0}$. We set

$$
W_{j}=\bigcup_{n=0}^{\infty} W_{j}(n)
$$

with $W_{j}(0)=\{e\}$ for every $j$.
Definition 5.3. Let $A_{i}$ denote the adjacency matrix of the rooted graph $\left(\mathcal{G}_{i}, e_{i}\right)$, where $i \in I$. Let us define their copies in $*_{i \in I}\left(\mathcal{G}_{i}, e_{i}\right)$ by the formulas

$$
\left(A_{i}(n)\right)_{w, w^{\prime}}= \begin{cases}1 & \text { if }\left\{w, w^{\prime}\right\}=\left\{x u, x^{\prime} u\right\} \text { for }\left\{x, x^{\prime}\right\} \in E_{j} \text { and } u \in W_{j}(n-1) \\ 0 & \text { otherwise }\end{cases}
$$

where $n \in \mathbb{N}$. By $P_{i}(n)$ we denote the canonical projection of $\mathcal{H}_{I}$ onto $l^{2}\left(W_{i}(n)\right)$ for $n \geq 1$, with $P_{i}(0)$ denoting the projection onto $l^{2}(e)=\mathbb{C} \delta(e)$ for every $i \in I$.

Theorem 5.1. The adjacency matrix $A\left(*_{i \in I} \mathcal{G}_{i}\right)$ of the free product of graphs admits a decomposition of the form $A\left(*_{i \in I} \mathcal{G}_{i}\right)=\sum_{i \in I} A^{(i)}$, where

$$
\begin{equation*}
A^{(i)}=\sum_{n=1}^{\infty} A_{i}(n)=\sum_{n=1}^{\infty} A_{i} P_{i}(n-1) \tag{5.1}
\end{equation*}
$$

are free with respect to the vacuum expectation $\varphi$ and the action of $A_{i}$ in the second sum is given by $A_{i} \delta(x u)=\delta\left(x^{\prime} u\right)$ whenever $\left\{x, x^{\prime}\right\} \in E_{i}(i \in I)$. Moreover, the series is strongly convergent for every $i \in I$.

Proof. First, let us observe that using local finitness of $\mathcal{G}_{i}$, we can write

$$
A^{(i)} \delta(w)=\sum_{\substack{w^{\prime}=x^{\prime} u \\\left\{x, x^{\prime}\right\} \in E_{i}}} \delta\left(w^{\prime}\right)=A_{i}(n) \delta(w)
$$

whenever $w=x u$, where $u \in W_{i}(n-1)$ and we allow $x, x^{\prime}$ to be arbitrary vertices from $V_{i}$, thus if $x=e_{1}$, we have $e_{1} u \equiv u$. Moreover, observe that $A_{i}(m) \delta(w)=0$ for such $w$ for any $m \neq n$. Writing

$$
\delta(w)= \begin{cases}\delta(x) \otimes \delta(u) & \text { whenever } w=x u \text { and } x \in V_{i}^{0} \\ \delta(u) & \text { if } w=e_{i} u\end{cases}
$$

we obtain

$$
\begin{aligned}
A^{(i)} \delta(w) & =\sum_{\substack{w^{\prime}=x^{\prime} u \\
\left\{x, x^{\prime}\right\} \in E_{i}, x^{\prime} \neq e_{i}}} \delta\left(x^{\prime}\right) \otimes \delta(u)+\mathbb{1}_{\left\{\left\{x, e_{i}\right\} \in E_{i}\right\}} \delta(u) \\
& \left.=\left(A_{i} \delta(x)\right)^{0} \otimes \delta(u)+\left\langle A_{i} \delta(x), \delta\left(e_{i}\right)\right)\right\rangle \delta(u)
\end{aligned}
$$

if $w=x u, u \in W_{i}$, where $1_{\{z \in A\}}=1$ if and only if $z \in A$ and otherwise is zero.

We can observe now that $A^{(i)}=\lambda\left(A_{i}\right)$, where $\lambda$ denotes the free product representation of the free product $\mathbb{C}\left[A_{1}\right] * \mathbb{C}\left[A_{2}\right]$ on $l^{2}\left(*_{i=1}^{n} V_{i}\right)$ in the sense of Avitzour ${ }^{3}$ and Voiculescu. ${ }^{33}$ Therefore, the $A^{(i)}$ are free with respect to $\varphi$.

As a consequence of the decomposition theorem, one can use free additive convolutions ${ }^{33}$ to compute spectral distributions of free products of rooted graphs in terms of spectral distributions of the factors. One also obtains asymptotic spectral properties of free powers $(\mathcal{G}, e)^{* n}$. The proofs of these starightforward facts are omitted.

Corollary 5.1. Let $A_{i}$ be the adjacency matrix of $\left(\mathcal{G}_{i}, e_{i}\right), i \in I=\{1, \ldots, n\}$, and let $\mu_{i}$ denote its spectral distribution, where $1 \leq i \leq n$. Then the spectral distribution of $A\left(*_{i=1}^{n} \mathcal{G}_{i}\right)$ in the state $\varphi$ is given by $\mu=\mu_{1} \boxplus \mu_{2} \boxplus \cdots \boxplus \mu_{n}$.

Corollary 5.2. Let $A$ be the adjacency matrix of $(\mathcal{G}, e)$ and let $A^{* n}$ denote the adjacency matrix of $(\mathcal{G}, e)^{* n}$. Then

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \varphi\left(\left(\frac{A^{* n}}{\sqrt{n k(e)}}\right)^{2 m}\right)=c_{m} \tag{5.2}
\end{equation*}
$$

where $c_{m}$ is the mth Catalan number for $m \in \mathbb{N}, c_{0}=1$ and $k(e)$ is the degree of the root $e$. The odd moments vanish.

Remark 5.1. The correspondence between the free product of graphs and free probability (in particular, Corollary 5.2) can be applied to establish a connection between free products of graphs and free additive convolutions of their spectral distributions.

Example 5.2. Using Corollary 5.2, we find spectral distributions for two standard examples of $n$-ary rooted trees and homogeneous trees which will be needed later. Thus, the spectral distributions of $n$-ary trees $\mathbb{T}_{n}$ have Cauchy transforms

$$
\begin{equation*}
G_{\nu_{n}}(z)=\frac{z-\sqrt{z^{2}-4 n}}{2 n} \tag{5.3}
\end{equation*}
$$

with densities given by Wigner laws

$$
\begin{equation*}
d \nu_{n}(x)=\frac{\sqrt{4 n-x^{2}}}{2 \pi n} d x \tag{5.4}
\end{equation*}
$$

with the supports on $[-2 \sqrt{n}, 2 \sqrt{n}]$. In a similar manner, we obtain the Cauchy transforms of the spectral distributions $\mu_{n}$ of homogenous trees $\mathbb{H}_{n}$ of the form

$$
\begin{equation*}
G_{\mu_{n}}(z)=\frac{(2-n) z+n \sqrt{z^{2}-4(n-1)}}{2\left(z^{2}-n^{2}\right)} \tag{5.5}
\end{equation*}
$$

which, with the help of the Stieltjes inversion formula, give the (absolutely continuous) measures with densities

$$
\begin{equation*}
d \mu_{n}(x)=\frac{n \sqrt{4(n-1)-x^{2}}}{2 \pi\left(n^{2}-x^{2}\right)} d x \tag{5.6}
\end{equation*}
$$

supported on $[-2 \sqrt{n-1}, 2 \sqrt{n-1}]$. In particular, the spectral distribution of $\mathbb{H}_{2} \cong$ $\mathbb{Z}$ in the vacuum state $\varphi$ associated with the vertex 0 is the arcsine law $d \mu_{2}(x)=$ $1 /\left(\pi \sqrt{4-x^{2}}\right) d x$.

## 6. Orthogonal Decompositions of Branches

Let us look at the concept of "branches" of the free product of graphs introduced by Quenell. ${ }^{30}$ They correspond to the so-called "subordination functions" studied first by Voiculescu ${ }^{36}$ and Biane. ${ }^{4}$ We rely on the recent general study of the free additive convolution and its decompositions given in Ref. 20, where we refer the reader for the main concepts, like s-freeness, as well as general proofs.

Definition 6.1. Let $\left(V_{i}, e_{i}\right)_{i \in I}$ be a finite family of rooted sets. By the $j$ th subordination branch of $*_{i \in I}\left(V_{i}, e_{i}\right)$, where $j \in I$, we shall understand the rooted set $\left(S_{j}, e\right)$, where

$$
S_{j}=\{e\} \cup\left\{v_{1} v_{2} \cdots v_{m} \in *_{i \in I} V_{i}: v_{m} \in V_{j}^{0}, m \in \mathbb{N}\right\}
$$

is the subset of $*_{i \in I} V_{i}$ consisting of the empty word and words which end with a letter from $V_{j}^{0}$.

Definition 6.2. Let $\left(\mathcal{G}_{i}, e_{i}\right)_{i \in I}$ be a finite family of rooted graphs. By the $j$ th subordination branch of $*_{i \in I}\left(\mathcal{G}_{i}, e_{i}\right)$, where $j \in I$, we shall understand the rooted graph $\left(\mathcal{B}_{j}, e\right)$, where $\mathcal{B}_{j} \equiv \mathcal{B}_{j}\left(\left(\mathcal{G}_{i}\right)_{i \in I}\right)$ is the subgraph of $*_{i \in I} \mathcal{G}_{i}$ restricted to the set $S_{j}$ defined above. As before, we often omit the roots in the notations.

