

Author's Accepted Manuscript

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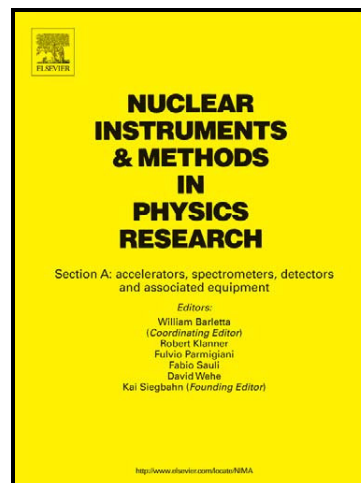
PII: S0168-9002(09)01225-X
DOI: doi:10.1016/j.nima.2009.06.024
Reference: NIMA 50221

To appear in: *Nuclear Instruments and Methods in Physics Research A*

Received date: 9 June 2009
Accepted date: 11 June 2009

Cite this article as: A. Pietropaolo, E. Perelli Cippo, G. Gorini, M. Tardocchi, E.M. Schooneveld, C. Andreani and R. Senesi, γ -ray background sources in the VESUVIO spectrometer at ISIS spallation neutron source, *Nuclear Instruments and Methods in Physics Research A*, doi:10.1016/j.nima.2009.06.024

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4 A. Pietropaolo ^{a,*}, E. Perelli Cippo ^b, G. Gorini ^a,
5 M. Tardocchi ^b, E. M. Schooneveld ^c C. Andreani ^d R. Senesi ^d

6 ^a*CNISM Milano-Bicocca, Università degli Studi di Milano-Bicocca Dipartimento di*
7 *Fisica "G. Occhialini", Piazza della Scienza 3, 20126 Milano, Italy and NAST*
8 *Center (Nanoscienze-Nanotecnologie-Strumentazione) Università degli Studi di*
9 *Roma Tor Vergata, via della Ricerca Scientifica 1, 00133 Roma, Italy*

10 ^b*Università degli Studi di Milano-Bicocca, Dipartimento di Fisica "G. Occhialini",*
11 *Piazza della Scienza 3, 20126 Milano, Italy*

12 ^c*ISIS Facility, Rutherford Appleton Laboratory, Chilton, Didcot, Oxfordshire OX11*
13 *OX11, United Kingdom*

14 ^d*Università degli Studi di Roma Tor Vergata, Dipartimento di Fisica and NAST*
15 *Center (Nanoscienze-Nanotecnologie-Strumentazione), via della Ricerca*
16 *Scientifica 1, 00133 Roma, Italy*

17 **Abstract**

18 An investigation of the gamma background was carried out in the VESUVIO spec-
19 trometer at the ISIS spallation neutron source. This study, performed with a Yttrium-
20 Aluminum-Perovskite (YAP) scintillator, follows high resolution pulse height mea-
21 surements of the gamma background carried out on the same instrument with the
22 use of a High-purity Germanium detector. In this experimental work, a mapping
23 of the gamma background was attempted, trying to find the spatial distribution
24 and degree of directionality of the different contributions identified in the previous
25 study. It is found that the gamma background at low times is highly directional
26 and mostly due to the gamma rays generated in the moderator-decoupler system.
27 The other contributions, consistently to the findings of a previous experiment, are
28 identified as a nearly isotropic one due to neutron absorption in the walls of the
29 experimental hall, and a directional one coming from the beam dump.

30 **PACS:**29.40.Mc, 98.70.Vc, 29.30.Hs

31 *Key words:* Neutron Instrumentation, gamma background, spallation neutron
32 source

