



Non-integrable Ising Models in Cylindrical Geometry: Grassmann Representation and Infinite Volume Limit

Giovanni Antinucci, Alessandro Giuliani  and Rafael L. Greenblatt

Abstract. In this paper, meant as a companion to Antinucci et al. (Energy correlations of non-integrable Ising models: the scaling limit in the cylinder, 2020. [arXiv: 1701.05356](https://arxiv.org/abs/1701.05356)), we consider a class of non-integrable 2D Ising models in cylindrical domains, and we discuss two key aspects of the multiscale construction of their scaling limit. In particular, we provide a detailed derivation of the Grassmann representation of the model, including a self-contained presentation of the exact solution of the nearest neighbor model in the cylinder. Moreover, we prove precise asymptotic estimates of the fermionic Green's function in the cylinder, required for the multiscale analysis of the model. We also review the multiscale construction of the effective potentials in the infinite volume limit, in a form suitable for the generalization to finite cylinders. Compared to previous works, we introduce a few important simplifications in the localization procedure and in the iterative bounds on the kernels of the effective potentials, which are crucial for the adaptation of the construction to domains with boundaries.

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1. Introduction

In this article, which is a companion to [4], we consider a class of non-integrable perturbation of the 2D nearest-neighbor Ising model in cylindrical geometry and discuss some of the key ingredients required in the multiscale construction of the scaling limit of the energy correlations in finite domains. The material presented here generalizes and simplifies the approach proposed by two of the authors in [18], where a similar problem in the translationally invariant setting was investigated. As discussed extensively in [4, Section 3.1], which we refer to for additional motivations and references, the methods of [18], as well as of several other related works on the Renormalization Group (RG) construction of the bulk scaling limit of non-integrable lattice models at the critical point, are insufficient for controlling the effects of the boundaries at the precision required for the construction of the scaling limit in finite domains. This is a serious obstacle in the program of proving conformal invariance of the scaling limit of statistical mechanics models [17]; the goal would be to prove results comparable to the remarkable ones obtained for the nearest neighbor 2D Ising model [10, 12, 29], but for a class of non-integrable models, such as perturbed Ising [2] or dimer models [22] in two dimensions, via methods that do not rely on the exact solvability of the microscopic model. In this paper and in its companion [4], we attack this program by constructing the scaling limit of the energy correlations of a class of non-integrable perturbations of the standard 2D Ising model in the simplest possible finite domain with boundary, that is, a finite cylinder.

Let us define the setting more precisely. For positive integers L and M , with L even, we let G_Λ be the discrete cylinder with sides L and M in the horizontal and vertical directions, respectively, with periodic boundary conditions in the horizontal direction and open boundary conditions in the vertical directions. We consider G_Λ as a graph with vertex set $\Lambda = \mathbb{Z}_L \times (\mathbb{Z} \cap [1, M])$, where $\mathbb{Z}_L = \mathbb{Z}/L\mathbb{Z}$ (in the following we shall identify the elements of \mathbb{Z}_L with $\{1, \dots, L\}$, unless otherwise stated) and edge set \mathfrak{B}_Λ consisting of all pairs of the form¹ $\{z, z + \hat{e}_j\}$ for $z \in \Lambda, j \in \{1, 2\}$ and \hat{e}_1, \hat{e}_2 the unit vectors in the two coordinate directions. For $x \in \mathfrak{B}_\Lambda$, we let $j(x)$ be the j for which $x = \{z, z + \hat{e}_j\}$ for some $z \in \Lambda$, so that $j(x) = 1$ for a horizontal bond and $j(x) = 2$ for a vertical bond. The model is defined by the Hamiltonian

$$H_\Lambda(\sigma) = - \sum_{x \in \mathfrak{B}_\Lambda} J_{j(x)} \epsilon_x - \lambda \sum_{X \subset \Lambda} V(X) \sigma_X, \tag{1.1}$$

where J_1, J_2 are two positive constants, representing the couplings in the horizontal and vertical directions, $\epsilon_x = \epsilon_x(\sigma) := \sigma_z \sigma_{z'}$ for $x = \{z, z'\}$; the spin variable σ belongs to $\Omega_\Lambda := \{\pm 1\}^\Lambda$, and $\sigma_X := \prod_{x \in X} \sigma_x$; V is a finite range, translationally invariant, even interaction, obtained by periodizing in the horizontal direction a Λ -independent, translationally invariant, potential on \mathbb{Z}^2 ; finally, λ is the strength of the interaction, which can be of either sign and, for most of the discussion below, the reader can think of as being small, compared to J_1, J_2 , but independent of the system size. In the following, we shall refer to model (1.1) with $\lambda \neq 0$ as to the ‘interacting’ model, in contrast with the standard nearest-neighbor model, which we will refer to as the ‘non-interacting’, one of several terminological conventions motivated by analogy with quantum field theory. The Hamiltonian defines a Gibbs measure $\langle \cdot \rangle_{\beta, \Lambda}$ depending on the inverse temperature $\beta > 0$, which assigns to any $F : \Omega_\Lambda \rightarrow \mathbb{R}$ the expectation value

$$\langle F \rangle_{\beta, \Lambda} := \frac{\sum_{\sigma \in \Omega_\Lambda} e^{-\beta H_\Lambda(\sigma)} F(\sigma)}{\sum_{\sigma \in \Omega_\Lambda} e^{-\beta H_\Lambda(\sigma)}}. \tag{1.2}$$

The truncated correlations, or cumulants, of the energy observable ϵ_x , denoted $\langle \epsilon_{x_1}; \dots; \epsilon_{x_n} \rangle_{\beta, \Lambda}$, are given by

$$\langle \epsilon_{x_1}; \dots; \epsilon_{x_n} \rangle_{\beta, \Lambda} := \frac{\partial^n}{\partial A_1 \dots \partial A_n} \log \langle e^{A_1 \epsilon_{x_1} + \dots + A_n \epsilon_{x_n}} \rangle_{\beta, \Lambda} \Big|_{A_1 = \dots = A_n = 0}. \tag{1.3}$$

For the formulation of the main result, let us fix once and for all an interaction V with the properties spelled out after (1.1), and assume that J_1/J_2 and L/M belong to a compact $K \subset (0, +\infty)$. We let $t_l := t_l(\beta) := \tanh \beta J_l$, with $l = 1, 2$, and recall that in the non-interacting case, $\lambda = 0$, the critical temperature $\beta_c(J_1, J_2)$ is the unique solution of $t_2(\beta) = (1 - t_1(\beta))/(1 + t_1(\beta))$. Note that there exists a suitable compact $K' \subset (0, 1)$ such that whenever $J_1/J_2 \in K$ and $\beta \in [\frac{1}{2}\beta_c(J_1, J_2), 2\beta_c(J_1, J_2)]$, then $t_1, t_2 \in K'$. From now on, we will

¹If $z = ((z)_1, (z)_2) \in \Lambda$ has horizontal coordinate $(z)_1 = L$, we use the convention that $z + \hat{e}_1 \equiv (1, (z)_2)$.

think K, K' to be fixed once and for all. Moreover, we parameterize the Gibbs measure in terms of t_l as follows:

$$\langle \cdot \rangle_{\beta, \Lambda} \equiv \langle \cdot \rangle_{\lambda, t_1, t_2; \Lambda}.$$

Given these premises, we are ready to state the main result proven in [4].