In the case of two graphs, it is easy to see that $\mathcal{G}_{1} * \mathcal{G}_{2}$ consists of two branches, $\mathcal{B}_{1}$ and $\mathcal{B}_{2}$, with common root $e$. The branch $\mathcal{B}_{1}=\mathcal{B}_{1}\left(\mathcal{G}_{1}, \mathcal{G}_{2}\right)$ "begins" with a copy of $\mathcal{G}_{1}$ and the branch $\mathcal{B}_{2}=\mathcal{B}_{2}\left(\mathcal{G}_{1}, \mathcal{G}_{2}\right)$ "begins" with a copy of $\mathcal{G}_{2}$. For instance, in the case of a binary tree $\mathbb{T}_{2}$, the branches $\mathcal{B}_{1}$ and $\mathcal{B}_{2}$ are the left and right "halves" of $\mathbb{T}_{2}$, respectively. However, the $n$-ary tree can itself be viewed as a branch of another free product. Moreover, it is then constructed in a "distance-adapted" manner, i.e. natural truncations of the product lead to natural truncations of the tree (see Example 6.1 and Fig. 2).

Example 6.1. Take two graphs as in Fig. 2 and let $\mathcal{B}_{1}$ and $\mathcal{B}_{2}$ be the subordination branches of the free product $\mathcal{G}_{1} * \mathcal{G}_{2}$. Figure 2 shows that $\mathbb{T}_{2} \cong \mathcal{B}_{1}\left(\cong \mathcal{B}_{2}\right)$.

In a similar way one can obtain the $n$-ary rooted tree as a branch of a free product of two copies of the "fork" graph with $n+1$ vertices $e, x_{1}, \ldots, x_{n}$ (i.e. such that $e \sim x_{k}$ for all $1 \leq k \leq n$ ).

Motivated by Ref. 20, we can view the branches of $\mathcal{G}_{1} * \mathcal{G}_{2}$ as products of rooted graphs. The needed notion of a product corresponds to freeness with subordination, or $s$-freeness, introduced and studied there. Moreover, their adjacency matrices, $A\left(\mathcal{B}_{1}\right)$ and $A\left(\mathcal{B}_{2}\right)$, can be decomposed as the sum of components which are "free with subordination" (or, "s-free") with respect to a pair of states $(\varphi, \psi)$.


Fig. 2. Binary tree $\mathbb{T}_{2} \cong \mathcal{B}_{1}\left(\mathcal{G}_{1} * \mathcal{G}_{2}\right)$.

Definition 6.3. Let $(\mathcal{A}, \varphi, \psi)$ be a unital algebra with a pair of linear normalized functionals. Let $\mathcal{A}_{1}$ be a unital subalgebra of $\mathcal{A}$ and let $\mathcal{A}_{2}$ be a non-unital subalgebra with an "internal" unit $1_{2}$, i.e. $1_{2} b=b=b 1_{2}$ for every $b \in \mathcal{A}_{2}$. We say that the pair $\left(\mathcal{A}_{1}, \mathcal{A}_{2}\right)$ is free with subordination, or simply $s$-free, with respect to $(\varphi, \psi)$ if $\psi\left(1_{2}\right)=1$ and it holds that
(i) $\varphi\left(a_{1} a_{2} \cdots a_{n}\right)=0$ whenever each $a_{j} \in \mathcal{A}_{i_{j}}^{0}$ and $i_{1} \neq i_{2} \neq \cdots \neq i_{n}$
(ii) $\varphi\left(w_{1} 1_{2} w_{2}\right)=\varphi\left(w_{1} w_{2}\right)-\varphi\left(w_{2}\right) \varphi\left(w_{2}\right)$ for any $w_{1}, w_{2} \in \operatorname{alg}\left(\mathcal{A}_{1}, \mathcal{A}_{2}\right)$,
where $\mathcal{A}_{1}^{0}=\mathcal{A}_{1} \cap \operatorname{ker} \varphi$ and $\mathcal{A}_{2}^{0}=\mathcal{A}_{2} \cap \operatorname{ker} \psi$. We say that the pair $(a, b)$ of random variables from $\mathcal{A}$ is $s$-free with respect to $(\varphi, \psi)$ if there exists $1_{2} \in \mathcal{A}$ such that the pair $\left(\mathcal{A}_{1}, \mathcal{A}_{2}\right)$, where $\mathcal{A}_{1}$ is the unital algebra generated by $a$ and $\mathcal{A}_{2}$ is the non-unital algebra generated by $1_{2}$ and $b$, is s-free with respect to $(\varphi, \psi)$.

The notion of s-freeness resembles freeness - in the GNS representation, the corresponding product of Hilbert spaces is the direct sum of $\mathbb{C} \xi$, where $\xi$ is the unit (vacuum) vector, and tensors $\mathcal{H}_{i_{1}} \otimes \mathcal{H}_{i_{2}} \otimes \cdots \otimes \mathcal{H}_{i_{n}}$, where $i_{1} \neq i_{2} \neq \cdots \neq i_{n}=1$. The branches, which in this context replace free products of graphs, can also be decomposed along the lines of Theorem 5.1 and can be called "s-free products" of $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$, or $\mathcal{G}_{2}$ and $\mathcal{G}_{1}$. Using a similar notation, we obtain the following decomposition theorem (cf. Ref. 20).

Theorem 6.1. The adjacency matrix of the branch $\mathcal{B}_{1} \equiv \mathcal{B}_{1}\left(\mathcal{G}_{1}, \mathcal{G}_{2}\right)$ can be decomposed as the sum $A\left(\mathcal{B}_{1}\right)=A^{(1)}+A^{(2)}$, where the strongly convergent series

$$
\begin{equation*}
A^{(1)}=\sum_{n \text { odd }} A_{1}(n), \quad A^{(2)}=\sum_{n \text { even }} A_{2}(n), \tag{6.1}
\end{equation*}
$$

are s-free with respect to $(\varphi, \psi)$, where $\varphi(\cdot)=\langle\cdot \delta(e), \delta(e)\rangle$ and $\psi(\cdot)=\langle\cdot \delta(v), \delta(v)\rangle$ for any $v \in V_{1}^{0}$. An analogous decomposition holds for the branch $\mathcal{B}_{2}\left(\mathcal{G}_{1} * \mathcal{G}_{2}\right)$ with the summations over odd and even $n$ interchanged.

Proof. We refer the reader to Ref. 20, where it was shown, in a general Hilbert space setting, that sums of operators of the above type are s-free with respect to $(\varphi, \psi)$ (one has to verify conditions (i)-(ii) of Definition 6.3, and in the case of graphs, it is basically the same proof).

In order to "decompose completely" the branches, by which we mean to decompose them in terms of graphs $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$, we will interpret (6.1) in terms of an inductive limit of a sequence of graphs which resembles (but is not the same as) the sequence of $m$-free products approximating the free product. In this fashion we will obtain the "complete" orthogonal decomposition of branches given by the following theorem.

Theorem 6.2. The branch $\mathcal{B}_{1}$ is the inductive limit of the sequence $\left(\mathcal{G}_{1} \vdash_{m} \mathcal{G}_{2}\right)_{m \in \mathbb{N}}$ given by the recursion

$$
\mathcal{G}_{1} \vdash_{1} \mathcal{G}_{2}=\mathcal{G}_{1} \vdash \mathcal{G}_{2}, \quad \mathcal{G}_{1} \vdash_{m} \mathcal{G}_{2}=\mathcal{G}_{1} \vdash\left(\mathcal{G}_{2} \vdash_{m-1} \mathcal{G}_{2}\right),
$$

where $m>1$. An analogous statement holds for the branch $\mathcal{B}_{2}$.
Proof. Without loss of generality, consider branch $\mathcal{B}_{1}$. Our sequence of iterates will remind the inductive way of defining the free product of graphs given in Ref. 39, although it is not symmetric with respect to $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$. Recall that $\mathcal{B}_{i}$ "begins" with a copy of $\mathcal{G}_{i}$. Therefore, let $\mathcal{B}_{1}(0)$ be equal to $\mathcal{G}_{1}$ and choose its root to be $e_{1}$. To get $\mathcal{B}_{1}(1)$, to every vertex of $\mathcal{G}_{1}$ but the root we glue by its root a copy of $\mathcal{B}_{2}(0)$. In such a graph we again choose the root $e_{1}$. This gives a rooted graph $\left(\mathcal{B}_{1}(1), e_{1}\right)$, which is, in fact, $\mathcal{G}_{1} \vdash \mathcal{G}_{2}$. In a similar fashion we obtain $\left(\mathcal{B}_{2}(1), e_{2}\right)$. Now, note that the $m$ th approximant of the branch $\mathcal{B}_{1}$ is obtained by glueing by its root a copy of $\left(\mathcal{B}_{2}(m-1), e_{2}\right)$ to every vertex of $\left(\mathcal{G}_{1}, e_{1}\right)$ but the root. In other words, we obtain

$$
\mathcal{B}_{1}(m)=\mathcal{G}_{1} \vdash \mathcal{B}_{2}(m-1) \quad \text { and } \quad \mathcal{B}_{2}(m)=\mathcal{G}_{2} \vdash \mathcal{B}_{1}(m-1)
$$

for $m \geq 1$. It is clear that the inductive limits of our iterates give the branches, namely

$$
\mathcal{B}_{i}=\bigcup_{m \geq 0} \mathcal{B}_{i}(m)
$$

for $i=1,2$ (with the root $e_{i}$ ), and this proves the assertion.
In order to obtain spectral distributions of the branches, one takes a sequence of alternating iterates of orthogonal convolutions - this method was introduced in Ref. 20, but here, in the graph context, is especially appealing and easy to justify.