33 1 Introduction

34 Spallation neutron sources [1] represent a great opportunity for both funda-
35 mental and applied science. The wide spectrum of the neutron beams pro-
36 duced from the interaction of high energy protons (in the GeV region) with
37 heavy metal targets can be used in many applications. At the ISIS spallation
38 source [2] these studies are accomplished using neutrons from meV to several
39 tens of eV; the experimental techniques used range from neutron diffraction
40 [3–7] to inelastic scattering [8–12], the latter being used e.g. to investigate
41 the dynamics of hydrogenated systems. High energy neutrons (in the MeV
42 region) are also used e.g. to simulate the spectrum of the atmospheric neu-
43 trons produced by the high energy primary cosmic rays. These neutrons are a
44 growing concern for the reliability of electronic devices [13,14]. Moreover, neu-
45 trons in the eV-keV region can be used to develop innovative neutron-based
46 imaging and tomographic techniques for cultural heritage applications [15].
47 One important aspect to be considered in all applications of pulsed neutron
48 beams is the clear identification of the background sources. For example, the
49 neutron detection techniques at epithermal energies on VESUVIO beamline
50 [16] at ISIS are based on (n,γ) conversion, so that the γ -ray background needs
51 to be investigated. In a previous experimental paper [17], a first investigation
52 of the γ background in the VESUVIO spectrometer was performed by record-
53 ing pulse height spectra with a High-Purity Germanium (HPGe) detector.
54 These measurements allowed a clear identification of several γ -ray lines and
55 the recognition of the possible γ -ray background sources.
56 In the present study, an Yttrium-Aluminum-Perovskite (YAP) scintillation
57 detector was used since a mapping of the γ -ray background was carried out
58 by assessing the contribution of different γ -ray to the time of flight spectra
59 recorded at different locations. Previous studies demonstrated that YAP is
60 well suited for γ -ray measurements, as it is very insensitive to neutrons [18].

61 2 Experimental set up

62 The measurements were performed at the VESUVIO beam line operating at
63 the ISIS spallation neutron source. The layout of the instrument is shown in
64 figure 1. The water moderator at 295 K is placed at about 11 m from the
65 sample position and operated in the so-called wing configuration [19] above
66 the spallation target. A gadolinium poisoning system is used to lower the
67 Maxwellian component, together with a decoupler to reduce the intensity
68 of the over-thermalized neutrons in the reflectors system. The beam dump,
69 placed at about 5 m beyond the sample position, is mostly composed of hy-

* corresponding author: antonino.pietropaolo@mib.infn.it, fax. +390264482367

70 drogen, iron and boron.

71 The YAP scintillation detector used for the measurements was initially placed
72 on the top of the sample tank (about 30 cm from the beam axis) and sur-
73 rounded by a lead shielding (10 cm thick) covering the whole solid angle. In
74 order to study the background along different directions, small apertures were
75 opened depending on the chosen direction to investigate. The signals from the
76 detector were sent to the data acquisition electronics to record time of flight
77 (tof) spectra. In this configuration, an investigation of the directionality of
78 the γ -ray background was done, as discussed in the next section. Another set
79 of measurements was performed by placing the detector, without shielding, at
80 different positions within the experimental hall along the z direction shown in
81 figure 1. In this way it was possible to record tof spectra at different places,
82 thus providing a mapping of the γ -ray background.

83 3 Results and discussion

84 Figure 2 shows three normalized tof spectra recorded by the YAP detector
85 when placed on the sample tank and surrounded by the lead shielding. For
86 these measurements no scattering sample was used. It can be noticed that the
87 low tof tails are different in the three cases: in the case of shielding with an
88 acceptance window towards the moderator, the rate is higher by a factor be-
89 tween 2 and 8 in the first 30-40 μ s, as compared to the other cases (see figure
90 3). The other two spectra instead (with the lateral and rear aperture in the
91 shielding) are similar, their ratio being about 1 over the whole time region. In
92 the three spectra of figure 2, the peaks riding on top of the continuum are due
93 to the radiative neutron capture in the shielding. Indeed, this was not made
94 of pure lead and contained impurities such as antimony which has several
95 neutron resonances in the epithermal region. The main contribution to the
96 continuum in the low tof region comes from the moderator-decoupler system.
97 Indeed, the γ -rays production mechanisms in the moderator and in the de-
98 coupling system are known to decay exponentially with different "relaxation
99 times" whose magnitudes are in the order of few tens and few hundreds of mi-
100 croseconds, respectively [17]. Thus, the difference in the observed count rate
101 may be possibly attributed to this highly directional component. Although the
102 background from the beam dump is presumably directional, no difference be-
103 tween the background in the cases with lateral and rear apertures is observed.
104 This is due to the large distance of the scintillator from the beam dump (about
105 5 meters), so that the isotropic contribution from the lateral walls dominates
106 over other contributions.
107 It is worth reminding here that the count rate due to the different background