Theorem 1.1. *Fix V as discussed above. Fix J_1, J_2 so that J_1/J_2 belongs to the compact K introduced above. There exist $\lambda_0 > 0$ and analytic functions $\beta_c(\lambda), t_1^*(\lambda), Z_1(\lambda), Z_2(\lambda)$, defined for $|\lambda| \leq \lambda_0$, such that, for any finite cylinder Λ with $L/M \in K$ and any m -tuple $\mathbf{x} = (x_1, \dots, x_m)$ of distinct elements of \mathfrak{B}_Λ , with m_1 horizontal elements, m_2 vertical elements, and $m = m_1 + m_2 \geq 2$,*

$$\langle \epsilon_{x_1}; \dots; \epsilon_{x_m} \rangle_{\lambda, t_1(\lambda), t_2(\lambda); \Lambda} = (Z_1(\lambda))^{m_1} (Z_2(\lambda))^{m_2} \langle \epsilon_{x_1}; \dots; \epsilon_{x_m} \rangle_{0, t_1^*(\lambda), t_2^*(\lambda); \Lambda} + R_\Lambda(\mathbf{x}), \tag{1.4}$$

where $t_1(\lambda) := \tanh(\beta_c(\lambda)J_1)$, $t_2(\lambda) := \tanh(\beta_c(\lambda)J_2)$ and $t_2^*(\lambda) := (1 - t_1^*(\lambda))/(1 + t_1^*(\lambda))$. Moreover, denoting by $\delta(\mathbf{x})$ the tree distance of \mathbf{x} , i.e., the cardinality of the smallest connected subset of \mathfrak{B}_Λ containing the elements of \mathbf{x} , and by $d = d(\mathbf{x})$ the minimal pairwise distance among the midpoints of the edges in \mathbf{x} and the boundary of Λ , for all $\theta \in (0, 1)$ and $\varepsilon \in (0, 1/2)$ and a suitable $C_{\theta, \varepsilon} > 0$, the remainder R_Λ can be bounded as

$$|R_\Lambda(\mathbf{x})| \leq C_{\theta, \varepsilon}^m |\lambda| m! \frac{1}{d^{m+\theta}} \left(\frac{d}{\delta(\mathbf{x})} \right)^{2-2\varepsilon}. \tag{1.5}$$

As a corollary of this theorem, one readily obtain the existence and explicit structure of the scaling limit for the ‘energy sector’ of the interacting model, with quantitative estimates on the speed of convergence; see [4, Corollary 1.2] and Appendix D.

The proof of Theorem 1.1 is based on a multiscale analysis of the generating function of the energy correlations, formulated in the form of a Grassmann (Berezin) integral. While the strategy of this proof is based on the same general ideas used in [18] in the translationally invariant setting, that is, on the methods of the *fermionic constructive RG*, the presence of boundaries introduces several technical and conceptual difficulties, whose solution requires to adapt, improve and generalize the ‘standard’ RG procedure (e.g., in the definition of the ‘localization procedure’, in the way in which the kernels of the ‘effective potentials’ are iteratively bounded and in which the resulting bounds are summed over the label specifications, etc.) As discussed in [4, Section 3.1], which we refer to for additional details, we expect that understanding how to implement RG in the presence of boundaries or, more in general, of defects breaking translational invariance, will have an impact on several related problems, such as the computation of boundary critical exponents in models in the Luttinger liquid universality class, the Kondo problem, the Casimir effect, and the phenomenon of many-body localization.

In this paper we give a full presentation of some of the key ingredients required in the proof of our main result, namely:

1. exact solution of the nearest neighbor model on the cylinder in its Grassmann formulation, including multiscale bounds on the bulk and edge parts of the fermionic Green's function (Sect. 2);
2. reformulation of the generating function of energy correlations of the interacting model as a Grassmann integral (Sect. 3);
3. tree expansion and iterative bounds on the kernels of the effective potentials of the interacting theory in the full plane limit, including the computation and proof of analyticity of the interacting critical temperature (Sect. 4).

The other ingredients, including most of the novel aspects of the RG construction in finite volume, such as the definition of the localization procedure in finite volume, the norm bounds on the edge part of the effective potentials and the asymptotically sharp estimates on the correlation functions in the cylinder, are deferred to [4]; see the end of [4, Section 3.1] for a detailed summary and roadmap of the proof of Theorem 1.1.

Before starting the technical presentation, let us anticipate in little more detail the contents of the following sections, thus clarifying the main results of this paper.

Section 2: exact solution of the model in the cylinder. The multiscale construction of the interacting theory in the domain Λ requires a very fine control of the non-interacting model at the critical point, and, in particular, of the structure of its fermionic Green's function, which we call the 'propagator'; the propagator is nothing but the inverse of a signed adjacency matrix A , whose definition we recall in Sect. 2 below [26, Chapter IX]. The key properties we need, and we prove in Sect. 2 below (with some—important!—technical aspects of the proofs deferred to Appendices A, B and C), see, in particular, Eq. (2.2.14) and Proposition 2.3 below, are the following:

- multiscale decomposition of the propagator and bulk–edge decomposition of the single-scale propagator;
- exponentially decaying pointwise bounds on the bulk and edge parts of the single-scale propagators, with optimal dimensional bounds (with respect to the scale index) on their L^∞ norms and on their decay rates;
- Gram representation² of the bulk and edge parts of the single-scale propagators, with optimal dimensional bounds (with respect to the scale index) on the norms of the Gram vectors.

In reference with the second item, let us remark that the exponential decay needed (and proved below) for the propagator between two points $z, z' \in \Lambda$, is in terms of the 'right' distance between z and z' , namely: the standard Euclidean distance on the cylinder between z and z' in the case of the bulk part of the single-scale propagator; the Euclidean distance on the cylinder between z, z' and the boundary of Λ , in the case of the edge part of the single-scale propagator. In particular, the exponential decay of the edge part

²We say that a matrix g admits a Gram representation, if its elements $g_{i,j}$ can be written as the scalar product of two vectors u_i and v_j in a suitable Hilbert space.

of the single-scale propagator in the distance of z, z' from the boundary of Λ is of crucial importance for proving improved dimensional bounds on the finite-size corrections to the thermodynamic and correlation functions of the interacting model, which are systematically used in the conclusion of the proof of Theorem 1.1 in the companion paper [4], see [4, Section 4].

The proof we give of these key properties is based on an exact diagonalization of the signed adjacency matrix A in terms of the roots of a set of polynomials (this calculation first appeared in [23], and a similar calculation for a rectangle appears in [25]). It is unlikely that such an explicit diagonalization can be obtained in more general domains than the torus, the straight cylinder or the rectangle. Therefore, in order to generalize Theorem 1.1 to more general domains, it would be desirable to prove the properties summarized in the three items above via a more robust method, not based on an explicit diagonalization of A . It remains to be seen whether the methods of discrete holomorphicity, which allowed to prove the convergence of the propagator in general domains to an explicit, conformally covariant, limiting function [10], may allow one to prove the desired properties in general domains.

Section 3: Grassmann representation of the generating function. In Sect. 3, we turn our attention to the generating function for the energy correlations

$$Z_\Lambda(\mathbf{A}) := \sum_{\sigma \in \Omega_\Lambda} \exp \left(\sum_{x \in \mathfrak{B}_\Lambda} [\beta J_{j(x)} + A_x] \epsilon_x + \beta \lambda \sum_{X \subset \Lambda} V(X) \sigma_X \right). \tag{1.6}$$

that, if computed at a configuration \mathbf{A} such that A_x is equal to A_i for $x = x_i$ and zero otherwise, reduces to the combination $\langle e^{A_1 \epsilon_{x_1} + \dots + A_n \epsilon_{x_n}} \rangle_{\beta, \Lambda}$ appearing in (1.3), up to an overall multiplicative constant, independent of \mathbf{A} .

In Proposition 3.1 and Eq.(3.22) (adapting a similar result for the torus in [18]), we show that the correlations without repeated bonds are the same as those obtained by replacing $Z_\Lambda(\mathbf{A})$ with a Grassmann integral of the form

$$\Xi_\Lambda(\mathbf{A}) := e^{\mathcal{W}(\mathbf{A})} \int P_c^*(\mathcal{D}\phi) P_m^*(\mathcal{D}\xi) e^{\mathcal{V}^{(1)}(\phi, \xi, \mathbf{A})}, \tag{1.7}$$

where P_c^* and P_m^* are Gaussian Grassmann measures associated with the critical, non-interacting Ising model at parameters $t_1^*, t_2^* := (1 - t_1^*)/(1 + t_1^*)$, with t_1^* a free parameter. Moreover, $\mathcal{W}(\mathbf{A})$ is a multilinear function of \mathbf{A} and $\mathcal{V}^{(1)}(\phi, \xi, \mathbf{A})$ is a Grassmann polynomial whose coefficients are multilinear functions of \mathbf{A} , both of which are defined in terms of explicit, convergent, expansions. As a corollary of Lemma 3.2, we additionally prove that the ‘kernels’ of $\mathcal{W}(\mathbf{A})$ and $\mathcal{V}^{(1)}(\phi, \xi, \mathbf{A})$ (i.e., the coefficients of their expansions in \mathbf{A}, ϕ, ξ , thought of as functions of the positions of the components of \mathbf{A}, ϕ, ξ on the cylinder) can be naturally decomposed into sums of a ‘bulk’ part (equal, essentially, to their infinite plane limit restricted to the cylinder, with the appropriate boundary conditions) plus an ‘edge’ part (their boundary corrections), exponentially decaying in the appropriate distances. In particular, the edge part of the kernels decays exponentially (on the lattice scale) away from

the boundary, a fact that will play a major role in the control of the boundary corrections to the correlation functions in [4].

