

Article

Optimization and Characterization of High-Harmonic Generation for Probing Solid Density Plasmas

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Abstract: The creation of high energy density plasma states produced during laser–solid interaction on a sub-picosecond timescale opens a way to create astrophysical plasmas in the lab to investigate their properties, such as the frequency-dependent refractive index. Available probes to measure absorption and phase-changes given by the complex refractive index of the plasma state are extreme-UV (EUV) and soft X-ray (XUV) ultra-short pulses from high harmonic generation (HHG). For demanding imaging applications such as single-shot measurements of solid density plasmas, the HHG probe has to be optimized in photon number and characterized in intensity and wavefront stability from shot-to-shot. In an experiment, a coherent EUV source based on HHG driven by a compact diode-pumped laser is optimized in photons per pulse for argon and xenon, and the shot-to-shot intensity stability and wavefront changes are characterized. The experimental results are compared to an analytical model estimating the HHG yield, showing good agreement. The obtained values are compared to available data for solid density plasmas to confirm the feasibility of HHG as a probe.

Keywords: XUV; high-harmonic generation; high energy density physics; diode-pumped lasers; pump-probe

1. Introduction

High energy density plasmas are created in high power laser facilities worldwide. When a high power laser (Intensity $> 10^{14}$ W/cm²) interacts with a solid target, the target undergoes a transition from solid to plasma state. With ultra-short laser heating, the plasma passes through an intermediate state called warm dense matter (WDM), in which the density is nearly equal to solid density, but the temperature is on the order of $\sim 10^5$ K. This transient state exists for picoseconds [1,2]. Knowing how this transition is created, and understanding the dynamics of this intermediate high energy density plasma state is important for modeling how the laser energy is absorbed by solids. In return, this helps in understanding the laser energy deposition on targets in inertial confinement fusion (ICF) [3] and for

modeling the equations of state of astrophysical objects where the main constituent is WDM [1,2,4]. For high-density plasma states, it is important to have a probe which can penetrate the plasma with a sub-picosecond temporal resolution. Available probes are X-rays produced by free-electron lasers (FEL) [5], extreme-UV (EUV) and soft X-ray (XUV) [6] pulses produced by the process of high-harmonic generation (HHG), and soft X-ray lasers seeded by HHG [7,8]. Although X-rays from FELs can provide coherent XUV pulses with a high photon number (2×10^{13} photons/shot) [9] and coherence and peak intensities up to 10^{18} W/cm² [10], its major drawbacks are size, cost, and availability. The widely-available source of coherent XUV pulses is HHG, which also has a high spatial coherence and additionally supports pulse durations up to ~ 100 attoseconds [11]. Compared to FELs, the photon number of HHG is smaller (10^9 photons/shot), but it is sufficient to carry out single-shot experiments [12]. In order to use XUV pulses from HHG for different applications, it is important to optimize the photon number and to characterize the stability and spatial properties. For experiments like plasma probing [7,13], holography [14], coherent diffraction imaging [12], and X-ray laser seeding [15], it is necessary to have spatially characterized and energy-optimized XUV pulses. Optimization in photon number and wavefront is achieved by adjusting the phase matching conditions of the different high-harmonic orders [16–19]. To take full advantage of the ultrafast resolution achievable with HHG, synchronization between the high-harmonics and pump pulse from the high power laser beam-line is crucial. Therefore, HHG and pump pulse ideally originate from the same laser. A significant fraction of current high power lasers are flash-lamp-pumped, which results in a low pumping efficiency. Additionally, the active medium is heated up unnecessarily and only low repetition rates can be supported. An emerging key technology is diode-pumped laser systems, which overcome both limitations thanks to their unparalleled electrical-to-optical efficiency, and allow more compact setups that require less maintenance.