Corollary 6.1. If $\mu$ and $\nu$ are spectral distributions of $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$, respectively, then the spectral distribution of $\mathcal{G}_{1} \vdash_{m} \mathcal{G}_{2}$ is given by $\mu \vdash_{m} \nu$, where the sequence $\left(\mu \vdash_{m} \nu\right)_{m \in \mathbb{N}}$ of distributions is given by the recursion

$$
\mu \vdash_{1} \nu=\mu \vdash \nu, \quad \mu \vdash_{m} \nu=\mu \vdash\left(\nu \vdash_{m-1} \mu\right),
$$

where $m>1$. If $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$ are uniformly locally finite, the spectral distribution of the branch $\mathcal{B}_{1}$ is given by the weak limit $\mu \mathbb{\boxtimes} \nu:=w-\lim _{m \rightarrow \infty}\left(\mu \vdash_{m} \nu\right)$. An analogous statement holds for the branch $\mathcal{B}_{2}$.

Proof. The spectral distribution of $\mathcal{G}_{1} \vdash_{m} \mathcal{G}_{2}$ is given by $\mu \vdash_{m} \nu$ by Theorem 4.1. Now, observe that moments of the same order $k$, where $k \leq 2 m$, in all graphs $\mathcal{G}_{1} \vdash_{n} \mathcal{G}_{2}$ (computed with respect to the root $e$ ), with $n \geq m$, are equal. This is because in the orthogonal product of graphs no copy of the second graph is glued to the root of the first graph and therefore, the distance from the root $e$ in $\mathcal{G}_{1} \vdash_{m} \mathcal{G}_{2}$ at which the graph differs from $\mathcal{G}_{1} \vdash_{m-1} \mathcal{G}_{2}$ is equal to $m+1$. Therefore, the moments of $\mu \vdash_{m} \nu$ converge to the corresponding moments of the spectral distributions of $\mathcal{B}_{1}$. If $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$ are uniformly locally finite, this implies weak convergence of measures.

Corollary 6.2. Under the assumptions of Corollary 6.1, the $K$-transform of $\mu \boxminus \nu$ can be expressed as

$$
K_{\mu \boxplus \nu}(z)=K_{\mu}\left(z-K_{\nu}\left(z-K_{\mu}\left(z-K_{\nu}(\cdots)\right)\right)\right),
$$

where the right-hand side is understood as the uniform limit on compact subsets of the complex upper half-plane. The K-transform of $\mu \vdash_{m} \nu$ is obtained by a truncation of the above formula to $m+1$ alternating transforms.

Proof. Since weak convergence of measures implies uniform convergence of the Cauchy transform on compact subsets of the complex upper half-plane, the assertion follows from a repeated application of (3.9) and Corollary 6.1.

Example 6.2. Consider two rooted graphs $\mathcal{G}_{1}, \mathcal{G}_{2}$, whose spectral distributions $\mu$, $\nu$ are associated with reciprocal Cauchy transforms of the form

$$
F_{\mu}(z)=z-\alpha_{0}-\frac{\omega_{0}}{z-\alpha_{1}}, \quad F_{\nu}(z)=z-\beta_{0}-\frac{\gamma_{0}}{z-\beta_{1}}
$$

respectively (this includes $\mathbb{K}_{n}$ and $\mathbb{F}_{n}$, whose free products were studied by other authors and also in Sec. 10). From Corollary 6.2 we easily obtain the K-transform

$$
K_{\mu \boxplus \nu}(z)=\alpha_{0}+\frac{\omega_{0}}{z-\alpha_{1}-\beta_{0}-\frac{\gamma_{0}}{z-\alpha_{0}-\beta_{1}-\frac{\omega_{0}}{z-\alpha_{1}-\beta_{0}-\frac{\gamma_{0}}{\cdots}}}}
$$

and thus, in view of (3.3), the distribution $\mu \boxminus \nu$ of branch $\mathcal{B}_{1}$ is associated with the sequences of Jacobi parameters

$$
\alpha=\left(\alpha_{0}, \alpha_{1}+\beta_{0}, \alpha_{0}+\beta_{1}, \alpha_{1}+\beta_{0}, \ldots\right), \quad \omega=\left(\omega_{0}, \gamma_{0}, \omega_{0}, \gamma_{0}, \ldots\right)
$$

which correspond to the so-called mixed periodic Jacobi continued fraction. ${ }^{13}$ For details on the corresponding measures, see Ref. 13. In particular, if $\mathcal{G}_{1}=\mathcal{G}_{2}=\mathbb{K}_{2}$ (3-vertex complete graph), we have $\alpha_{0}=\beta_{0}=0, \alpha_{1}=\beta_{1}=1 \omega_{0}=\gamma_{0}=2$, which gives $\mu \boxtimes \nu$ associated with the sequences of Jacobi parameters $\alpha=(0,1,1, \ldots)$ and $\omega=(2,2, \ldots)$. Its Cauchy transform is

$$
G(z)=\frac{z+1-\sqrt{z^{2}-2 z-7}}{2 z+4}
$$

and the measure has density $d \mu(x)=\sqrt{7+2 x-x^{2}} /(\pi(2 x+4))$ on the interval $[1-2 \sqrt{2}, 1+2 \sqrt{2}]$.

## 7. "Complete" Decompositions of Free Products

In this section we derive new decompositions of the free product of graphs, which are based on the orthogonal decomposition of branches of Sec. 6. We rely on the general theory of the decompositions of the free additive convolution ${ }^{20}$ and apply it to the context of graph products.

We start from two lemmas, which rephrase the results of Quenell using the language of quantum probability. This reduces certain proofs presented in Ref. 30 to the basic properties of monotone and boolean convolutions.

Lemma 7.1. The free product of rooted graphs admits the decomposition

$$
\begin{equation*}
\mathcal{G}_{1} * \mathcal{G}_{2} \cong \mathcal{B}_{1} \star \mathcal{B}_{2} \tag{7.1}
\end{equation*}
$$

which we call the star decomposition of $\mathcal{G}_{1} * \mathcal{G}_{2}$.
Proof. Notice that $V_{1} * V_{2}=S_{1} \cup S_{2}$. Moreover, it follows immediately from Definition 6.2 that the free product $\mathcal{G}_{1} * \mathcal{G}_{2}$ is obtained by glueing together the branches $\mathcal{B}_{1}$ and $\mathcal{B}_{2}$ at their roots. From the definition of the star product we know that this is the star product of $\mathcal{B}_{1}$ and $\mathcal{B}_{2}$.

Lemma 7.2. The free product of rooted graphs admits the decompositions

$$
\begin{equation*}
\mathcal{G}_{1} * \mathcal{G}_{2} \cong \mathcal{G}_{1} \triangleright \mathcal{B}_{2} \cong \mathcal{G}_{2} \triangleright \mathcal{B}_{1} \tag{7.2}
\end{equation*}
$$

which we call the comb decompositions of $\mathcal{G}_{1} * \mathcal{G}_{2}$.
Proof. Without loss of generality, consider the first relation. Observe that we can view the branch $\mathcal{B}_{1}$ as one replica of graph $\mathcal{G}_{1}$, to which we glue "orthogonally" replicas of branch $\mathcal{B}_{2}$. Therefore

$$
\mathcal{B}_{1} \cong \mathcal{G}_{1} \vdash \mathcal{B}_{2} \quad \text { and } \quad \mathcal{B}_{2} \cong \mathcal{G}_{2} \vdash \mathcal{B}_{1} .
$$

By Lemma 7.1, we can obtain $\mathcal{G}_{1} * \mathcal{G}_{2}$ by glueing $\mathcal{B}_{1}$ and $\mathcal{B}_{2}$ together at their roots identified with $e$. Equivalently, (one replica of) $\mathcal{B}_{2}$ is glued to $e$ and branch $\mathcal{B}_{1}$ is replaced by $\mathcal{G}_{1} \vdash \mathcal{B}_{2}$, which means that a replica of $\mathcal{B}_{2}$ is glued to every vertex
$v \in V_{1}^{0} \subset V_{1} * V_{2}$. In other words, a replica of $\mathcal{B}_{2}$ is glued to every vertex of $V_{1}^{0} \cup\{e\} \cong V_{1}$, which gives the comb product of $\mathcal{G}_{1}$ and $\mathcal{B}_{2}$, which proves our assertion.