Let us remark that, in addition to t_1^* and to the inverse temperature β , the representation (1.7) has another free parameter, Z (entering the definition of $\mathcal{V}^{(1)}$); while, for the validity of (1.7), these parameters can be chosen arbitrarily in certain intervals, in order for this representation to produce a convergent expansion for the critical energy correlations of the interacting model, uniformly in the system size, we will need to fix t_1^*, β, Z appropriately (a posteriori, they will be fixed uniquely by our construction, see below). Parameters of this kind are known as *counterterms* in the RG terminology.

Section 4: the RG expansion for the effective potentials in the full plane limit.

Equation (1.7) is the starting point for a multiscale expansion, which is fully presented in the companion paper [4], see in particular [4, Section 3], but which we summarize here in order to provide the context for Sect. 4, where we carry out an auxiliary expansion for the full plane limit of the kernels of the ‘effective potentials’. Such an auxiliary expansion, among other things, fixes the values of t_1^*, β, Z which are actually used in Theorem 1.1, see Sect. 4.5 below.

The goal is to iteratively compute (1.7) in terms of a sequence of effective potentials, defined as follows: at the first step we let

$$e^{\mathcal{W}^{(0)}(\mathbf{A}) + \mathcal{V}^{(0)}(\phi, \mathbf{A})} \propto \int P_m^*(\mathcal{D}\xi) e^{\mathcal{V}^{(1)}(\phi, \xi, \mathbf{A})}, \tag{1.8}$$

where \propto means ‘up to a multiplicative constant independent of \mathbf{A} ’; the polynomials $\mathcal{W}^{(0)}, \mathcal{V}^{(0)}$ are specified uniquely by the normalization $\mathcal{W}^{(0)}(\mathbf{0}) = \mathcal{V}^{(0)}(0, \mathbf{A}) = 0$.

We are left with computing the integral of $e^{\mathcal{V}^{(0)}(\phi, \mathbf{A})}$ with respect to the Gaussian integration $P_c^*(\mathcal{D}\phi)$ with propagator \mathfrak{g}_c^* . As anticipated above, in Sect. 2.2 we decompose the critical propagator \mathfrak{g}_c^* as $\mathfrak{g}^{(\leq h)} + \sum_{j=h+1}^0 \mathfrak{g}^{(j)}$, for any $h < 0$; correspondingly, in light of the addition formula for Grassmann integrals (see, e.g., [21, Proposition 1]), we introduce the sequences $P^{(\leq h)}$ and $P^{(h)}$ of Gaussian Grassmann integrations, whose propagators are $\mathfrak{g}^{(\leq h)}$ and $\mathfrak{g}^{(h)}$, respectively, and satisfy, for any Grassmann function f ,

$$\int P^{(\leq h)}(\mathcal{D}\phi) f(\phi) = \int P^{(\leq h-1)}(\mathcal{D}\phi) P^{(h)}(\mathcal{D}\varphi) f(\phi + \varphi). \tag{1.9}$$

We can then iteratively define $\mathcal{V}^{(h)}$ and $\mathcal{W}^{(h)}$ with $\mathcal{W}^{(h)}(\mathbf{0}) = \mathcal{V}^{(h)}(0, \mathbf{A}) \equiv 0$ and

$$e^{\mathcal{W}^{(h-1)}(\mathbf{A}) + \mathcal{V}^{(h-1)}(\phi, \mathbf{A})} \propto \int P^{(h)}(\mathcal{D}\varphi) e^{\mathcal{V}^{(h)}(\phi + \varphi, \mathbf{A})}. \tag{1.10}$$

The iteration continues until the scale $h^* = -\lfloor \log_2(\min\{L, M\}) \rfloor$ is reached, at which point we let

$$e^{\mathcal{W}^{(h^*-1)}(\mathbf{A})} \propto \int P^{(\leq h^*)}(\mathcal{D}\phi) e^{\mathcal{V}^{(h^*)}(\phi, \mathbf{A})}, \tag{1.11}$$

giving

$$\Xi_\Lambda(\mathbf{A}) \propto \exp \left(\mathcal{W}(\mathbf{A}) + \sum_{h=h^*-1}^0 \mathcal{W}^{(h)}(\mathbf{A}) \right). \tag{1.12}$$

In order to obtain bounds on the kernels of $\mathcal{W}^{(h)}(\mathbf{A})$ leading to an expansion for the energy correlations that is uniform in the system size, at each step it is necessary to isolate from $\mathcal{V}^{(h)}$ the contributions that tend to expand (in an appropriate norm) under iterations: these, in the RG terminology, are the relevant and marginal terms, which we collect in the so-called *local part* of $\mathcal{V}^{(h)}$, denoted by $\mathcal{L}\mathcal{V}^{(h)}$. In other words, at each step of the iteration, we rewrite $\mathcal{V}^{(h)} = \mathcal{L}\mathcal{V}^{(h)} + \mathcal{R}\mathcal{V}^{(h)}$, where, in our case, $\mathcal{L}\mathcal{V}^{(h)}$ includes: three terms that are quadratic in the Grassmann variables and independent of \mathbf{A} , depending on a sequence of h -dependent parameters which we denote $\underline{\nu} = \{(\nu_h, \zeta_h, \eta_h)\}_{h \leq 1}$ and call the *running coupling constants*; and two terms that are quadratic in the Grassmann variables and linear in \mathbf{A} , depending on another sequence of effective parameters, $\{Z_{1,h}, Z_{2,h}\}_{h \leq 0}$, called the *effective vertex renormalizations*. Moreover, $\mathcal{R}\mathcal{V}^{(h)}$ is the so-called *irrelevant*, or *renormalized*, part of the effective potential, which is not the source of any divergence.

Such a decomposition corresponds to a systematic reorganization, or ‘resummation’, of the expansions arising from the multiscale computation of the generating function. The goal will be to show that, by appropriately choosing the parameters t_1^*, β, Z , which the right side of (1.7) depends on (and which are related via a simple invertible mapping to the initial values of the running coupling constants, ν_0, ζ_0, η_0), the whole sequence $\underline{\nu}$ remains bounded, uniformly in h^* ; see Sect. 4.5. Under these conditions, we will be able to show that the resulting expansions for multipoint energy correlations are convergent, uniformly in h^* . Our estimates are based on writing the quantities involved as sums over terms indexed by Gallavotti-Nicolò (GN) trees [13–15], which emerge naturally from the multiscale procedure; the relevant aspects of the definitions of the GN trees will be reviewed in Sect. 4.3 below.

In order to obtain L, M independent values of these parameters, we study the iteration in the limit $L, M \rightarrow \infty$ in Sect. 4; we can also restrict to $\mathbf{A} = \mathbf{0}$, since this already includes all of the potentially divergent terms. This would superficially appear to involve a number of complications such as defining an infinite-dimensional Grassmann integral, but in fact the multiscale computation of the generating function, when understood as an iteration for the kernels of $\mathcal{V}^{(h)}$, denoted by $V_\Lambda^{(h)}$, has a perfectly straightforward infinite-volume version, which is stated and analyzed in Sect. 4. The convergence as $L, M \rightarrow \infty$ of the finite volume kernels $V_\Lambda^{(h)}$ to the solution $V_\infty^{(h)}$ of the infinite-volume recursive equations for the kernels is one of the main subjects of [4], especially [4, Section 3].

Section 4 is a reformulation of [18, Section 3]. We nonetheless present it at length, partly because the treatment of the propagator on the cylinder in Sect. 2 imposes a different choice of variables which makes the translation of some statements awkward, but mainly in order to take the opportunity to

make a number of technical improvements and simplify some unnecessarily obscure aspects of what is already a complicated argument.