In the presented work, we use a diode-pumped Yb:CaF₂/Yb:YAG laser for HHG. The advantage of this laser system is the possibility of achieving high pulse energies at a high repetition rate [20,21], which allows for the simultaneous creation of a plasma and HHG. However, the laser pulses from Yb amplifiers are relatively long compared to conventional laser systems based on titanium-sapphire amplifiers, for which HHG is typically optimized in photon number and wavefront [22]. The number of photons and stability of HHG are optimized for the Yb amplifier-based laser system delivering ~ 1 ps pulses, and the wavefront is characterized. The experimental results for the photon yield in HHG are compared to an analytical model. Finally, the prospects for probing solid density plasmas with the obtained high-harmonic pulses are discussed.

2. Experimental Setup

Figure 1 shows a block diagram of the HHG experimental setup in the Laboratory for Intense Lasers (L2I) at the Instituto de Plasmas e Fusão Nuclear (IPFN). The in-house built chirped pulse amplification system consists of two diode-pumped amplification stages [23]. A regenerative amplifier with a Yb:CaF₂ crystal and a multi-pass amplifier with a Yb:YAG crystal are used as amplification stages, and both stages are pumped with diodes. The output of the laser system after compression is ~ 30 mJ pulse energy with a pulse duration of ~ 1000 fs at a wavelength of 1030 nm with a repetition rate of 2 Hz. This technology based on diode-pumped doped Yb amplifiers has the potential to be upgraded to joule-level pulse energies at a higher repetition rate [21]. This will enable the generation of pump and probe from the same laser. For HHG, a motorised plano-convex lens (f -number = $f/12$) inside the vacuum chamber is used to focus the incoming IR beam onto the gas cell. The generated XUV pulse is filtered from the residual IR pulse using a 150 nm-thick aluminum filter with supporting mesh. After the aluminum filter, an XUV spectrometer is placed which consists of a grazing incidence silica spherical mirror (grazing angle = 5° , radius of curvature = 10 m) followed by a gold transmission grating (1000 lines/mm) and an XUV-sensitive CCD camera (Princeton Instruments Pixis XO 1024B, Acton, MA, USA). For direct measurement of the XUV intensity profile, the XUV camera is placed directly after the filters without spectrometer. For wavefront measurement, the XUV camera is replaced

by a XUV wavefront sensor with alignment for tip/tilt and the vertical/horizontal axis [24,25] and a 300 nm-thick Al filter without supporting mesh is used.

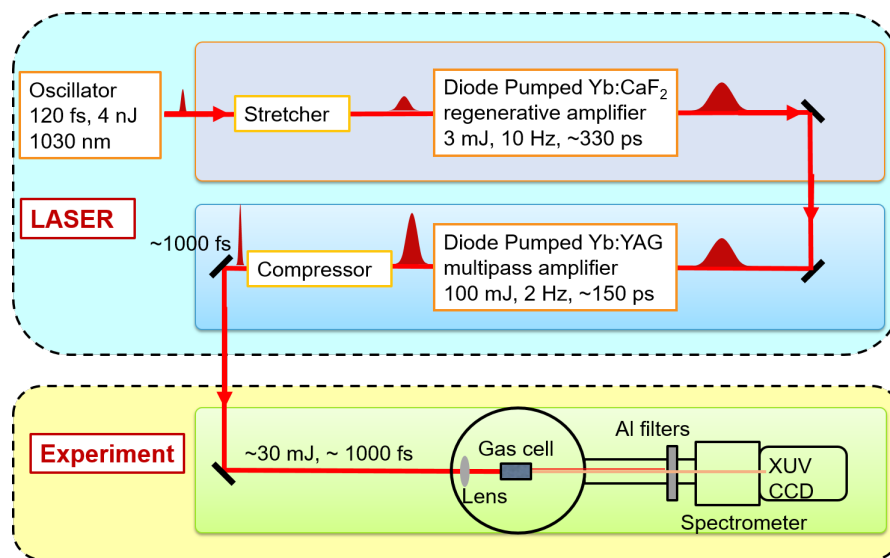


Figure 1. Schematic of high-harmonic generation (HHG) beamline with diode-pumped laser system at the Laboratory for Intense Lasers (L2I), Instituto de Plasmas e Fusão Nuclear (IPFN).