108 components was written as [17]:

$$109 \quad B(t) = B_{iso}(t) + B_{dir}(t) + B'_z(t) + B''_z(t) + B_f(t) \quad (1)$$

110 where $B_{iso}(t)$ represents the isotropic contribution (mostly) due to the walls
 111 surrounding the spectrometer, $B_{dir}(t)$ is the term coming from the beam dump,
 112 $B'_z(t)$ and $B''_z(t)$ are due to a sort of γ -ray halo produced by the moderator-
 113 reflector-decoupler system, while $B_f(t)$ (the "gamma flash" component) rep-
 114 resents a transient process at very short times. For a thorough explanation of
 115 the different terms the reader is referred to ref. [17].

116 For the other set of measurements, *i.e.* those performed at different positions
 117 within the experimental hall, we calculated the count rates C_α and C_β in two
 118 different tof intervals α and β , delimited by dashed lines in the upper panel
 119 of figure 4.

120 The spectrum shown in panel (a) of the figure is relative to a measurement
 121 at intermediate z (*i.e.* near the sample tank), while that shown in panel (b)
 122 is relative to a measurement near the beam dump. In the former, no evident
 123 features are present, while in the latter, resonance peaks can be well identified
 124 against the continuum for tof values below $200 \mu\text{s}$.

125 C_α and C_β show opposite trends with varying z , as shown in figure 5. In par-
 126 ticular, C_α is decreasing with z , ranging from 1.6 kHz at 50 cm to about 0.9
 127 kHz at 380 cm (*i.e.* closer to the beam dump position), while C_β varies from
 128 about 0.9 kHz at 50 cm to about 1.9 kHz at 380 cm. For $z = 380$ cm, C_α is
 129 found to be close to 1.9 kHz if the total counts are considered. This value is
 130 obtained if the structures (resonance peaks) are considered, while 0.9 kHz is
 131 obtained if one performs the counts sum over values obtained by interpolating
 132 the continuum beneath the peaks. Despite the collimation system is made of
 133 a material similar to the beam dump, the peaks structure is covered by the
 134 higher background level present at low tof.

135 As far as the C_α parameter is concerned, as a first guess one should expect a
 136 $1/(z+L)^2$ behavior as $z+L$ ($L \simeq 900$ cm) is the distance from the correspond-
 137 ing background source (the moderator-reflector system). The continuous line
 138 if figure 5 represent the $1/(z+L)^2$ function normalized to the count rate at z
 139 $= 225$ cm.

140 The count rate C_α , being calculated in a time region around $50 \mu\text{s}$ is domi-
 141 nated by the contributions of the γ -rays coming from moderator and reflector
 142 [17]. The differences between the experimental values of C_α and the continu-
 143 ous line are due to the contribution of the γ -rays produced in the collimation
 144 system (see figure 1) at low z and to the increasing contribution of the di-
 145 rectional γ -rays from the beam dump at higher z . For C_β , that is calculated
 146 in the time region around $350 \mu\text{s}$, the contribution of moderator and reflector
 147 is almost negligible [17]. A slow linear increase in the count rate between $z=$
 148 50 and 300 cm is obtained, followed by a rapid increase in the proximity of
 149 the beam dump. This is compatible with a picture where the contribution

150 of the isotropic component due to the surrounding walls around the detector
151 overwhelms the count rate (and it is almost constant along z) superimposed
152 to a varying count rate due to the decreasing distance between beam dump
153 and detector.

154 4 Conclusions

155 It is found that the γ -ray background at short tof values is highly directional
156 and mostly due to the γ -rays generated in the moderator-decoupler system.
157 The other contributions are identified as a nearly isotropic one due to neutron
158 absorption in the walls of the experimental hall, and a directional one com-
159 ing from the beam dump. The isotropic source provides an almost uniform
160 background independent of position, while the beam dump, being a localized
161 source, mostly contributes in its proximity.

162 A previous measurement performed with a HpGe detector gave only time in-
163 tegrated information, while in the present case the information was integrated
164 over the whole γ -ray energies. The results obtained in the present paper vali-
165 date the assumption made in the previous measurement about the identifica-
166 tion of the possible γ -rays background sources. Furthermore they provide an
167 estimate of the relative importance of the different backgrounds components
168 at different times.