Previously, e.g., in [3, 15, 16, 18], the localization operator (and consequently the remainder) was defined in terms of the Fourier transform of the functions involved. This has the advantage of providing a simple procedure for parametrizing the local part of the effective potential by a finite number of running coupling constants, but is quite difficult to apply to non-translation-invariant systems (in [3] this led to a peculiar restriction on the dependence of the interaction on the system size). Moreover, it makes the treatment of finite size corrections awkward and leads to a convoluted definition of the derivative operators in the remainder (see [6] and [19] for the treatment of finite-size corrections on a finite torus via the ‘standard’ definition of localization operator). To deal with this, in Sect. 4.2 we introduce a localization operator defined directly in terms of lattice functions and write the remainder in terms of discrete derivatives using a lattice interpolation procedure. Such definitions naturally admit finite volume counterparts, discussed in [4, Section 3.1].

The strategy used to estimate the interpolation factors in the above cited works also involves decomposing them into components which can be matched with propagators; this involves a number of complications, since it cannot be done in a strictly iterative fashion (this is the problem discussed in [6, Section 3.3]). When we handle this issue in Sect. 4.4, see in particular Proposition 4.6, we instead show iteratively that the coefficient function of the effective interaction satisfies a norm bound morally equivalent to exponential decay in the position variables (associated with exponential decay in the scale-decomposed propagators, see Proposition 2.3); this then makes it possible to bound the contribution of the interpolation operator immediately, avoiding technical issues such as the ‘accumulation of derivatives’ (see [6, Section 3.3] and [15, end of Section 8.4]) or the proof that the Jacobian associated with the change of variables arising from the interpolation procedure is equal to 1 (see [6, (3.119)]). In these aspects, the strategy used in this paper for iteratively estimating the kernels of the effective potentials overlaps with [20], which was developed in parallel with the present work.

2. The Nearest-Neighbor Model

In this section we review some aspects of the exact solution of the nearest-neighbor model ($\lambda = 0$), which will play a central role in the multiscale computation of the generating function for the energy correlations of the non-integrable model, to be discussed in the following sections. In particular, after having recalled the Grassmann representation of the partition function, we explain how to diagonalize the Grassmann action; next, we compute the Grassmann propagator and define its multiscale decomposition, to be used in the following; finally, we compute the scaling limit of the propagator, with quantitative bounds on the remainder.

2.1. Diagonalization of the Free Action

2.1.1. Introduction to the Grassmann Variables and Representation. Let us recall the form of the Hamiltonian $H_\Lambda(\sigma)$ (1.1) in the integrable case $\lambda = 0$:

$$H_\Lambda(\sigma) = - \sum_{l=1}^2 J_l \sum_{z \in \Lambda} \sigma_z \sigma_{z+\hat{e}_l},$$

with the understanding that $\sigma_{z+\hat{e}_2} = 0$ for $z = (z_1, z_2)$ such that³ $z_2 = M$ and $\sigma_{z+\hat{e}_1} = \sigma_{z+(1-L)\hat{e}_1}$ for z such that $z_1 = L$.

As is well known [26, Chapter VI.3], the partition function at inverse temperature $\beta > 0$, $Z_\Lambda = \sum_{\sigma \in \Omega_\Lambda} e^{-\beta H_\Lambda(\sigma)}$, can be written as a Pfaffian, which admits the following representation in terms of Grassmann variables, see, e.g., [27] or [16, Appendix A1]:

$$Z_\Lambda = 2^{LM} (\cosh \beta J)^{L(2M-1)} \int \mathcal{D}\Phi e^{\mathcal{S}_{t_1, t_2}(\Phi)}, \tag{2.1.1}$$

where $\Phi = \{(\bar{H}_z, H_z, \bar{V}_z, V_z)\}_{z \in \Lambda}$ is a collection of $4LM$ Grassmann variables (we will also use the notation $\{\Phi_i\}_{i \in \mathcal{I}}$ for \mathcal{I} a suitable label set with $4LM$ elements), $\mathcal{D}\Phi$ denotes the Grassmann ‘differential’,

$$\mathcal{D}\Phi = \prod_{z \in \Lambda} d\bar{H}_z dH_z d\bar{V}_z dV_z,$$

and

$$\begin{aligned} \mathcal{S}_{t_1, t_2}(\Phi) := & \sum_{z \in \Lambda} (t_1 \bar{H}_z H_{z+\hat{e}_1} + t_2 \bar{V}_z V_{z+\hat{e}_2} + \bar{H}_z H_z + \bar{V}_z V_z + \bar{V}_z \bar{H}_z \\ & + V_z \bar{H}_z + H_z \bar{V}_z + V_z H_z) \end{aligned} \tag{2.1.2}$$

where $t_l = \tanh \beta J_l$ for $l = 1, 2$, and $H_{(L+1, (z)_2)}$, $V_{((z)_1, M+1)}$ should be interpreted as representing $-H_{(1, (z)_2)}$ and 0 , respectively. The identification of $H_{(L+1, (z)_2)}$ with $-H_{(1, (z)_2)}$ corresponds to *anti-periodic* boundary conditions in the horizontal direction for the Grassmann variables: these are the right boundary conditions to consider for a cylinder with an even number of sites in the periodic direction, see [26, Eq.(2.6d)]. For later reference, we let $E_x = \bar{H}_z H_{z+\hat{e}_1}$ for a horizontal edge x with endpoints $z, z + \hat{e}_1$, and $E_x = \bar{V}_z V_{z+\hat{e}_2}$ for a vertical edge x with endpoints $z, z + \hat{e}_2$. Sometimes, we will call $\{(\bar{H}_z, H_z)\}_{z \in \Lambda}$ the *horizontal* variables, and $\{(\bar{V}_z, V_z)\}_{z \in \Lambda}$ the *vertical* ones.

The quadratic form $\mathcal{S}_{t_1, t_2}(\Phi)$ can be written as $\mathcal{S}_{t_1, t_2}(\Phi) = \frac{1}{2}(\Phi, A\Phi)$ for a suitable $4LM \times 4LM$ anti-symmetric matrix A (here (\cdot, \cdot) indicates the standard scalar product for vectors whose components are labeled by indices in \mathcal{I} , i.e., $(\Phi, A\Phi) = \sum_{i, j \in \mathcal{I}} \Phi_i A_{ij} \Phi_j$). In terms of this matrix A , (2.1.1) can be rewritten as

$$Z_\Lambda = 2^{LM} (\cosh \beta J_1)^{LM} (\cosh \beta J_2)^{L(M-1)} \text{Pf} A.$$

³In this section we shall denote the components of $z \in \Lambda$ by z_1, z_2 . We warn the reader that in the following sections the symbols z_1, z_2 will mostly be used, instead, for the first two elements of an n -ple z in Λ^n or in $(\mathbb{Z}^2)^n$, for which we will use the notation $z = (z_1, \dots, z_n)$.

[We recall that the Pfaffian of a $2n \times 2n$ antisymmetric matrix A is defined as

$$\text{Pf}A := \frac{1}{2^n n!} \sum_{\pi} (-1)^{\pi} A_{\pi(1),\pi(2)} \cdots A_{\pi(2n-1),\pi(2n)}; \tag{2.1.3}$$

where the sum is over permutations π of $(1, \dots, 2n)$, with $(-1)^{\pi}$ denoting the signature. One of the properties of the Pfaffian is that $(\text{Pf}A)^2 = \det A$.] For later purpose, we also need to compute the averages of arbitrary monomials in the Grassmann variables Φ_i , with $i \in \mathcal{I}$. These can all be reduced to the computation of the inverse of A , thanks to the ‘fermionic Wick rule’:

$$\langle \Phi_{i_1} \cdots \Phi_{i_m} \rangle := \frac{1}{\text{Pf}A} \int \mathcal{D}\Phi \Phi_{i_1} \cdots \Phi_{i_m} e^{\frac{1}{2}(\Phi, A\Phi)} = \text{Pf}G, \tag{2.1.4}$$

where, if m is even, G is the $m \times m$ matrix with entries

$$G_{jk} = \langle \Phi_{i_j} \Phi_{i_k} \rangle = -[A^{-1}]_{i_j, i_k} \tag{2.1.5}$$

(if m is odd, the r.h.s. of (2.1.4) should be interpreted as 0). Often $\langle \Phi_i \Phi_j \rangle$ is referred to as the (ij) component of the) *propagator* of the Grassmann field Φ , or as the *covariance* of $\mathcal{D}\Phi e^{\mathcal{S}(\Phi)}$; such a form (with a quadratic function in the exponent) is known as a *Grassmann Gaussian measure*.