3. Results

3.1. HHG and Optimization

The high-harmonic photon count was optimized for two noble gases: argon and xenon. The optimization was done by changing the focal length of the lens, gas cell length, position of focus, gas pressure, and beam diameter before lens. The optimized total energy of all harmonic orders was calculated by considering the quantum efficiency and gain of the camera and central frequency of the high-harmonic spectrum. Furthermore, the Al filter used with a thickness of 150 nm on a nickel support grid was cross-calibrated with a second Al filter, which resulted in a transmission of 40% for wavelengths ~10–60 nm. In the case of both gases, the aperture before the lens was reduced to 9 mm and a pulse energy of ~10 mJ measured after the aperture. The gas-cell had a length of 10 mm, and the focus was placed at the exit of the gas-cell for optimized HHG. The optimization parameters and obtained results are listed in Table 1.

Table 1. Optimized operating conditions for HHG in argon and xenon calculated for the whole HH spectrum.

Gas	Rayleigh Length	Focal Length	Pressure	HH Photon Number	HH Energy
Units	mm	mm	mbar	Photons per Shot	nJ
Argon	3.3	500	55	7×10^7	0.44 ± 0.05
Xenon	6.4	700	15	3.6×10^8	1.44 ± 0.16

Figure 2 shows the optimized high harmonic beam cross-section and spectrum for argon and xenon. The periodic structure in the high harmonic beam cross-section results from a near-field diffraction effect from the nickel support grid of the filter.

The spatial profile and shot-to-shot stability of the high harmonics beam from argon and xenon was also measured, and the obtained values are listed in Table 2. Figure 3 shows the shot-to-shot stability analysis.