Corollary 7.1. The following relations hold:

$$
\begin{aligned}
& F_{\mu \boxplus \nu}(z)=F_{\mu}\left(F_{\nu \boxplus \mu}(z)\right)+F_{\nu}\left(F_{\mu \boxplus \nu}(z)\right)-z, \\
& F_{\mu \boxplus \nu}(z)=F_{\mu}\left(F_{\nu \boxplus \mu}(z)\right)=F_{\nu}\left(F_{\mu \boxplus \nu}(z)\right),
\end{aligned}
$$

where $\mu$ and $\nu$ are spectral distributions of rooted graphs $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$, respectively.
Proof. These are straightforward consequences of Lemmas 7.1-7.2 and (3.8), (3.10).

Remark 7.1. In terms of moment generating functions $M_{\mu}(z)$, related to $F_{\mu}(z)$ by (3.2) and (3.4), formulas of Corollary 7.1 give the results of Quenell ${ }^{30}$ (see also Ref. 24). In our notation, $M_{\mu}(z)$ correspond to return generating functions of type $R_{e}(z)$ (from which the first return generating functions of type $S_{e}(z)$ are easily obtained). Moreover, these results can be easily generalized to a finite number of rooted graphs.

Let us now use the "complete" orthogonal decomposition of branches (Theorem 6.2) and Lemmas 7.1-7.2 to derive "complete" decompositions of the free product of graphs. We begin with a decomposition of $m$-free products.

Theorem 7.1. Let $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$ be rooted graphs with spectral distributions $\mu$ and $\nu$. Then their $m$-free product admits the decomposition

$$
\mathcal{G}_{1} *^{(m)} \mathcal{G}_{2}=\left(\mathcal{G}_{1} \vdash_{m} \mathcal{G}_{2}\right) \star\left(\mathcal{G}_{2} \vdash_{m} \mathcal{G}_{1}\right)
$$

called the star-orthogonal decomposition, and its spectral distribution is given by the $m$-free convolution $\mu \boxplus_{m} \nu:=\left(\mu \vdash_{m} \nu\right) \uplus\left(\nu \vdash_{m} \mu\right)$.

Proof. The proof consists in describing how to obtain the iterates of the free product $\mathcal{G}_{1} * \mathcal{G}_{2}$ in an inductive manner by appropriate glueing. Thus, in the first step we obtain $\mathcal{G}_{1} *^{(1)} \mathcal{G}_{2}$ by glueing one copy of $\mathcal{G}_{1}$ to one copy of $\mathcal{G}_{2}$ by means of identifying $e_{1}$ with $e_{2}$ and choosing it to be the root $e$. This gives $\left(\mathcal{G}_{1} \vdash \mathcal{G}_{2}\right) \star\left(\mathcal{G}_{2} \vdash\right.$ $\left.\mathcal{G}_{1}\right)$. Note that in the $m$ th step we can obtain the graph $\mathcal{G}_{1} *^{(m)} \mathcal{G}_{2}$ by glueing a copy of $\mathcal{G}_{2} *^{(m-1)} \mathcal{G}_{1}$ to every vertex of $\mathcal{G}_{1}$ but the root $e_{1}$, and vice versa, a copy of $\mathcal{G}_{1} *{ }^{(m-1)} \mathcal{G}_{2}$ to every vertex of $\mathcal{G}_{2}$ but the root $e_{2}$, and then, by glueing the two graphs obtained in that way at their roots ( $e_{1}$ and $e_{2}$, respectively). These rules of glueing correspond to the orthogonal and star products and thus the first assertion is proved. The second assertion is then a consequence of Corollary 6.1.

Below we state our results on the decompositions of the free product of uniformly locally finite rooted graphs. The limits of products of rooted graphs are understood
as inductive limits (towers of graphs with the same root). Results concerning convolutions have been proven in Ref. 20 for compactly supported probability measures (see also Ref. 17, where it is shown that $\mu \boxplus \nu=w-\lim _{m \rightarrow \infty} \mu \boxplus_{m} \nu$, with a different, purely algebraic, definition of the $m$-free convolution).

Theorem 7.2. Let $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$ be rooted graphs with spectral distributions $\mu$ and $\nu$. Then their free product admits the decomposition

$$
\mathcal{G}_{1} * \mathcal{G}_{2} \cong\left(\mathcal{G}_{1} \vdash\left(\mathcal{G}_{2} \vdash\left(\mathcal{G}_{1} \vdash \cdots\right)\right)\right) \star\left(\mathcal{G}_{2} \vdash\left(\mathcal{G}_{1} \vdash\left(\mathcal{G}_{2} \vdash \cdots\right)\right)\right)
$$

called the star-orthogonal decomposition. If $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$ are uniformly locally finite, its spectral distribution is given by $\mu \boxplus \nu=w-\lim _{m \rightarrow \infty}\left(\left(\mu \vdash_{m} \nu\right) \uplus\left(\nu \vdash_{m} \mu\right)\right)$.

Proof. The first statement follows from Theorem 6.1 and Lemma 7.1. The weak limit formula for $\mu \boxplus \nu$ is a consequence of Corollary 6.1.

Theorem 7.3. Under the assumptions of Theorem 7.2, the free product of rooted graphs admits the decomposition

$$
\mathcal{G}_{1} * \mathcal{G}_{2} \cong \mathcal{G}_{1} \triangleright\left(\mathcal{G}_{2} \vdash\left(\mathcal{G}_{1} \vdash\left(\mathcal{G}_{2} \vdash \cdots\right)\right)\right)
$$

called the comb-orthogonal decomposition. If $\mathcal{G}_{1}$ and $\mathcal{G}_{2}$ are uniformly locally finite, its spectral distribution is given by $\mu \boxplus \nu=w-\lim _{m \rightarrow \infty} \mu \triangleright\left(\nu \vdash_{m} \mu\right)$.

Proof. The first statement follows from Theorem 6.2 and Lemma 7.2. The formula for $\mu \boxplus \nu$ is a consequence of Corollary 6.1.

Corollary 7.2. Under the assumptions of Theorem 7.2, the Cauchy transform of $\mu \boxplus \nu$ can be expressed as

$$
G_{\mu \boxplus \nu}(z)=G_{\mu}\left(z-K_{\nu}\left(z-K_{\mu}\left(z-K_{\nu}(\cdots)\right)\right)\right),
$$

where the right-hand side is understood as the uniform limit on compact subsets of the complex upper half-plane.

Proof. This is a consequence of repeated application of (3.9) and Theorem 7.3.
If $\mathcal{G}_{1}$ or $\mathcal{G}_{2}$ is not uniformly locally finite, the statement of Theorem 7.3 concerning spectral distributions holds in the weaker sense of convergence of moments.

Example 7.1. Take two graphs of type given in Example 6.2. In contrast to the s-free product considered there, one cannot obtain an explicit formula for the continued Jacobi fraction corresponding to the free product (essentially, due to the presence of $G_{\mu}$ in the beginning of the above formula). Instead, we immediately obtain the algebraic formula

$$
G_{\mu \boxplus \nu}(z)=\frac{1}{z-\alpha_{0}-\frac{\omega_{0}}{z-\alpha_{1}-K_{\nu \boxplus \mu}(z)}}
$$

which allows us to find an analytic form of $G_{\mu \boxplus \nu}(z)$ once we have an analytic formula for $K_{\nu \boxplus \mu}(z)$. This type of algebraic computation was used, for instance, in Ref. 10, where we refer the reader for a general explicit formula for the Green function (equivalent to the Cauchy transform). Explicit computations based on this formula for specific graphs are given in Secs. 9-10.

## 8. Quantum Decomposition of Adjacency Matrices

In this section we will introduce a new type of "quantum decomposition" of the adjacency matrix $A(\mathcal{G})$ of a graph $\mathcal{G}$ in which the distance is measured with respect to a set of vertices instead of a single vertex.

The results of this section can be applied to any graph, not only free products of graphs to which this paper is devoted. In the latter case, our quantum decomposition is of different category than those studied in the previous sections, but it bears some resemblance to the comb-orthogonal decomposition since in some sense it "begins" with one copy of one of the graphs. However, as in the standard "quantum decomposition", ${ }^{11,12}$ its components, called "quantum components", use infinitely many copies of both graphs and, moreover, cannot be represented as subgraphs of the product graph. More importantly, it allows us to obtain more complete information on the spectral properties of many graphs, including certain free products, which cannot be obtained by means of other decompositions, including the standard "quantum decomposition".