In the following sections, we shall explain how to compute the Pfaffian of A and its inverse A^{-1} , via a block diagonalization procedure.

2.1.2. Diagonalization of \mathcal{S}_{t_1, t_2} .

Horizontal direction diagonalization and Schur reduction. By exploiting the periodic boundary conditions in the horizontal direction, we can block diagonalize the Grassmann action by performing a Fourier transform in the same direction: for each $z_2 \in \{1, 2, \dots, M\}$ we define

$$\begin{aligned} H_{z_2}(k_1) &= \sum_{z_1=1}^L e^{ik_1 z_1} H_{(z_1, z_2)}, & \bar{H}_{z_2}(k_1) &= \sum_{z_1=1}^L e^{ik_1 z_1} \bar{H}_{(z_1, z_2)}, \\ V_{z_2}(k_1) &= \sum_{z_1=1}^L e^{ik_1 z_1} V_{(z_1, z_2)}, & \bar{V}_{z_2}(k_1) &= \sum_{z_1=1}^L e^{ik_1 z_1} \bar{V}_{(z_1, z_2)}, \end{aligned} \tag{2.1.6}$$

with $k_1 \in \mathcal{D}_L$, where

$$\mathcal{D}_L := \left\{ \frac{\pi(2m-1)}{L} : m = -\frac{L}{2} + 1, \dots, \frac{L}{2} \right\}, \tag{2.1.7}$$

in terms of which

$$\begin{aligned} \mathcal{S}_{t_1, t_2}(\Phi) &= \frac{1}{L} \sum_{\substack{k_1 \in \mathcal{D}_L \\ z_2=1, \dots, M}} \left[(1 + t_1 e^{-ik_1}) \bar{H}_{z_2}(-k_1) H_{z_2}(k_1) + \bar{V}_{z_2}(-k_1) V_{z_2}(k_1) \right. \\ &\quad + t_2 \bar{V}_{z_2}(-k_1) V_{z_2+1}(k_1) + \bar{V}_{z_2}(-k_1) \bar{H}_{z_2}(k_1) + V_{z_2}(-k_1) \bar{H}_{z_2}(k_1) \\ &\quad \left. + H_{z_2}(-k_1) \bar{V}_{z_2}(k_1) + V_{z_2}(-k_1) H_{z_2}(k_1) \right]. \end{aligned} \tag{2.1.8}$$

[Note that as a consequence of the convention that $V_{(z_1, M+1)} = 0$, $V_{M+1}(k_1)$ should also be interpreted as 0.] The terms in the second line of (2.1.8), mixing

the horizontal with the vertical variables, can be eliminated by a linear transformation corresponding to a Schur reduction of the coefficient matrix (cf. [26, p. 120]):

$$\begin{aligned} \begin{bmatrix} \overline{H}_{z_2}(k_1) \\ H_{z_2}(k_1) \end{bmatrix} &= \begin{bmatrix} \xi_{+,z_2}(k_1) \\ \xi_{-,z_2}(k_1) \end{bmatrix} + \begin{bmatrix} (1+t_1 e^{ik_1})^{-1} & -(1+t_1 e^{ik_1})^{-1} \\ (1+t_1 e^{-ik_1})^{-1} & (1+t_1 e^{-ik_1})^{-1} \end{bmatrix} \begin{bmatrix} \phi_{+,z_2}(k_1) \\ \phi_{-,z_2}(k_1) \end{bmatrix}, \\ \begin{bmatrix} \overline{V}_{z_2}(k_1) \\ V_{z_2}(k_1) \end{bmatrix} &= \begin{bmatrix} \phi_{+,z_2}(k_1) \\ \phi_{-,z_2}(k_1) \end{bmatrix}; \end{aligned} \tag{2.1.9}$$

Defining a related set of Grassmann variables on Λ by $\phi_{\omega,z} = \frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} e^{-ik_1 z_1} \phi_{\omega,z_2}(k_1)$ and analogously for $\xi_{\omega,z}$, we then obtain

$$\begin{aligned} \begin{bmatrix} \overline{H}_z \\ H_z \end{bmatrix} &= \begin{bmatrix} \xi_{+,z} \\ \xi_{-,z} \end{bmatrix} + \sum_{y=1}^L \begin{bmatrix} s_+(z_1-y) & -s_+(z_1-y) \\ s_-(z_1-y) & s_-(z_1-y) \end{bmatrix} \begin{bmatrix} \phi_{+, (y,z_2)} \\ \phi_{-, (y,z_2)} \end{bmatrix}, \\ \begin{bmatrix} \overline{V}_z \\ V_z \end{bmatrix} &= \begin{bmatrix} \phi_{+,z} \\ \phi_{-,z} \end{bmatrix}, \end{aligned} \tag{2.1.10}$$

where $s_{\pm}(z_1) := \frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} \frac{e^{-ik_1 z_1}}{1+t_1 e^{\pm ik_1}}$. By the Poisson summation formula,

$$s_{\pm}(z_1) = \sum_{n \in \mathbb{Z}} (-1)^n s_{\infty, \pm}(z_1 + nL), \quad \text{with} \quad s_{\infty, \pm}(z) = \int_{-\pi}^{\pi} \frac{dk_1}{2\pi} \frac{e^{-ik_1 z_1}}{1+t_1 e^{\pm ik_1}}. \tag{2.1.11}$$

It is straightforward to check, via a complex shift of the path of integration over k_1 , that $s_{\infty, \pm}$ (and, therefore, s_{\pm}) decays exponentially in z_1 ; more precisely, $|s_{\infty, \pm}(z_1)| \leq e^{-\alpha|z_1|}(1-t_1 e^{\alpha})^{-1}$ for any $\alpha \in [0, -\log t_1]$, and $|s_{\pm}(z_1) - s_{\infty, \pm}(z_1)| \leq e^{-\alpha L}(1-t_1 e^{\alpha})^{-1}$ whenever $|z_1| \leq L/2$.

In terms of the new variables, the Grassmann action reads $\mathcal{S}_{t_1, t_2}(\Phi) = \mathcal{S}_m(\xi) + \mathcal{S}_c(\phi)$ (the labels ‘ m ’ and ‘ c ’ stand for ‘massive’ and ‘critical’, for reasons that will become clear soon), where

$$\mathcal{S}_m(\xi) = \frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} \sum_{z_2=1}^M (1+t_1 e^{-ik_1}) \xi_{+,z_2}(-k_1) \xi_{-,z_2}(k_1), \tag{2.1.12}$$

$$\begin{aligned} \mathcal{S}_c(\phi) &= \frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} \sum_{z_2=1}^M \left[-b(k_1) \phi_{+,z_2}(-k_1) \phi_{-,z_2}(k_1) + t_2 \phi_{+,z_2}(-k_1) \phi_{-,z_2+1}(k_1) \right. \\ &\quad \left. - \frac{i}{2} \Delta(k_1) \phi_{+,z_2}(-k_1) \phi_{+,z_2}(k_1) + \frac{i}{2} \Delta(k_1) \phi_{-,z_2}(-k_1) \phi_{-,z_2}(k_1) \right], \end{aligned} \tag{2.1.13}$$

with

$$\Delta(k_1) := \frac{2t_1 \sin k_1}{|1+t_1 e^{ik_1}|^2}, \quad b(k_1) := \frac{1-t_1^2}{|1+t_1 e^{ik_1}|^2}, \tag{2.1.14}$$

where, as a consequence of the convention used above for V_z , the term in $\mathcal{S}_c(\phi)$ involving $\phi_{M+1,-}(k_1)$ should be interpreted as being equal to zero. Since \mathcal{S}_m and \mathcal{S}_c involve independent sets of Grassmann variables, the Gaussian integral

appearing in the partition function factors into a product of two integrals, and the propagators associated with the two terms can be calculated separately.