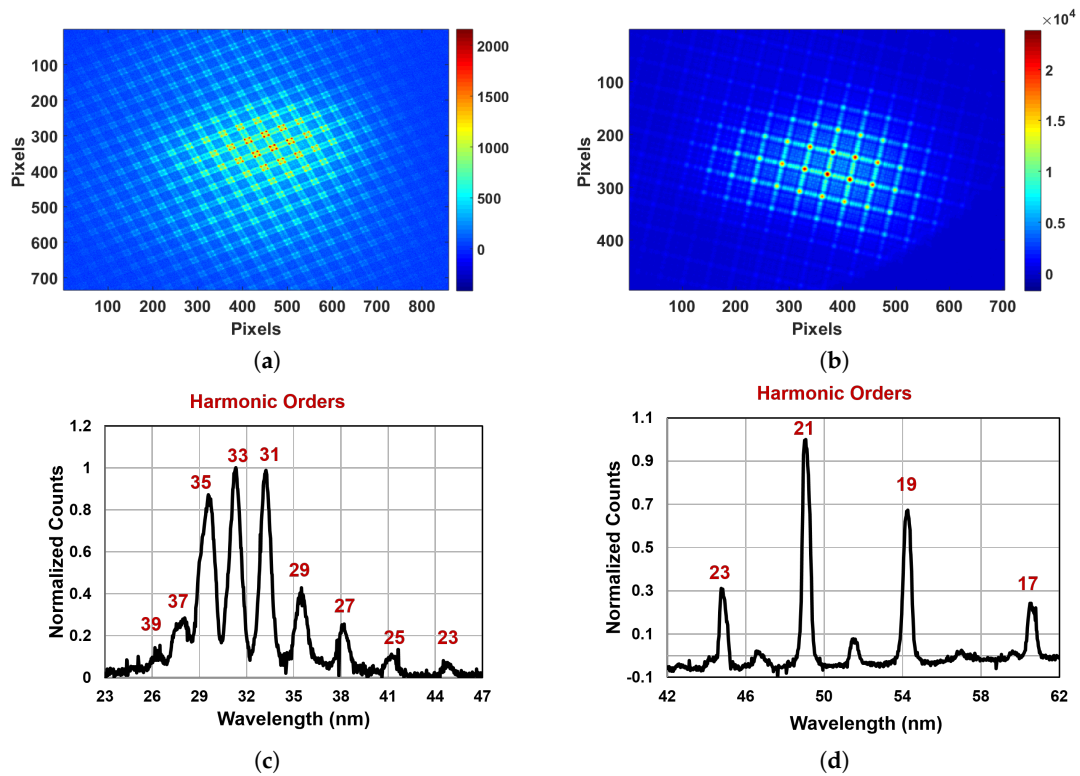


Figure 2. Beam cross-section (a,b) and spectrum (c,d) for HHG from argon and xenon respectively under optimized conditions in terms of number of photons and divergence. The smaller peaks <0.1 counts in the xenon spectrum result from the support grid of the grating.

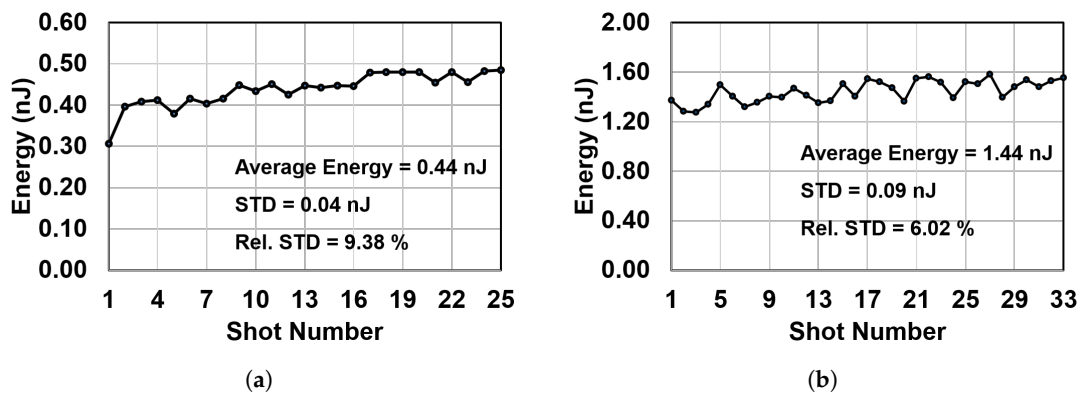


Figure 3. Single-shot energy stability analysis for high-harmonics generated in (a) argon and (b) xenon.

Table 2. Divergence and shot to shot energy stability of HHG in argon and xenon at L2I (calculated for the whole HH spectrum).

Gas	Divergence Horizontal	Divergence Vertical	Energy Stability
Units	mrاد	mrاد	%
Argon	2.48	2.18	9.4
Xenon	1.89	1.25	6.0

Experimental results are compared with a numerical calculation of the HH intensity considering macroscopic phase matching [17]. In non-optimized conditions, HHs were produced in argon using a lens with a focal length of 0.7 m. For the calculation of the HH intensities, the neutral dispersion,

plasma dispersion, Gouy phase, atomic phase, and re-absorption were considered. The ionization fraction for the dispersion was calculated using the ADK model [26]. For the experiment, a beam-waist of 64 μm was estimated. The focus was 1.3 mm after the gas-cell, which led to a peak intensity of $\sim 1.5 \times 10^{14} \text{ W/cm}^2$ in the focus and a calculated average ionization fraction of 4.8% in the gas-cell.

In Figure 4, the experimental measured intensities of HH order 23 to 39 are integrated and compared to the analytical calculated intensities for different pressures of argon in the gas-cell. Both show a maximum at the same pressure of 55 mbar. Furthermore, the calculated intensities reproduce the experimental curve. However, at pressures above 70 mbar, the calculations show a higher HH signal than the experimental values. This can be explained by the higher re-absorption of the HHs, since the vacuum system was not able to extract the argon efficiently from the chamber of the gas-cell at these pressures.