The set of vertices with respect to which distance is measured will be denoted by $\mathcal{V}_{0}$. The sequence of sets

$$
\mathcal{V}_{n}=\left\{v \in V ; d\left(v, \mathcal{V}_{0}\right)=n\right\},
$$

where $d\left(v, \mathcal{V}_{0}\right)=\min \left\{d\left(v, v_{0}\right) ; v_{0} \in \mathcal{V}_{0}\right\}$ and $n \in \mathbb{N}^{*}:=\mathbb{N} \cup\{0\}$, will be called the distance partition of the set $V$. The associated Hilbert space decomposition is of the form

$$
\begin{equation*}
l^{2}(V)=\bigoplus_{n \in \mathbb{N}^{*}} l^{2}\left(\mathcal{V}_{n}\right) \tag{8.1}
\end{equation*}
$$

which, in turn, leads to the quantum decomposition of $A$, by which we shall understand the triple $\left(A^{+}, A^{0}, A^{-}\right)$of operators on $l^{2}(V)$ given by

$$
\begin{equation*}
A^{+} \delta(x)=\sum_{\substack{y \sim x \\ y \in \mathcal{V}_{n+1}}} \delta(y), \quad A^{0} \delta(x)=\sum_{\substack{y \sim x \\ y \in \mathcal{V}_{n}}} \delta(y), \quad A^{-} \delta(x)=\sum_{\substack{y \neq x \\ y \in \mathcal{\nu}_{n-1}}} \delta(y) \tag{8.2}
\end{equation*}
$$

whenever $x \in \mathcal{V}_{n}$. Clearly, we have $A=A^{+}+A^{0}+A^{-}$, which justifies the above terminology, and, moreover, $\left(A^{+}\right)^{*}=A^{-}$and $\left(A^{0}\right)^{*}=A^{0}$. Finally, a nonzero vector $\xi \in l^{2}(V)$ will be called a vacuum vector if $A^{-} \xi=0$. Of interest to us will be vacuum vectors of special type.

Definition 8.1. A vector $\xi \in l^{2}(V)$ will be called a $J$-vacuum vector with respect to the quantum decomposition $\left(A^{+}, A^{0}, A^{-}\right)$(or, simply, a J -vacuum vector) if it is
a vacuum vector and for every $n \in \mathbb{N}^{*}$ it holds that

$$
\begin{align*}
A^{-} A^{+}\left(A^{+n} \xi\right) & =\omega_{n} A^{+n} \xi,  \tag{8.3}\\
A^{0}\left(A^{+n} \xi\right) & =\alpha_{n} A^{+n} \xi, \tag{8.4}
\end{align*}
$$

where $\alpha_{n} \in \mathbb{R}$ and $\omega_{n} \geq 0$ and where we use the convention that $A^{+m} \xi=0$ implies $\omega_{m}=\alpha_{m}=0$ and thus $\omega_{n}=\alpha_{n}=0$ for $n \geq m$.

In this fashion we can associate with each J-vacuum vector the Jacobi parameters written in the form of a pair of sequences $(\alpha, \omega)$ (called from now on $J$ sequences), where $\alpha=\left(\alpha_{n}\right)_{n \in \mathbb{N}^{*}}$ and $\omega=\left(\omega_{n}\right)_{n \in \mathbb{N}^{*}}$.

Proposition 8.1. J-vacuum vectors associated with non-identical J-sequences are orthogonal.

Proof. Denote these vectors by $\xi$ and $\xi^{\prime}$ and the associated J-sequences by $(\alpha, \omega)$ and $\left(\alpha^{\prime}, \omega^{\prime}\right)$. Assume that $\alpha_{n} \neq \alpha_{n}^{\prime}$ for some $n \in \mathbb{N}^{*}$. Without loss of generality we can assume that $\alpha_{n} \neq 0$, which implies that $\omega_{n-1} \neq 0$. Then

$$
\left\langle A^{+n} \xi, A^{+n} \xi^{\prime}\right\rangle=\frac{1}{\alpha_{n}}\left\langle A^{0} A^{+n} \xi, A^{+n} \xi^{\prime}\right\rangle=\frac{\alpha_{n}^{\prime}}{\alpha_{n}}\left\langle A^{+n} \xi, A^{+n} \xi^{\prime}\right\rangle
$$

and therefore, $A^{+n} \xi \perp A^{+n} \xi^{\prime}$, which gives

$$
0=\left\langle A^{+n} \xi, A^{+n} \xi^{\prime}\right\rangle=\left\langle A^{-n} A^{+n} \xi, \xi^{\prime}\right\rangle=\omega\left\langle\xi, \xi^{\prime}\right\rangle,
$$

where $\omega=\omega_{0} \omega_{1} \cdots \omega_{n-1} \neq 0$ and thus $\xi \perp \xi^{\prime}$. Similar computations for the case when $\omega_{n} \neq \omega_{n}^{\prime}$ for given $n \in \mathbb{N}$ lead to orthogonality $\xi \perp \xi^{\prime}$ as well.

For a given distance partition of $V$, any set $\Xi$ of vectors from $l^{2}(V)$ will be called distance-adapted if $\Xi=\bigcup_{n \in \mathbb{N}^{*}} \Xi_{n}$, where $\Xi_{n} \subseteq l^{2}\left(\mathcal{V}_{n}\right)$. Such sets are convenient to deal with and for that reason we show in the proposition given below that if we have a set of mutually orthogonal J-vacuum vectors, which we call an orthogonal J-vacuum set, we can always choose one which is distance-adapted. For a given quantum decomposition of $A$, we denote by $[\mathcal{A} x]$ and $\left[\mathcal{A}^{+} x\right]$ the closed linear subspaces generated by vectors $\left\{A^{n} x\right\}_{n=0}^{\infty}$ and $\left\{A^{+n} x\right\}_{n=0}^{\infty}$, respectively.

Proposition 8.2. For a given distance partition of $V$ and the associated quantum decomposition of $A$, let $\Xi$ be an orthogonal J-vacuum set. Then there exists an orthogonal J-vacuum set $\Theta$, which is distance-adapted and such that

$$
\bigoplus_{\xi \in \Xi}[\mathcal{A} \xi] \subseteq \bigoplus_{\xi \in \Theta}[\mathcal{A} \xi]
$$

Proof. Let $\Xi=\left\{\xi_{i} ; i \in I\right\}$, for a countable set of indices $I$. According to the decomposition (8.1), we have

$$
\xi_{i}=\sum_{n=0}^{\infty} \xi_{i}^{(n)}
$$

for every $i \in I$, where $\xi_{i}^{(n)} \in l^{2}\left(\mathcal{V}_{n}\right)$ and $n \in \mathbb{N}^{*}$. It is not difficult to observe that each $\xi_{i}^{(n)}$ is a J-vacuum vector or $\xi_{i}^{(n)}=0$. For every $n \in \mathbb{N}^{*}$, we choose from the set $\left\{\xi_{i}^{(n)} ; i \in I\right\}$ a maximal linearly independent set, which we denote $\Gamma_{n}=\left\{\gamma_{1}, \ldots, \gamma_{k_{n}}\right\}$. If $\xi_{i}^{(n)}=0$, we set $\Gamma_{n}=\emptyset$. Of course, $k_{n} \leq\left|\mathcal{V}_{n}\right|$. We divide $\Gamma_{n}$ into disjoint classes

$$
\Gamma_{n}=\Gamma_{n}(1) \cup \Gamma_{n}(2) \cup \cdots \cup \Gamma_{n}\left(l_{n}\right)
$$

subject to the condition: $\gamma_{i}, \gamma_{j} \in \Gamma_{n}(l) \Longleftrightarrow \gamma_{i}, \gamma_{j}$ are associated with the same J-sequences. From Proposition 8.1 it follows that vectors from different classes are orthogonal. Let $\Theta_{n}(l)$ be the set obtained by applying the Gram-Schmidt orthogonalization to class $\Gamma_{n}(l)$. Then

$$
\Theta_{n}=\Theta_{n}(1) \cup \Theta_{n}(2) \cup \cdots \cup \Theta_{n}\left(l_{n}\right)
$$

is an orthogonal set. Moreover, each element of $\Theta_{n}$ is a J-vacuum vector since it is a linear combination of J -vacuum vectors associated with the same J -sequence. Finally, we take $\Theta$ to be the union of the $\Theta_{n}$, i.e. $\Theta=\bigcup_{n \in \mathbb{N}^{*}} \Theta_{n}$. From the construction of the set $\Theta$ it follows easily that it satisfies the conditions stated above.

Definition 8.2. For a given distance-adapted J-vacuum set $\Xi$ with decomposition $\Xi=\bigcup_{n \in \mathbb{N}^{*}} \Xi_{n}$, define a sequence of mutually orthogonal sets by the recurrence

$$
\begin{equation*}
B_{0}=\Xi_{0}, \quad B_{n+1}=\left(A^{+} B_{n} \cup \Xi_{n+1}\right) \backslash\{0\} \tag{8.5}
\end{equation*}
$$

The set $\Xi$ will be called generating if for every $n \in \mathbb{N}^{*}$ the set $B_{n}$ is a basis in $l^{2}\left(\mathcal{V}_{n}\right)$. This notion will turn out useful in the theorem given below.

Theorem 8.1. If $\Xi$ is an orthogonal J-vacuum set which is generating and distance-adapted, then we have the direct sum decomposition $l^{2}(V)=\bigoplus_{\xi \in \Xi}[\mathcal{A} \xi]$.