The ‘massive’ propagator The calculations for \mathcal{S}_m are trivial. Let the anti-symmetric matrix A_m be defined by $\mathcal{S}_m(\xi) = \frac{1}{2}(\xi, A_m \xi)$. Recall that \mathcal{S}_m was defined in (2.1.12), from which

$$\text{Pf}A_m = \prod_{k_1 \in \mathcal{D}_L} (1 + t_1 e^{ik_1})^M,$$

and the propagator is given by the appropriate entry of A_m^{-1} , which in the form used in (2.1.12) is block-diagonal with 2×2 blocks such that

$$\langle \xi_{\omega, z_2}(k_1) \xi_{\omega', z'_2}(k'_1) \rangle = L \delta_{z_2, z'_2} \delta_{\omega, -\omega'} \delta_{k_1, -k'_1} \frac{\omega}{1 + t_1 e^{i\omega k_1}}.$$

Therefore, going back to x -space,

$$\langle \xi_{\omega, z} \xi_{\omega', z'} \rangle = \omega \delta_{\omega, -\omega'} s_{\omega}(z_1 - z'_1) \delta_{z_2, z'_2}, \tag{2.1.15}$$

where $s_{\omega}(z_1)$ was defined right after (2.1.10). For later reference, the matrix formed by the elements in (2.1.15) will be referred to as the massive propagator and denoted by

$$\mathfrak{g}_m(z, z') = \delta_{z_2, z'_2} \begin{bmatrix} 0 & s_+(z_1 - z'_1) \\ -s_-(z_1 - z'_1) & 0 \end{bmatrix}. \tag{2.1.16}$$

Recalling that s_{\pm} decays exponentially, see comments after (2.1.11), we see that $\mathfrak{g}_m(z, z')$ decays exponentially as well, and so corresponds to a massive field in the language of quantum field theory.

The ‘critical’ propagator The antisymmetric matrix A_c defined by $\mathcal{S}_c(\phi) = \frac{1}{2}(\phi, A_c \phi)$ can be placed into an explicit block-diagonal form by an ansatz resembling a Fourier sine transform with shifted frequencies; this involves a lengthy but elementary calculation which is detailed in Appendix A. Here we simply state the result for the critical case

$$t_1 t_2 + t_1 + t_2 = 1 \Leftrightarrow t_2 = \frac{1 - t_1}{1 + t_1} \Leftrightarrow t_1 = \frac{1 - t_2}{1 + t_2}, \tag{2.1.17}$$

which is the only case of relevance for the present work⁴. In this case we obtain

$$\begin{bmatrix} \langle \phi_{+, z} \phi_{+, z'} \rangle & \langle \phi_{+, z} \phi_{-, z'} \rangle \\ \langle \phi_{-, z} \phi_{+, z'} \rangle & \langle \phi_{-, z} \phi_{-, z'} \rangle \end{bmatrix} = \begin{bmatrix} g_{++}(z, z') & g_{+-}(z, z') \\ g_{-+}(z, z') & g_{--}(z, z') \end{bmatrix} = \mathfrak{g}_c(z, z') \tag{2.1.18}$$

where

$$\begin{aligned} \mathfrak{g}_c(z, z') := & \frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} \sum_{k_2 \in \mathcal{Q}_M(k_1)} \frac{1}{2N_M(k_1, k_2)} e^{-ik_1(z_1 - z'_1)} \\ & \times \left\{ e^{-ik_2(z_2 - z'_2)} \hat{\mathfrak{g}}(k_1, k_2) - e^{-ik_2(z_2 + z'_2)} \begin{bmatrix} \hat{g}_{++}(k_1, k_2) & \hat{g}_{+-}(k_1, -k_2) \\ \hat{g}_{-+}(k_1, k_2) & e^{2ik_2(M+1)} \hat{g}_{--}(k_1, k_2) \end{bmatrix} \right\} \end{aligned} \tag{2.1.19}$$

⁴The case $t_2 < (1 - t_1)/(1 + t_1)$, corresponding to the paramagnetic phase can be considered in the same way, while the case $t_2 > (1 - t_1)/(1 + t_1)$ is somewhat more complicated since one of the frequencies may be complex [23].

with

$$\begin{aligned} \hat{g}(k_1, k_2) &:= \begin{bmatrix} \hat{g}_{++}(k_1, k_2) & \hat{g}_{+-}(k_1, k_2) \\ \hat{g}_{-+}(k_1, k_2) & \hat{g}_{--}(k_1, k_2) \end{bmatrix} \\ &:= \frac{1}{D(k_1, k_2)} \begin{bmatrix} -2it_1 \sin k_1 & -(1 - t_1^2)(1 - B(k_1)e^{-ik_2}) \\ (1 - t_1^2)(1 - B(k_1)e^{ik_2}) & +2it_1 \sin k_1 \end{bmatrix} \end{aligned} \tag{2.1.20}$$

where

$$D(k_1, k_2) := 2(1 - t_2)^2(1 - \cos k_1) + 2(1 - t_1)^2(1 - \cos k_2), \tag{2.1.21}$$

$$B(k_1) := t_2 \frac{|1 + t_1 e^{ik_1}|^2}{1 - t_1^2}, \tag{2.1.22}$$

$\mathcal{Q}_M(k_1)$ is the set of solutions of the following equation, thought of as an equation for k_2 at k_1 fixed, in the interval $(-\pi, \pi)$:

$$\sin k_2(M + 1) = B(k_1) \sin k_2 M, \tag{2.1.23}$$

and

$$N_M(k_1, k_2) = \frac{\frac{d}{dk_2} (B(k_1) \sin k_2 M - \sin k_2(M + 1))}{B(k_1) \cos k_2 M - \cos k_2(M + 1)}. \tag{2.1.24}$$

Remark 2.1. From the above formula it is immediately clear that

$$\hat{g}_{++}(k_1, k_2) = \hat{g}_{++}(k_1, -k_2) = -\hat{g}_{++}(-k_1, k_2) = \hat{g}_{--}(-k_1, k_2) \tag{2.1.25}$$

and

$$\hat{g}_{+-}(k_1, k_2) = \hat{g}_{+-}(-k_1, k_2) = -\hat{g}_{-+}(k_1, -k_2); \tag{2.1.26}$$

furthermore, Eq. (2.1.23) is equivalent to

$$\hat{g}_{+-}(k_1, k_2) = -e^{-2ik_2(M+1)} \hat{g}_{-+}(k_1, k_2), \tag{2.1.27}$$

which therefore holds for all $q \in \mathcal{Q}_M(k_1)$. Moreover, $N_M(k_1, k_2) = N_M(-k_1, k_2) = N_M(k_1, -k_2)$. As we will see in Sect. 2.4 below, these relationships are closely related to the symmetries of the Ising model on a cylindrical lattice.

Remark 2.2. The definition (2.1.19) can be extended to all $z, z' \in \mathbb{R}^2$; then in particular, using the relationships listed in the previous remark, we have

$$g_{++}((z_1, 0), z') = g_{++}(z, (z'_1, 0)) = g_{+-}((z_1, 0), z') = g_{-+}(z, (z'_1, 0)) = 0, \tag{2.1.28}$$

and

$$\begin{aligned} g_{+-}(z, (z'_1, M + 1)) &= g_{-+}((z_1, M + 1), z') = g_{--}((z_1, M + 1), z') \\ &= g_{--}(z, (z'_1, M + 1)) = 0, \end{aligned} \tag{2.1.29}$$

for all z, z' .

2.2. The Critical Propagator: Multiscale Decomposition and Decay Bounds

In this section, we decompose $\mathbf{g}_c(z, z')$ into a sum of terms satisfying bounds which are the main inputs of the multiscale expansion.

2.2.1. Multiscale and Bulk/Edge Decompositions. Let

$$f_\eta(k_1, k_2) := D(k_1, k_2)e^{-\eta D(k_1, k_2)}, \tag{2.2.1}$$

with $D(k_1, k_2)$ as in (2.1.21). Note that

$$\int_0^\infty f_\eta(k_1, k_2) \, d\eta = 1 \tag{2.2.2}$$

as long as k_1 and k_2 are not both integer multiples of 2π (and so whenever $k_2 \in \mathcal{Q}_M(k_1)$). Comparing with Eq. (2.1.20), we see immediately that $f_\eta(k_1, k_2)\hat{\mathbf{g}}(k_1, k_2)$ is an entire function of both k_1 and k_2 . The reader may find it helpful in what follows to bear in mind that, for large η , $f_\eta(k_1, k_2)$ is peaked in a region where $D(k_1, k_2)$ is of the order η^{-1} , and so k_1, k_2 are of order $\eta^{-1/2}$.