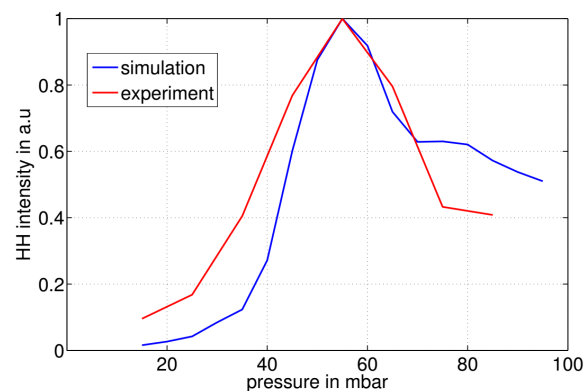


Figure 4. Comparison between simulated and experimental HH intensity depending on the gas-cell pressure with all HH orders integrated. The laser is focused 1.3 mm after the exit surface of the gas-cell.

The ionization fraction for three different laser intensities with varying pulse durations is plotted in Figure 5. It shows that the ionization fraction increases as the pulse duration increases. This results in a higher density of free electrons, which causes the phase-mismatch. Thus, the efficiency of HHG decreases. For the phase matching calculations, the electron density is obtained from calculations using the ADK model, which uses an average ionization. This assumption leads to a good agreement between calculated and experimental values, as shown in Figure 4. This also shows the validity of the model for use in the optimization of the HHG photon yield in the case of long pulse lasers.

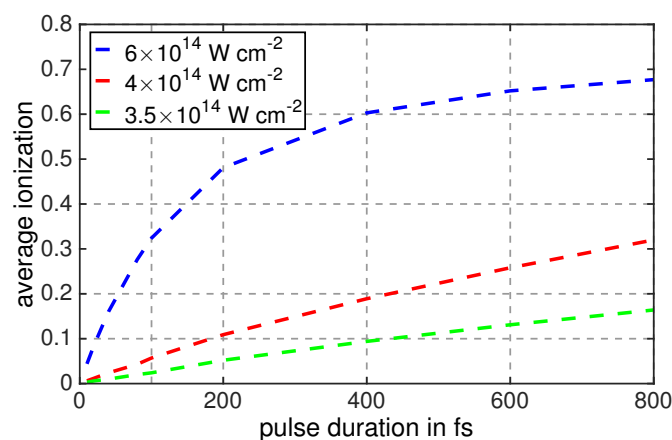


Figure 5. Calculated ionization fractions in argon produced by a laser at three different peak intensities and for different pulse durations.

3.2. Wavefront Measurement

Wavefront (WF) measurements were carried out for HHG in argon with the same optimized generation conditions as listed in Table 1, except the pressure was increased to 65 mbar. This was done to reduce the ellipticity of the beam while still maintaining the photon counts. The Al filter with supporting mesh was replaced by a 300 μm filter without mesh with a theoretical transmission of 49% [27] to avoid diffraction, which disturbs the wavefront measurement. The wavefront data was acquired by integrating the HH signal for 20 s. In a previous study, it was shown that the integration of wavefronts of multiple shots will not change the root mean square (RMS) wavefront error [19]. The XUV wavefront sensor (WFS) has a Hartmann plate with 57×57 apertures with a pitch of 250 μm . The sensitivity of the WFS is $\lambda/30$ for the spectral range used in the experiment. The full width half maximum (FWHM) of the HH beam diameter is about 2 mm \times 3 mm at the CCD of the WFS. In total, 13 consecutive HH wavefronts were recorded in the same conditions for HHG and compared to estimate the stability of the wavefront. The HH wavefronts were reconstructed using the software MrBeam [24]. Prior to wavefront reconstruction, a dark image was subtracted and hot pixels were removed for each measurement. Hot pixels arose due to the long exposure time. For the reconstruction, a rectangle of 32×23 apertures was selected and centered on recorded HH intensity, which covers the full aperture of the beam. An example for a recorded wavefront is shown in Figure 6. For each measurement, the RMS wavefront error, astigmatism, and coma was calculated; furthermore, the RMS of the beam slopes in the selected area was also calculated. The smallest variation of the RMS wavefront error was 1.14 nm, which corresponds to $\lambda/28$ and is therefore within the detection limit of the WFS. Additionally, the mean values and their standard deviation were calculated. Figure 7 shows the evolution of the RMS wavefront error, astigmatism, and coma for the different acquisitions, and the mean values are summarized in Table 3.