169 The results obtained, indicate that a further contribution to equation 1 should
170 be added to yield a more complete expression of the background-induced count
171 rate. This further contribution, that can be indicated as $B_{coll}(t)$, takes into
172 consideration the γ -rays coming from the neutron collimation system of the
173 instrument.

174 The approach used here in the VESUVIO beam line for background investiga-
175 tion can be extended at instruments operating at pulsed neutron sources and
176 especially at the new beam lines of TS2 at ISIS.

177 Acknowledgements

178 This work was supported within the CNR-CCLRC Agreement No. 01/9001
179 concerning collaboration in scientific research at the spallation neutron source
180 ISIS. The financial support of the Consiglio Nazionale delle Ricerche in this
181 research is hereby acknowledged.

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Fig. 1. Layout of the VESUVIO beam line at the ISIS spallation neutron source

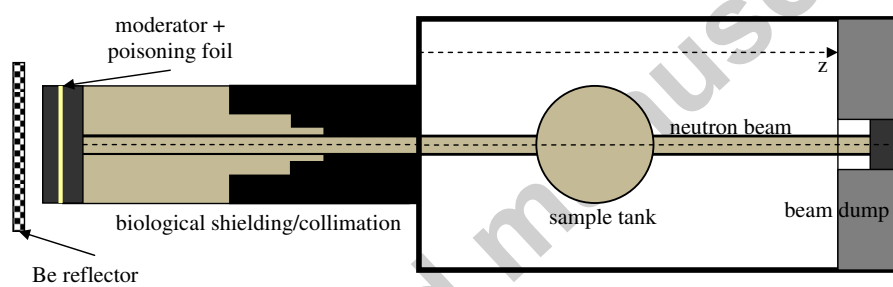


Fig. 2. time of flight spectra recorded by the YAP scintillator in different shielding configuration: (a) shielding aperture towards the moderator (total current $I = 413 \mu\text{Ah}$), (b) lateral aperture ($I = 100 \mu\text{Ah}$) and (c) rear aperture, i.e. towards the beam dump ($I = 1315 \mu\text{Ah}$).