Proof. First, let us show that for $\xi \in \Xi$ it holds that $[\mathcal{A} \xi]=\left[\mathcal{A}^{+} \xi\right]$. Notice that for $m<n$ we have

$$
\begin{aligned}
\left\langle A^{+n} \xi, A^{+m} \xi\right\rangle & =\left\langle A^{+(n-m-1)} \xi, A^{-(m+1)} A^{+m} \xi\right\rangle \\
& =\omega\left\langle A^{+(n-m-1)} \xi, A^{-} \xi\right\rangle=0,
\end{aligned}
$$

and therefore $\left\{A^{+n} \xi\right\}_{n=0}^{\infty}$ is an orthogonal set. Applying the Gram-Schmidt orthogonalization to the set $\left\{A^{n} \xi\right\}_{n=0}^{\infty}$, we obtain $\left\{A^{+n} \xi\right\}_{n=0}^{\infty}$, and thus $[\mathcal{A} \xi]=\left[\mathcal{A}^{+} \xi\right]$. Next, observe that for different $\xi, \xi^{\prime} \in \Xi$ and arbitrary $m, n \geq 0$, the vectors $A^{+n} \xi$, $A^{+m} \xi^{\prime}$ are orthogonal, which follows from a straightforward induction. Therefore, $\left[\mathcal{A}^{+} \xi\right] \perp\left[\mathcal{A}^{+} \xi^{\prime}\right]$. Finally, we know by assumption that $B_{n}$ is a basis in $l^{2}\left(\mathcal{V}_{n}\right)$, and therefore

$$
l^{2}(V)=\bigoplus_{n=0}^{\infty} l^{2}\left(\mathcal{V}_{n}\right)=\bigoplus_{n=0}^{\infty} \operatorname{span}\left(B_{n}\right) \subseteq \bigoplus_{\xi \in \Xi}\left[\mathcal{A}^{+} \xi\right]=\bigoplus_{\xi \in \Xi}[\mathcal{A} \xi]
$$

which completes the proof since the reverse implication is obvious.

Let us observe that Theorem 8.1 gives an interacting Fock space decomposition of $l^{2}(V)$ since $[\mathcal{A} \xi]$ is an interacting Fock space for each $\xi \in \Xi$, in which the set $\left\{A^{+n} \xi\right\}_{n=0}^{\infty}$ is a basis. The results of this Section give us sufficient conditions for an orthogonal decomposition of $l^{2}(V)$ of Theorem 8.1 to exist and that in turn allows us to get detailed information about the spectral properties of the adjacency matrix $A$, including spectral distributions associated with all vacuum vectors $\xi \in \Xi$ which appear in this decomposition. In particular, this also gives the spectrum of the considered graph.

## 9. Trees

In this section we use the theory of Sec. 8 to consider the simplest examples of $n$-ary trees and homogenous trees. Here, the set $\mathcal{V}_{0}$ will consist of one root, which considerably simplifies the spectral analysis and the corresponding quantum decomposition agrees with that used in the approach of Hora and Obata. ${ }^{11,12}$ However, our analysis goes a little further since we study spectral distributions associated with all cyclic vectors.

## 9.1. $\boldsymbol{n}$-ary trees $\mathbb{T}_{\boldsymbol{n}}$

By Theorem 8.1, it suffices to find a distance-adapted generating J-vacuum set. Denote by $W_{n}$ the set of words in $n$ letters $a_{1}, a_{2}, \ldots, a_{n}$, including the empty word. Then there exists a natural bijection between $W_{n}$ and $V\left(\mathbb{T}_{n}\right)$. If $w$ labels a vertex of $\mathbb{T}_{n}$ for which $d(w, e)=m$, then $a_{1} w, a_{2} w, \ldots, a_{n} w$ denote "sons" of $w$. Let $\Xi_{0}=\{\delta(e)\}$ and

$$
\Xi_{m}=\left\{\sum_{j=1}^{k}\left(\delta\left(a_{j} w\right)-\delta\left(a_{k+1} w\right)\right), 1 \leq k \leq n-1, w \in W_{n},|w|=m-1\right\}
$$

for every natural $m \geq 1$. Clearly, $\Xi=\bigcup_{m=0}^{\infty} \Xi_{m}$ is a distance-adapted orthogonal J-vacuum set. In fact, it is easy to see that it is an orthogonal set of vacuum vectors. Next, observe that the cardinalities of sets

$$
B_{m}=A^{+}\left(B_{m-1}\right) \cup \Xi_{m}, \quad m \geq 1
$$

with $B_{0}=\Xi_{0}$, satisfy the recurrence

$$
\left|B_{m}\right|=\left|B_{m-1}\right|+(n-1) n^{m-1}
$$

since $\left|\Xi_{m}\right|=(n-1) n^{m-1}$, which gives $\left|B_{m}\right|=n^{m}=\left|\mathcal{V}_{m}\right|$ and thus $B_{m}$ is a basis of $\mathcal{V}_{m}$. By Theorem 8.1, we have a direct sum decomposition of $l^{2}(V)$ into the sum of $[\mathcal{A} \xi], \xi \in \Xi$. Finally, every $\delta(w)$ (and thus every $\xi \in \Xi$ ) is an eigenvector of both $A^{-} A^{+}$and $A^{0}$ with eigenvalues $n$ and 0 , respectively, for every $w \in V\left(\mathbb{T}_{m}\right)$. Therefore, every $\xi$ is a J-vacuum vector and the associated J-sequences are $\omega_{k}(\xi)=$ $n$ and $\alpha_{k}(\xi)=0$ for every $k$. This shows that to every $\xi$ corresponds the same Cauchy transform, namely that of the form (5.3) and thus the measure (5.4). Thus,
the spectrum of $\mathbb{T}_{n}$ agrees with the support of that measure, which is the interval $[-2 \sqrt{n}, 2 \sqrt{n}]$.

### 9.2. Homogenous trees $\mathbb{H}_{\boldsymbol{n}}$

A similar approach can be used for homogeneous trees. Since $\mathbb{H}_{n}$ is a symmetric graph, any vertex can be chosen to be the root denoted $e$.

The case $n=1$ is straightforward, we have $\Xi=\{\delta(e)\}$ and the decomposition of Theorem 8.1 is simply $l^{2}(V)=[\mathcal{A} \delta(e)]$. Therefore, let $n \geq 2$ and denote by $a_{1}$, $a_{2}, \ldots, a_{n}$ the "sons" of $e$ and label all the vertices with distance bigger than 2 from the root using only $a_{1}, a_{2}, \ldots, a_{n-1}$. The situation is very similar to that of the $n$-ary trees. Let $\Xi_{0}=\{\delta(e)\}$ and

$$
\begin{aligned}
& \Xi_{1}=\left\{\sum_{j=1}^{k}\left(\delta\left(a_{j}\right)-\delta\left(a_{k+1}\right)\right), 1 \leq k \leq n-1\right\} \\
& \Xi_{m}=\left\{\sum_{j=1}^{k}\left(\delta\left(a_{j} w\right)-\delta\left(a_{k+1} w\right)\right), 1 \leq k \leq n-2, w \in W_{n-1},|w|=m-1\right\}
\end{aligned}
$$

(for $m=2$, the procedure of finding cyclic vectors stops at $\Xi_{1}$ ), where $W_{m-1}$ is the set of words in $a_{1}, a_{2}, \ldots, a_{n-1}$ of length equal to $m-1$. The set $\Xi=\bigcup_{m=0}^{\infty} \Xi_{m}$ is clearly a distance-adapted orthogonal set of vacuum vectors. To show that $\Xi$ is generating, take the sequence $\left(B_{m}\right)$ defined by (8.5) and observe that we have the recurrence

$$
\left|B_{m}\right|=\left|B_{m-1}\right|+n(n-1)^{m-2}(n-2)
$$

since $\left|\Xi_{m}\right|=n(n-1)^{m-2}(n-2)$, which is solved by $\left|B_{m}\right|=n(n-1)^{m-1}=\left|\mathcal{V}_{m}\right|$, and thus $B=\bigcup_{m \in \mathbb{N}^{*}} B_{m}$ is a basis in $l^{2}(V)$.

The vector $\delta(e)$ (and thus $\xi_{0}$ ) is an eigenvector of both $A^{-} A^{+}$and $A^{0}$ with eigenvalues $n$ and 0 , respectively. This gives J-sequences $\omega_{0}\left(\xi_{0}\right)=n$ and $\omega_{k}\left(\xi_{0}\right)=$ $n-1$ for $k \geq 1$ with $\alpha_{k}(\xi)=0$ for all $k$, and thus the Cauchy transform

$$
G_{\xi_{0}}(z)=\frac{-2 z+n z-n \sqrt{4-4 n+z^{2}}}{2(n-z)(n+z)}
$$

If $n=1$, the measure $\mu_{\xi_{0}}$ consists of two atoms at $\pm 1$ of mass $1 / 2$ each. For $n \geq 2$, $\mu_{\xi_{0}}$ is absolutely continuous with respect to Lebesgue measure and has density

$$
f_{\xi_{0}}(x)=\left|\frac{n \sqrt{4 n-4-x^{2}}}{2 \pi(n-x)(n+x)}\right|
$$

on the interval $[-2 \sqrt{n-1}, 2 \sqrt{n-1}]$.
In turn, vectors $\delta(w) \neq \delta(e)$ (and thus $\xi \neq \xi_{0}$ ) are eigenvectors of both $A^{-} A^{+}$ and $A^{0}$ with eigenvalues $n-1$ and 0 , respectively. Therefore, $\xi \neq \xi_{0}$ (for $n \geq 2$ )
are J-vacuum vectors with the Jacobi parameters of the form $\omega_{k}(\xi)=n-1$ and $\alpha_{k}(\xi)=0$ for all $k$, which give the Cauchy transform

$$
G_{\xi}(z)=\frac{z-\sqrt{z^{2}-4(n-1)}}{2(n-1)}
$$

and thus $\mu_{\xi}$ is the Wigner measure on the interval $[-2 \sqrt{n-1}, 2 \sqrt{n-1}]$. Although we get two different spectral distributions, their contributions to the spectrum $\operatorname{spec}\left(\mathbb{H}_{n}\right)=[-2 \sqrt{n-1}, 2 \sqrt{n-1}]$ are the same.