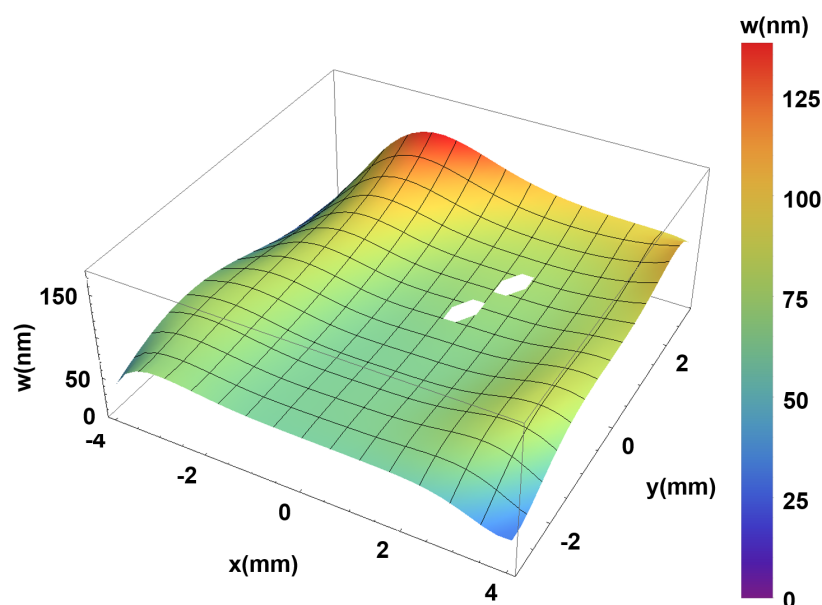


Figure 6. An example of a reconstructed high harmonic wavefront (acquisition 1) measured with the Hartmann type XUV wavefront sensor. One white spot arose from a closed hole in the center of the Hartmann plate. It is used for alignment of the WFS to the beam. The other white spot occurred due to a small particle covering one hole of the plate accidentally.

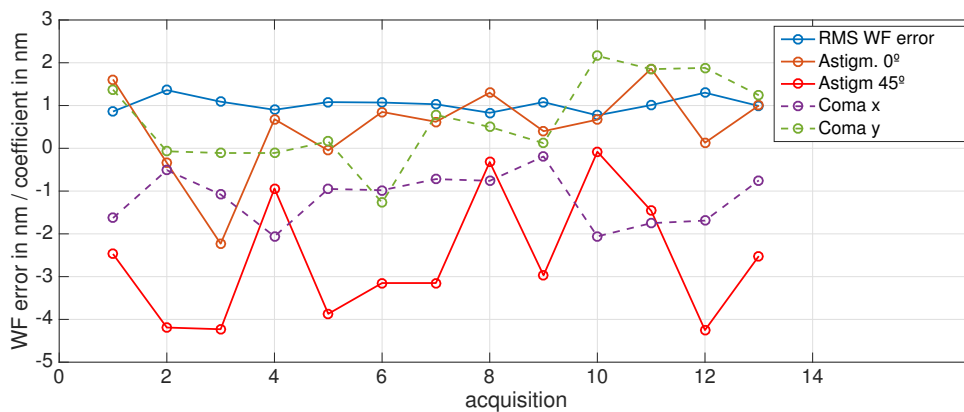


Figure 7. The evolution of the root mean square (RMS) wavefront (WF) error and the Zernike coefficients of the main aberrations.