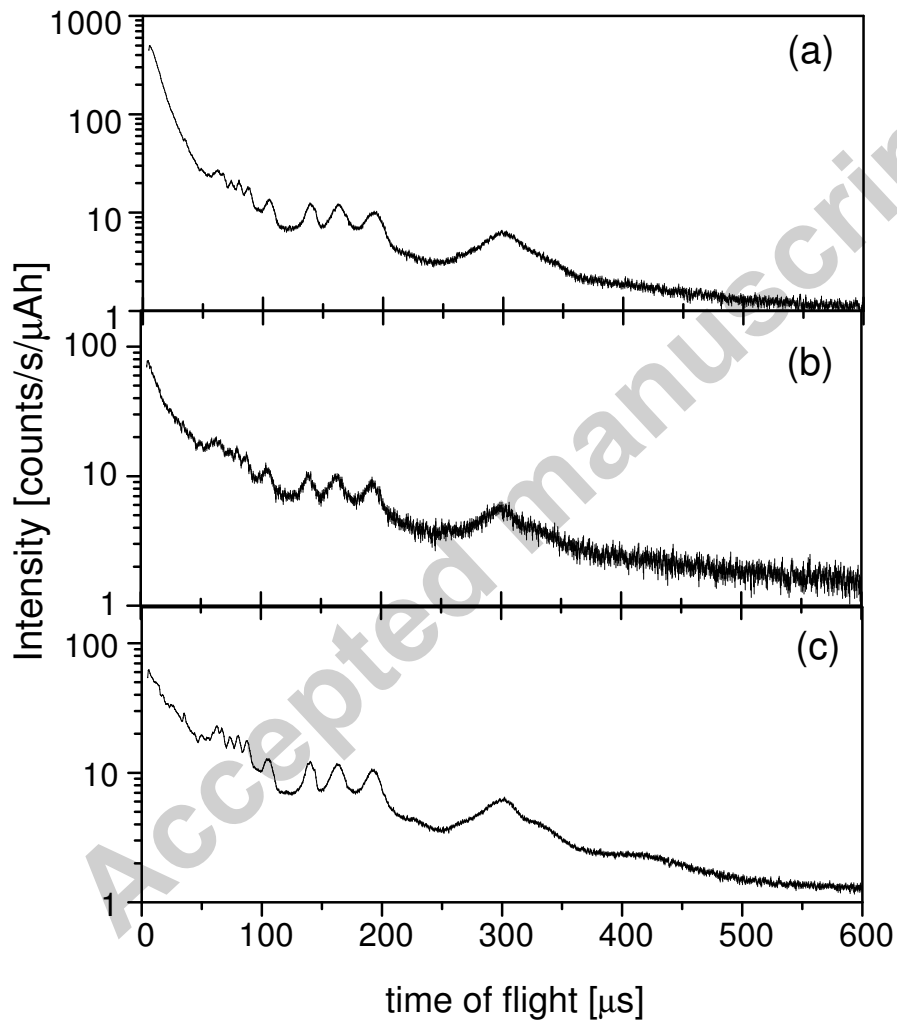


Fig. 3. Intensity ratio in the tof region up to $600 \mu\text{s}$ between spectra recorded with: (a) front and lateral shielding windows; (b) front and rear windows and (c) lateral and rear windows.