## 10. Other Examples

In this section we apply the same approach as in Sec. 9 to two types of examples: free products of complete graphs $\mathbb{K}_{n} * \mathbb{K}_{m}$ and free products of a complete graph with a graph of "fork"-type $\mathbb{K}_{n} * \mathbb{F}_{m}$. In the first case, we deal with free products of symmetric graphs and therefore their spectra can be determined from one spectral distribution (see Ref. 10 for a detailed study). However, since our approach examines spectral distributions associated with all cyclic vectors, even in this case we obtain new information (for instance, about the "multiplicity" of the spectrum).

### 10.1. Free products $\mathbb{K}_{\boldsymbol{n}} * \mathbb{K}_{\boldsymbol{m}}$

By $\mathbb{K}_{n}$ we denote the complete graph with $n+1$ vertices, i.e. such in which each pair of vertices forms an edge. Observe that any choice of a root gives an isomorphic rooted graph since all vertices are equivalent. Take two complete graphs $\mathbb{K}_{n}$ and $\mathbb{K}_{m}$ with vertices $x_{0}, x_{1}, \ldots, x_{n}$ (choose $x_{0}$ as the root) and $y_{0}, y_{1}, \ldots, y_{m}$ (choose $y_{0}$ as the root), respectively, where $n, m \geq 1$. Then $\mathbb{K}_{n} * \mathbb{K}_{m}=(V, E, e)$, where the distance partition of the set of vertices $V$ can be given by the recursion

$$
\begin{align*}
\mathcal{V}_{0} & =\left\{e, x_{1}, \ldots, x_{n}\right\}, \\
\mathcal{V}_{2 k+1} & =\left\{y_{i} w ; w \in \mathcal{V}_{2 k}, i=1, \ldots, m\right\},  \tag{10.1}\\
\mathcal{V}_{2 k+2} & =\left\{x_{i} w ; w \in \mathcal{V}_{2 k+1}, i=1, \ldots, n\right\},
\end{align*}
$$

where $k=0,1, \ldots$ (see Figs. 3 and 4). Directly from this construction it follows that

$$
\begin{equation*}
\left|\mathcal{V}_{2 k}\right|=(n+1) m^{k} n^{k} \quad \text { and } \quad\left|\mathcal{V}_{2 k+1}\right|=(n+1) m^{k+1} n^{k} . \tag{10.2}
\end{equation*}
$$

Let $\left(A^{+}, A^{-}, A^{0}\right)$ be the quantum decomposition of the incidence matrix of $\mathbb{K}_{n} * \mathbb{K}_{m}$. We shall find the corresponding orthogonal J-vacuum set $\Xi$ which is generating and distance-adapted. Let

$$
\Xi_{0}=\left\{\sum_{j=0}^{n} \delta\left(x_{j}\right)\right\} \cup\left\{\sum_{j=0}^{i-1}\left(\delta\left(x_{j}\right)-\delta\left(x_{i}\right)\right) ; i=1, \ldots, n\right\}
$$



Fig. 3. $\quad \mathbb{K}_{2} * \mathbb{K}_{2}$.


Fig. 4. $\quad \mathbb{K}_{1} * \mathbb{K}_{2}$.

$$
\begin{align*}
& \Xi_{2 k+1}=\left\{\sum_{j=1}^{i-1}\left(\delta\left(y_{j} w\right)-\delta\left(y_{i} w\right)\right) ; w \in \mathcal{V}_{2 k}, i=2, \ldots, m\right\},  \tag{10.3}\\
& \Xi_{2 k+2}=\left\{\sum_{j=1}^{i-1}\left(\delta\left(x_{j} w\right)-\delta\left(x_{i} w\right)\right) ; w \in \mathcal{V}_{2 k+1}, i=2, \ldots, n\right\},
\end{align*}
$$

where $x_{0}$ is identified with $e$. This gives

$$
\left|\Xi_{2 k+1}\right|=(m-1)\left|\mathcal{V}_{2 k}\right| \quad \text { and } \quad\left|\Xi_{2 k+2}\right|=(n-1)\left|\mathcal{V}_{2 k-1}\right|
$$

for $k \in \mathbb{N}^{*}$, which, by (10.2) leads to

$$
\begin{equation*}
\left|\Xi_{k}\right|=\left|\mathcal{V}_{k}\right|-\left|\mathcal{V}_{k-1}\right| \tag{10.4}
\end{equation*}
$$

An elementary computation shows that $\Xi_{k}$ is an orthogonal set for every $k \in \mathbb{N}^{*}$. Moreover, it is clear that these sets are mutually orthogonal and that they contain only vacuum vectors (use $A^{-} \delta\left(y_{j} w\right)=\delta(w)$ and $A^{-} \delta\left(x_{j} w^{\prime}\right)=\delta\left(w^{\prime}\right)$ ) and therefore $\Xi$ is an orthogonal set of vacuum vectors.

We will show now that it is an orthogonal J-vacuum set. Let $\xi \in \Xi_{2 k+1}$ and $r \in \mathbb{N}$. We have

$$
A^{+r} \xi=A^{+r} \sum_{j=1}^{i-1}\left(\delta\left(y_{j} w\right)-\delta\left(y_{i} w\right)\right)=\sum_{j=1}^{i-1} \sum_{\substack{|u|=r \\ u y_{i} w \in V}}\left(\delta\left(u y_{j} w\right)-\delta\left(u y_{i} w\right)\right)
$$

for some $w \in \mathcal{V}_{2 k}$. If $r$ is odd, then each $u$ in the above sum begins with a vertex from $\mathbb{K}_{n}$ and therefore each $\delta\left(u y_{k} w\right)$ in the above sum (and thus also $A^{+r} \xi$ ) is an eigenvector of $A^{-} A^{+}$with eigenvalue $m$. Moreover, the sums $\sum_{\substack{|u|=r \\ u y_{k} w \in V}} \delta\left(u y_{k} w\right)$
(and thus also $A^{+r} \xi$ ) are eigenvectors of $A^{0}$ with eigenvalue $n-1$ for each $k$. Therefore, $\omega_{r}(\xi)=m$ and $\alpha_{r}(\xi)=n-1$ for $r$ odd. In the case of $r$ even, each $u$ begins with a vertex from $\mathbb{K}_{m}$ in the above reasoning and thus $\omega_{r}(\xi)=n$ and $\alpha_{r}(\xi)=m-1$. Analogous computations holds when $\xi \in \Xi_{2 k}$. Therefore, all vectors from $\Xi$ are J-vacuum vectors.

To show that $\Xi$ is generating, we need to check that for every $k \in \mathbb{N}^{*}$ the set $B_{k}$, defined recursively by (8.5), is a basis in $l^{2}\left(\mathcal{V}_{k}\right)$. For that purpose it is enough to show that $\left|B_{k}\right|=\left|\mathcal{V}_{k}\right|$ since $B_{k}$ is a set of non-zero orthogonal vectors. This is shown by induction. For $k=0$ we have $\left|B_{0}\right|=\left|\Xi_{0}\right|=n+1=\left|\mathcal{V}_{0}\right|$. Assume now that $\left|B_{k}\right|=\left|\mathcal{V}_{k}\right|$ for some $k \in \mathbb{N}$. Then from (10.4) we get

$$
\left|B_{k+1}\right|=\left|A^{+} B_{k}\right|+\left|\Xi_{k+1}\right|=\left|\mathcal{V}_{k}\right|+\left|\mathcal{V}_{k+1}\right|-\left|\mathcal{V}_{k}\right|=\left|\mathcal{V}_{k+1}\right|
$$

and our claim holds by induction. This proves that $\Xi$ is a generating J-vacuum set.
In order to find the spectrum of $A$, we find the probability measures $\mu_{\xi}$ with moments $\mu_{\xi}(n)=\left\langle A^{n} \xi, \xi\right\rangle$ for every $\xi \in \Xi$ and then take the union of their supports. Let

$$
\xi_{0}=\delta(e)+\delta\left(x_{1}\right)+\cdots+\delta\left(x_{n}\right)
$$

The associated J-sequences are of the form

$$
\omega_{k}\left(\xi_{0}\right)=\left\{\begin{array}{ll}
n & k \text { odd } \\
m & k \text { even }
\end{array}, \quad \alpha_{k}\left(\xi_{0}\right)= \begin{cases}n & k=0 \\
m-1 & k \text { odd } \\
n-1 & k \text { even positive }\end{cases}\right.
$$

and they give a continued-fraction representation of the Cauchy transform $G_{\xi_{0}}$ of the measure $\mu_{\xi_{0}}$, which leads to

$$
\begin{gathered}
G_{\xi_{0}}(z)=\frac{1-m n+m z+n z-z^{2}+\sqrt{w(z)}}{2(-1+n-z)(m+n-z)} \\
w(z)=\left(1-2 n+m n+(2-m-n) z+z^{2}\right)^{2}-4 m(1-m+z)(1-n+z) .
\end{gathered}
$$

Applying the Stieltjes inversion formula to the transform $G_{\xi_{0}}$, we obtain the absolutely continuous part of $\mu_{\xi_{0}}$ of the form

$$
f_{\xi_{0}}(x)=\left|\frac{\sqrt{-w(x)}}{2 \pi(x+1-n)(x-m-n)}\right|, \quad x \in I_{m, n}
$$

on the set $I_{m, n}$ being the union of two closed intervals with ends at $\frac{1}{2}(m+n-2 \pm$ $\left.\sqrt{4(\sqrt{m} \pm \sqrt{n})^{2}+(m-n)^{2}}\right)$ and disjoint interiors. In addition, $\mu_{\xi_{0}}$ has an atom at $n-1$ of mass $\frac{1}{1+m} \max \{0, m-n\}$.