Thanks to (2.2.2), f_η induces the following multiscale decomposition of \mathbf{g}_c defined in (2.1.19). Let $h^* := -\lfloor \log_2(\min\{L, M\}) \rfloor$; then, for any $h^* \leq h < 0$,

$$\mathbf{g}_c(z, z') = \mathbf{g}^{(\leq h)}(z, z') + \sum_{j=h+1}^0 \mathbf{g}^{(j)}(z, z'), \tag{2.2.3}$$

where

$$\mathbf{g}^{(0)}(z, z') := \int_0^1 \mathbf{g}^{[\eta]}(z, z') \, d\eta, \tag{2.2.4}$$

$$\mathbf{g}^{(h)}(z, z') := \int_{2^{-2h-2}}^{2^{-2h}} \mathbf{g}^{[\eta]}(z, z') \, d\eta \quad \text{for } h^* < h < 0, \tag{2.2.5}$$

$$\mathbf{g}^{(\leq h)}(z, z') := \int_{2^{-2h-2}}^\infty \mathbf{g}^{[\eta]}(z, z') \, d\eta, \tag{2.2.6}$$

and

$$\begin{aligned} \mathbf{g}^{[\eta]}(z, z') &:= \frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} \sum_{k_2 \in \mathcal{Q}_M(k_1)} \frac{1}{2N_M(k_1, k_2)} e^{-ik_1(z_1 - z'_1)} f_\eta(k_1, k_2) \\ &\times \left\{ e^{-ik_2(z_2 - z'_2)} \hat{\mathbf{g}}(k_1, k_2) - e^{-ik_2(z_2 + z'_2)} \begin{bmatrix} \hat{g}_{++}(k_1, k_2) & \hat{g}_{+-}(k_1, -k_2) \\ \hat{g}_{-+}(k_1, k_2) & e^{2ik_2(M+1)} \hat{g}_{--}(k_1, k_2) \end{bmatrix} \right\}. \end{aligned} \tag{2.2.7}$$

Note that the single-scale propagators preserve the cancellations at the boundary spelled out in Remark 2.2 above, namely, denoting the components of $\mathbf{g}^{(h)}$ by $g_{\omega\omega'}^{(h)}$, with $\omega, \omega' \in \{\pm\}$, in analogy with (2.1.18),

$$\begin{aligned} g_{++}^{(h)}((z_1, 0), z') &= g_{++}^{(h)}(z, (z'_1, 0)) = g_{+-}^{(h)}((z_1, 0), z') = g_{-+}^{(h)}(z, (z'_1, 0)) = 0, \\ g_{+-}^{(h)}(z, (z'_1, M+1)) &= g_{-+}^{(h)}((z_1, M+1), z') = g_{--}^{(h)}((z_1, M+1), z') \\ &= g_{--}^{(h)}(z, (z'_1, M+1)) = 0, \end{aligned} \tag{2.2.8}$$

and analogously for $g_{\omega\omega'}^{(\leq h)}$. Note also that, taking $L, M \rightarrow \infty$, the cutoff propagator $\mathfrak{g}_c^{[\eta]}(z, z')$ tends to its infinite-plane counterpart, provided that z, z' are chosen ‘well inside the cylinder’; in particular, if $z_{L,M} := (L/2, \lfloor M/2 \rfloor)$, then

$$\begin{aligned} & \lim_{L,M \rightarrow \infty} \mathfrak{g}^{[\eta]}(z_{L,M} + z, z_{L,M} + z') \\ &= \int_{[-\pi, \pi]^2} \frac{dk_1 dk_2}{(2\pi)^2} e^{-i[k_1(z_1 - z'_1) + k_2(z_2 - z'_2)]} f_\eta(k_1, k_2) \hat{\mathfrak{g}}(k_1, k_2) \quad (2.2.9) \\ &=: \mathfrak{g}_\infty^{[\eta]}(z - z'). \end{aligned}$$

For later purposes, we need to decompose the cutoff propagator $\mathfrak{g}^{[\eta]}$ into a ‘bulk’ part which is minimally sensitive to the size and shape of the cylinder, plus a remainder which we call the ‘edge’ part. The bulk part is simply chosen to be the restriction of $\mathfrak{g}_\infty^{[\eta]}$ to the cylinder, with the appropriate (anti-periodic) boundary conditions in the horizontal direction:

$$\mathfrak{g}_B^{[\eta]}(z, z') := s_L(z_1 - z'_1) \mathfrak{g}_\infty^{[\eta]}(\text{per}_L(z_1 - z'_1), z_2 - z'_2) \quad (2.2.10)$$

where, recalling that $z_1, z'_1 \in \{1, \dots, L\}$,

$$s_L(z_1 - z'_1) := \begin{cases} +1, & |z_1 - z'_1| < L/2 \\ 0, & |z_1 - z'_1| = L/2 \\ -1, & |z_1 - z'_1| > L/2 \end{cases} \quad (2.2.11)$$

and

$$\text{per}_L(z_1) := z_1 - L \left\lfloor \frac{z_1}{L} + \frac{1}{2} \right\rfloor. \quad (2.2.12)$$

The edge part is, by definition, the difference between the full cutoff propagator and its bulk part:

$$\mathfrak{g}_E^{[\eta]}(z, z') := \mathfrak{g}^{[\eta]}(z, z') - \mathfrak{g}_B^{[\eta]}(z, z'). \quad (2.2.13)$$

Using these expressions, we define $\mathfrak{g}_B^{(h)}$ and $\mathfrak{g}_E^{(h)}$ via the analogues of (2.2.4)–(2.2.5), with the subscript c replaced by B and E , respectively. As a consequence, for any $h^* \leq h < 0$,

$$\mathfrak{g}_c(z, z') = \mathfrak{g}^{(\leq h)}(z, z') + \sum_{j=h+1}^0 (\mathfrak{g}_B^{(j)}(z, z') + \mathfrak{g}_E^{(j)}(z, z')). \quad (2.2.14)$$

As already observed in Remark 2.2, all the functions involved in this identity can be naturally extended to all $x, y \in \mathbb{R}^2$ (and, therefore, in particular, to all $z, z' \in \mathbb{Z}^2$), by interpreting the right side of (2.2.7), etc., as a function on $\mathbb{R}^2 \times \mathbb{R}^2$.

2.2.2. Decay Bounds and Gram Decomposition: Statement of the Main Results. Given the multiscale and bulk–edge decomposition (2.2.14), we now intend to prove suitable decay bounds for the single-scale bulk and edge propagators, as well as to show the existence of an inner product representation (‘Gram representation’) thereof. These will be of crucial importance in the non-perturbative multiscale bounds on the partition and generating functions,

discussed in the rest of this work, and they are summarized in the following proposition.

Given a function $f : (\mathbb{Z}^2)^2 \rightarrow \mathbb{C}$ (or a function $f : \Lambda^2 \rightarrow \mathbb{C}$ extendable to $(\mathbb{Z}^2)^2$, in the sense explained after (2.2.14)), we let $\partial_{1,j}$ be the discrete derivative in direction j with respect to the first argument, defined by $\partial_{1,j}f(z, z') := f(z + \hat{e}_j, z') - f(z, z')$, with \hat{e}_j the j -th Euclidean basis vector; an analogous definition holds for $\partial_{2,j}$.

Proposition 2.3. *There exist constants c, C such that, for any integer h in $[h^* + 1, 0]$, any $\mathbf{r} = (r_{1,1}, r_{1,2}, r_{2,1}, r_{2,2}) \in \mathbb{Z}_+^4$, and any $z, z' \in \Lambda^2$,*

1.

$$\|\partial^{\mathbf{r}} \mathbf{g}^{(h)}(z, z')\| \leq C^{1+|\mathbf{r}|_1} \mathbf{r}! 2^{(1+|\mathbf{r}|_1)h} e^{-c2^h \|z-z'\|_1} \tag{2.2.15}$$

where the matrix norm in the left side (recall that $\partial^{\mathbf{r}} \mathbf{g}^{(h)}(x, y)$ is a 2×2 matrix) is the max norm, i.e., the maximum over the matrix elements,

$$\partial^{\mathbf{r}} := \prod_{i,j=1}^2 \partial_{i,j}^{r_{i,j}}, \quad \mathbf{r}! = \prod_{i,j=1}^2 r_{i,j}!, \quad \text{and } \|z\|_1 = |\text{per}_L(z_1)| + |z_2|, \text{ see (2.2.12),}$$

Moreover, if $z, z' \in \Lambda$ are such that $|\text{per}_L(z_1 - z'_1)| < L/2 - |r_{1,1}| - |r_{2,1}|$,

2.

$$\|\partial^{\mathbf{r}} \mathbf{g}_E^{(h)}(z, z')\| \leq C^{1+|\mathbf{r}|_1} \mathbf{r}! 2^{(1+|\mathbf{r}|_1)h} e^{-c2^h d_E(z, z')} \tag{2.2.16}$$

where $d_E(z, z') := \min\{|\text{per}_L(z_1 - z'_1)| + \min\{z_2 + z'_2, 2(M + 1) - z_2 - z'_2\}, L - |\text{per}_L(z_1 - z'_1)| + |z_2 - z'_2|\}$.