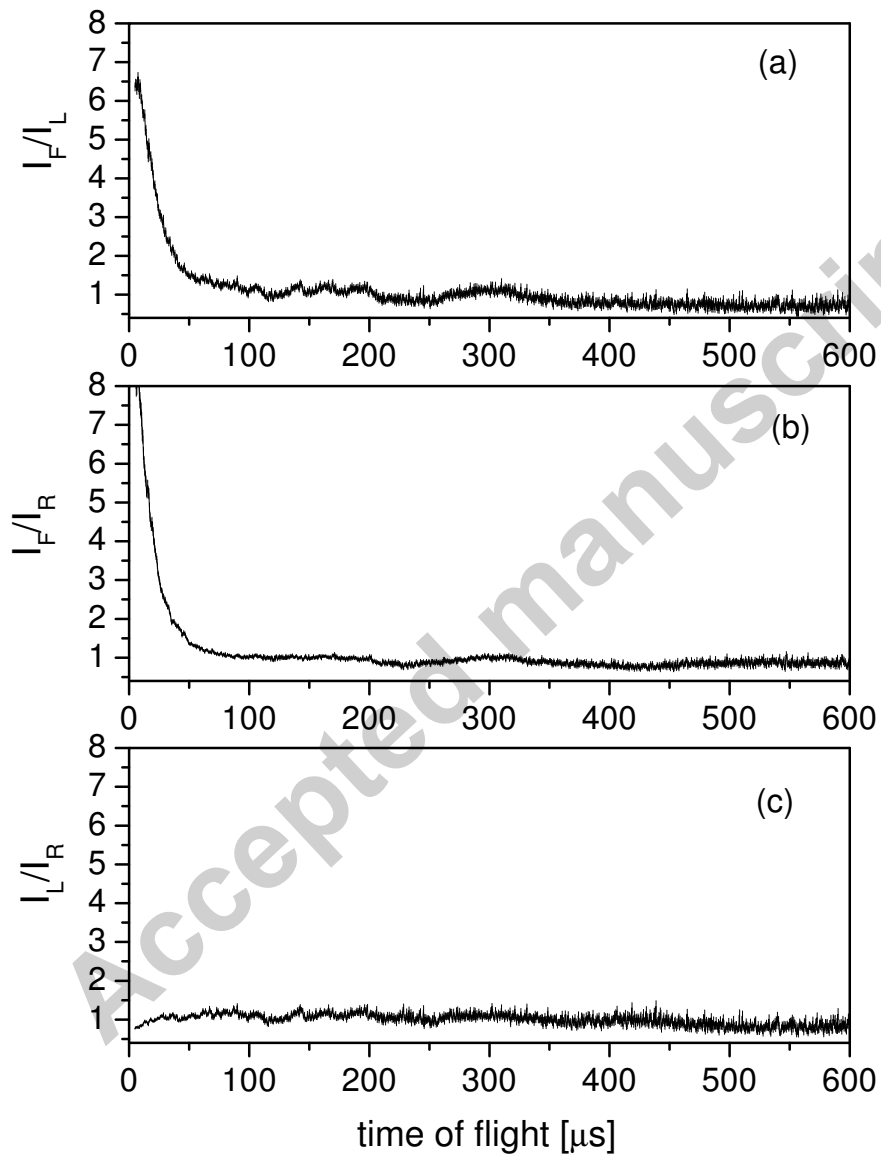


Fig. 4. Time of flight spectra recorded by the YAP detector at two different z positions, namely (a) in the proximity of the sample tank ($z=225$ cm) and (b) near the beam dump ($z=400$ cm). The vertical dashed lines define the tof intervals chosen to calculate the count rates (shown in figure 5) in the two regions.

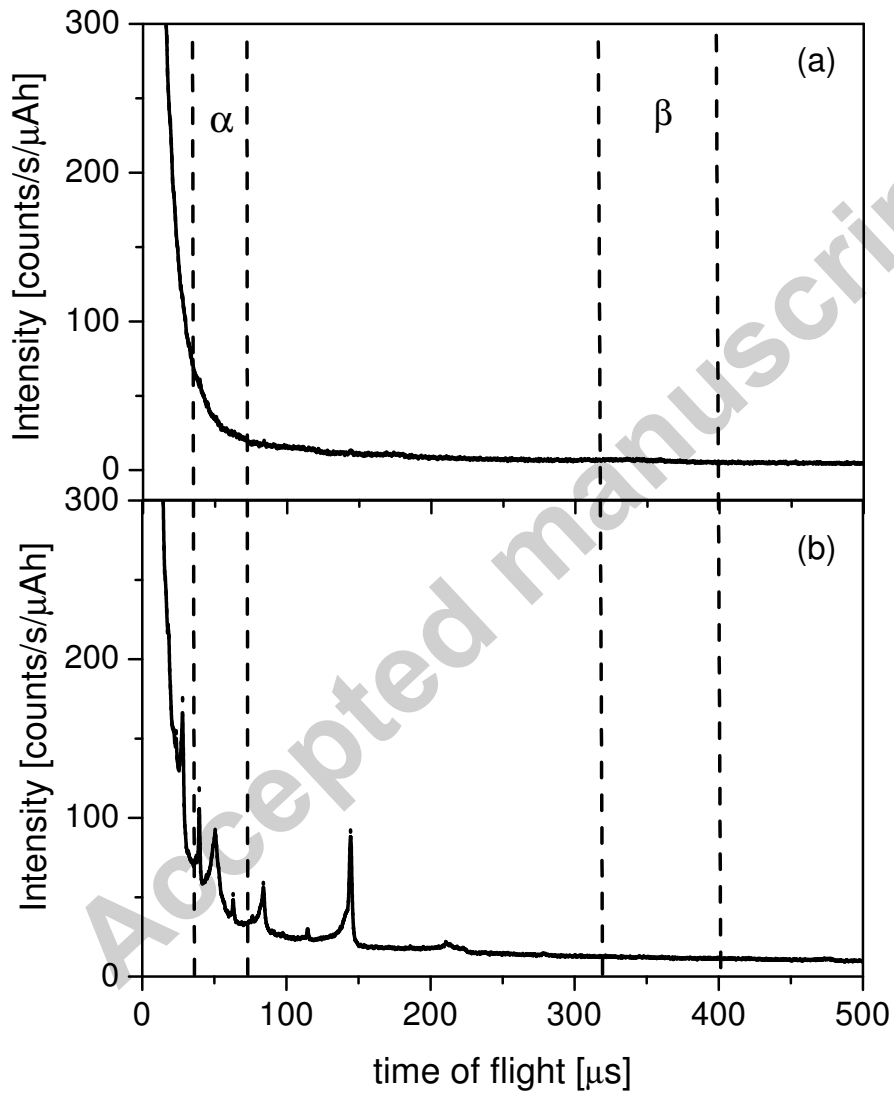


Fig. 5. Trend of the count rates (see text for details) in the regions identified in figure 4. The dashed line describes the trend of the intensity, normalized to the value at $z=225$ cm) from the moderator-decoupler system.