Now, let $\xi \in \Xi_{2 k}$ and $\xi \neq \xi_{0}$. Then the J-sequences are of the form

$$
\omega_{k}(\xi)=\left\{\begin{array}{ll}
n & k \text { odd } \\
m & k \text { even }
\end{array}, \quad \alpha_{k}(\xi)= \begin{cases}-1 & k=0 \\
m-1 & k \text { odd } \\
n-1 & k \text { even positive }\end{cases}\right.
$$

which give the Cauchy transform

$$
G_{\xi}(z)=\frac{1-2 m+2 n-m n+(2-m+n) z+z^{2}+\sqrt{w(z)}}{2 n(1-m+z)(2+z)}
$$

Again, the Stieltjes inversion formula gives the absolutely continuous part of $\mu_{\xi}$ of the form

$$
f_{\xi}(x)=\left|\frac{\sqrt{-w(x)}}{2 n(1-m+x)(2+x)}\right|, \quad x \in I_{m, n}
$$

Besides, $\mu_{\xi}$ has atoms at -2 and $m-1$ of masses $\frac{m n-1}{n(1+m)}$ and $\frac{1}{n(1+m)} \max \{0, n-m\}$, respectively.

For $\xi \in \Xi_{2 k+1}$, the J-sequences are of the form

$$
\omega_{k}(\xi)=\left\{\begin{array}{ll}
m & k \text { odd } \\
n & k \text { even }
\end{array}, \quad \alpha_{k}(\xi)= \begin{cases}-1 & k=0 \\
n-1 & k \text { odd } \\
m-1 & k \text { even positive }\end{cases}\right.
$$

which give the Cauchy transform

$$
\begin{gathered}
G_{\xi}(z)=\frac{1-2 n+2 m-m n+(2-n+m) z+z^{2}+\sqrt{w(z)}}{2 m(1-n+z)(2+z)} . \\
\quad f_{\xi}(x)=\left|\frac{\sqrt{-w(x)}}{2 m(1-n+x)(2+x)}\right|, \quad x \in I_{m, n} .
\end{gathered}
$$

Besides, $\mu_{\xi}$ has atoms at -2 and $n-1$ of masses $\frac{m n-1}{m(1+n)}$ and $\frac{1}{m(1+n)} \max \{0, m-n\}$, respectively.

We conclude that for any $m, n \geq 1$ the continuous part of the spectrum of $\mathbb{K}_{n} * \mathbb{K}_{m}$ agrees with $I_{m, n}$. The point spectrum is given by

$$
\begin{cases}\emptyset & \text { if } m=n=1 \\ \{-2\} & \text { if } m=n>1 \\ \{-2, m-1\} & \text { if } m<n \\ \{-2, n-1\} & \text { if } m>n\end{cases}
$$

### 10.2. Free products $\mathbb{K}_{\boldsymbol{n}} * \mathbb{F}_{\boldsymbol{m}}$

By $\mathbb{F}_{m}$ we denote the fork of degree $m$, i.e. a connected simple graph with $m+1$ vertices, say $y_{0}, y_{1}, \ldots, y_{m}$, in which the only edges are given by $y_{i} \sim y_{0}$ for $i \neq 0$. As before, by $\mathbb{K}_{n}$ we denote a complete graph with vertices $x_{0}, x_{1}, \ldots, x_{n}$. It is easy to see that the set of vertices of $\mathbb{K}_{n} * \mathbb{F}_{m}$ coincides with that in the previous example. Besides, on $V$ we introduce the same distance partition as in (8.1) (see Fig. 5). Let $\Xi=\sum_{i=0}^{\infty} \Xi_{i}$ be the set of vectors defined as in (10.3). That $\Xi$ is a generating J-vacuum set is shown as in the case of $\mathbb{K}_{n} * \mathbb{K}_{m}$.


Fig. 5. $\mathbb{K}_{2} * \mathbb{F}_{2}$.

Let us compute now the measures $\mu_{\xi}$ with moments $\mu_{\xi}(k)=\left\langle A^{k} \xi, \xi\right\rangle$, for all $\xi \in \Xi$. The Jacobi parameters associated with vector $\xi_{0}$ are of the form

$$
\omega_{k}\left(\xi_{0}\right)=\left\{\begin{array}{ll}
n & k \text { odd } \\
m & k \text { even }
\end{array}, \quad \alpha_{k}\left(\xi_{0}\right)= \begin{cases}n & k=0 \\
0 & k \text { odd } \\
n-1 & k \text { even positive }\end{cases}\right.
$$

The Cauchy transform $G_{\xi_{0}}$ of measure $\mu_{\xi_{0}}$ has the form

$$
G_{\xi_{0}}(z)=\frac{n-m-z-n z+z^{2}-\sqrt{v(z)}}{2\left(m-n^{2}+2 n z-z^{2}\right)}
$$

where $v(z)=\left(m-n+z-n z+z^{2}\right)^{2}-4 m z(z+1-n)$. The measure $\mu_{\xi_{0}}$ is absolutely continuous with respect to Lebesgue measure and its density is of the form

$$
f_{\xi_{0}}(x)=\left|\frac{\sqrt{-v(x)}}{2 \pi\left(m-n^{2}+2 n x-x^{2}\right)}\right|, \quad x \in J_{m, n}
$$

where $J_{m, n}$ denotes the set, which is a union of two closed intervals with disjoint interiors and ends at $\frac{1}{2}\left(n-1 \pm \sqrt{4(\sqrt{m} \pm \sqrt{n})^{2}+(n-1)^{2}}\right)$.

Let now $\xi \in \Xi_{2 k}$ and $\xi \neq \xi_{0}$. Then the Jacobi parameters are of the form

$$
\omega_{k}(\xi)=\left\{\begin{array}{ll}
n & k \text { odd } \\
m & k \text { even }
\end{array}, \quad \alpha_{k}(\xi)= \begin{cases}-1 & k=0 \\
0 & k \text { odd } \\
n-1 & k \text { even positive }\end{cases}\right.
$$

which gives the Cauchy transform

$$
G_{\xi}(z)=\frac{z^{2}+z+n z-m+n-\sqrt{v(z)}}{2 n\left(1-m+2 z+z^{2}\right)}
$$

The absolutely continuous part of the measure $\mu_{\xi}$ is given by

$$
f_{\xi}(x)=\left|\frac{\sqrt{-v(x)}}{2 \pi n\left(1-m+2 x+x^{2}\right)}\right|, \quad x \in J_{m, n}
$$

Besides, $\mu_{\xi}$ has two atoms at $\pm \sqrt{m}-1$ of mass $\frac{n-1}{2 n}$ each.

For $\xi \in \Xi_{2 k+1}$, the Jacobi parameters are of the form

$$
\omega_{k}(\xi)=\left\{\begin{array}{ll}
m & k \text { odd } \\
n & k \text { even }
\end{array}, \quad \alpha_{k}(\xi)= \begin{cases}n-1 & k \text { odd } \\
0 & k \text { even }\end{cases}\right.
$$

which gives the Cauchy transform

$$
G_{\xi}(z)=\frac{m-n+z-n z+z^{2}-\sqrt{v(z)}}{2 m z}
$$

The absolutely continuous part of $\mu_{\xi}$ is given by

$$
f_{\xi}(x)=\left|\frac{\sqrt{-v(x)}}{2 \pi m z}\right|, \quad x \in J_{m, n}
$$

Besides, $\mu_{\xi}$ has an atom at 0 of mass $\frac{1}{m} \max \{0, m-n\}$.
We conclude that the continuous spectrum of $\mathbb{K}_{n} * \mathbb{F}_{m}$ coincides with $J_{m, n}$. Besides, $A$ has a discrete spectrum of the form

$$
\begin{cases}\emptyset & \text { if } m=n=1 \\ \{0\} & \text { if } m>n=1 \\ \{-2,0\} & \text { if } n>m=1 \\ \{\sqrt{m}-1,-\sqrt{m}-1\} & \text { if } n \geq m>1 \\ \{\sqrt{m}-1,-\sqrt{m}-1,0\} & \text { if } m>n>1\end{cases}
$$

For a general study of measures associated with mixed periodic Jacobi continued fractions, we refer the reader to Ref. 13.

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