Table 3. Mean values for the RMS wavefront error and the Zernike coefficients over 13 acquisitions.

Parameter	Mean	Standard Deviation
RMS Wavefront Error	4.90 nm	1.14 nm
Astigmatism 0°	−41.50 nm	18.48 nm
Astigmatism 45°	5.12 nm	9.55 nm
Coma X	9.37 nm	6.86 nm
Coma Y	−21.85 nm	35.03 nm
RMS Slope	0.60 mrad	0.01 mrad

4. Discussion

HHG from argon and xenon were optimized in terms of energy with single shot energies of 0.44 ± 0.05 nJ and 1.44 ± 0.16 nJ, respectively. The XUV pulses from HHG were characterized in spectrum, spatial profile, wavefront, and the shot-to-shot stability in energy and integrated wavefront stability. A shot-to-shot stability of 9.4% was measured for argon, and 6.0% for HHG in xenon. From the spectral analysis of the HHG from xenon, we can find that the brightest HH for xenon is order 21 with $\sim 44\%$ of the total HH energy. In the spectrum of argon HH order 35, 33, and 31 have nearly equal intensities. All three together share $\sim 71\%$ of the total HH energy. Analytical calculations for the HH emission confirm the optimized parameters for pulse energy, gas pressure, focal length, and beam diameter. Therefore, the numerical calculations considering the macroscopic phase matching conditions are an important tool for low repetition rate laser facilities, where extensive optimization scans are not practical.

In [13] values for the absorption and deflection of a XUV probe pulse in Al heated to warm-dense conditions are estimated. The referenced values result from heating Al with an FEL, however WDM can also be produced with optical lasers as shown in [28–30]. At a temperature of 10 eV, an absorption coefficient of $\alpha \approx 4.5 \times 10^6$ is estimated for photons with an energy of 30 eV (~ 41 nm), which leads to an absorption of $\sim 75\%$ in a 300 nm-thick foil. More absorption is expected for lower photon energies. In xenon, the HH order 23 with a wavelength of 44.7 nm has an energy of 13% of the whole spectrum, which corresponds to ~ 0.18 nJ. This corresponds to $\sim 2 \times 10^6$ counts in the XUV CCD camera used in the experiment. For an area of 100×100 pixels on the CCD, this relates to 200 counts per pixel, which is sufficient to be recorded in a single-shot acquisition. The deflection $\phi(x)$ of a beam in a plasma is given by the gradient of the refractive index ∇n and the length of the plasma L by

$$\phi(x) = \int_0^L \nabla n(x, z) dz, \tag{1}$$

where x is the transverse direction to the plasma and z is the perpendicular direction. Following [13], the largest deflection of an HH probe beam—larger than the plasma created at WDM conditions—is between the cold Al foil and warmest part of the solid plasma at a temperature of 10 eV. For a plasma with a diameter of 20 μm , a maximum beam deflection of 0.93 mrad is expected under these conditions. The measured mean slope of the characterized HH source was 0.6 mrad with a standard deviation of 10 μrad , which is smaller than the expected beam deflections. Therefore, the plasma-induced deflections are more significant than the fluctuations of the probe beam.

The narrower amplification bandwidth of diode-pumped Yb amplifiers leads to a Fourier-transform-limited pulse duration of a few hundred fs [31,32], which is not a typical parameter space for the optimization of HHG. The obtained results for the HH photon number per pulse and wavefront show the potential of using laser systems with Yb amplifiers for simultaneously creating plasmas and a probe based on HHG. Additionally, the relatively long pulse duration is advantageous for the plasma creation for X-ray lasers, which can be seeded by HHG to achieve a higher brilliance [15].

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