Finally, there exists a Hilbert space \mathcal{H}_{LM} with inner product (\cdot, \cdot) including elements $\gamma_{\omega, s, z}^{(h)}, \tilde{\gamma}_{\omega, s, z}^{(h)}, \gamma_{\omega, s, z}^{(\leq h)}, \tilde{\gamma}_{\omega, s, z}^{(\leq h)}$ (for $\mathbf{s} = (s_1, s_2) \in \mathbb{Z}_+^2, x \in \Lambda$) such that whenever $h^* \leq h \leq 0$,

3. $\partial^{(\mathbf{s}, \mathbf{s}')} g_{\omega\omega'}^{(h)}(z, z') \equiv \left(\tilde{\gamma}_{\omega, \mathbf{s}, z}^{(h)}, \gamma_{\omega', \mathbf{s}', z'}^{(h)} \right)$ and

$$\partial^{(\mathbf{s}, \mathbf{s}')} g_{\omega\omega'}^{(\leq h)}(z, z') \equiv \left(\tilde{\gamma}_{\omega, \mathbf{s}, z}^{(\leq h)}, \gamma_{\omega', \mathbf{s}', z'}^{(\leq h)} \right), \text{ and}$$

4. $\left| \gamma_{\omega, \mathbf{s}, z}^{(h)} \right|^2, \left| \tilde{\gamma}_{\omega, \mathbf{s}, z}^{(h)} \right|^2, \left| \gamma_{\omega, \mathbf{s}, z}^{(\leq h)} \right|^2, \left| \tilde{\gamma}_{\omega, \mathbf{s}, z}^{(\leq h)} \right|^2 \leq C^{1+|\mathbf{s}|_1} \mathbf{s}! 2^{h(1+2|\mathbf{s}|_1)}$ where $|\cdot|$ is the norm generated by the inner product (\cdot, \cdot) .

Combining points 3 and 4, we see that

Corollary 2.4. *For all $z, z' \in \Lambda, \mathbf{r} \in \mathbb{Z}_+^4$, and $h^* \leq h \leq 0$,*

$$\|\partial^{\mathbf{r}} \mathbf{g}^{(\leq h)}(z, z')\| \leq C^{1+|\mathbf{r}|_1} \mathbf{r}! 2^{(1+|\mathbf{r}|_1)h}. \tag{2.2.17}$$

Remark 2.5. Since $\mathbf{g}_B^{(h)} = \mathbf{g}^{(h)} - \mathbf{g}_E^{(h)}$, from items 1 and 2 it follows that the bulk, single-scale, propagator $\mathbf{g}_B^{(h)}$ satisfies the same estimate as (2.2.15) for all $x, y \in \Lambda$ allowed in Item 2, i.e., whenever $|\text{per}_L(z_1 - z'_1)| < L/2 - |r_{1,1}| - |r_{2,1}|$. The latter restriction on z, z' just comes from the requirement that the discrete derivatives do not act on the discontinuous functions s_L and per_L entering the definition of $\mathbf{g}_B^{(h)}$ (and, therefore, of $\mathbf{g}_E^{(h)}$); in fact, one can easily check from the proof that, if $\mathbf{r} = \mathbf{0}$, then (2.2.16) is valid for all $z, z' \in \Lambda$, without further restrictions.

Remark 2.6. All the estimates stated in the proposition are uniform in L, M , therefore, they remain valid for the $L, M \rightarrow \infty$ limit of the propagators. In particular, $\mathfrak{g}_\infty^{(h)}(z - z')$ and $\mathfrak{g}_\infty^{(\leq h)}(z - z')$ satisfy the same estimates as (2.2.15) and (2.2.17), respectively. Similarly, the Gram representation stated in items 3 and 4 is also valid for $\mathfrak{g}_\infty^{(h)}$. Therefore, if $|z_1 - z'_1| \leq L/2 - |s_1| - |s'_1|$, also $\partial^{(s,s')} \mathfrak{g}_E^{(h)}(\mathbf{x}, \mathbf{y}) = \partial^{(s,s')} \mathfrak{g}^{(h)}(z, z') - \partial^{(s,s')} \mathfrak{g}_\infty^{(h)}(z - z')$ admits a Gram representation, with qualitatively the same Gram bounds (and, of course, such a representation can be extended by anti-periodicity to all z, z' such that $|\text{per}_L(z_1 - z'_1)| \leq L/2 - |s_1| - |s'_1|$).

Remark 2.7. If we rename the massive propagator in (2.1.16) as $\mathfrak{g}_m(z, z') =: \mathfrak{g}^{(1)}(z, z')$ and then use it to define $\mathfrak{g}_\infty^{(1)}, \mathfrak{g}_B^{(1)}$, and $\mathfrak{g}_E^{(1)}$ in the same way that we did above with $\mathfrak{g}^{[n]}$, it is straightforward to check that the estimates in items 1 and 2 of the proposition remain valid for $h = 1$. Similarly, the reader can check that the proof of items 3 and 4 given below can be straightforwardly applied to the case $h = 1$, as well.

Remark 2.8. Again since $D(k_1, k_2)$ is exactly the denominator in the definition (2.1.20), it is easily seen that $f_\eta(k_1, k_2) \hat{\mathfrak{g}}(k_1, k_2)$ is an entire function of t_1 . All of the bounds in Proposition 2.3 are obtained by writing the relevant quantities as absolutely convergent integrals or sums in k_1, k_2 , and η ; since these bounds are locally uniform in t_1 as long as it is bounded away from 0 and 1, we also see that all of the propagators are analytic functions of t_1 with all other arguments held fixed.

The proof of Proposition 2.3 is in Appendix B.

2.3. Asymptotic Behavior of the Critical Propagator

Although our main result, Theorem 1.1, involves correlation functions for a finite lattice, since the continuum limit of the energy correlation functions of the non-interacting model is well understood [11, 24], as a result we also obtain a characterization of the scaling limit [4, Corollary 1.2]. For completeness, we give here a description of the scaling limit of the critical propagator, from which the non-interacting energy correlation functions are easily calculated.

We rescale the lattice as follows: fix two positive constants ℓ_1, ℓ_2 (no condition on the ratio ℓ_1/ℓ_2), and let $L = 2\lfloor a^{-1}\ell_1/2 \rfloor$, $M = \lfloor a^{-1}\ell_2 \rfloor$ for $a > 0$ the lattice mesh. Let $\Lambda^a := a\Lambda$ and Λ_{ℓ_1, ℓ_2} the continuum cylinder. We also let $\|\cdot\|$ indicate the Euclidean distance on the cylinder Λ_{ℓ_1, ℓ_2} . Given $z, z' \in \Lambda_{\ell_1, \ell_2}$, we let

$$\mathfrak{g}_{c,a}(z, z') = a^{-1} \mathfrak{g}_c(\lfloor a^{-1}z \rfloor, \lfloor a^{-1}z' \rfloor). \tag{2.3.1}$$

The main result of this section concerns the limiting behavior of $\mathfrak{g}_{c,a}$ as $a \rightarrow 0$.

Proposition 2.9. *Given $\ell_1, \ell_2 > 0$, there exist $C, c > 0$ such that for all $x, y \in \Lambda_{\ell_1, \ell_2}$ such that $z \neq z'$ and $a > 0$ for which $a(\min\{\ell_1, \ell_2, \|z - z'\|\})^{-1} \leq c$,*

$$\mathfrak{g}_{c,a}(z, z') = \mathfrak{g}_{\text{scal}}(z, z') + R_a(z, z'), \tag{2.3.2}$$

and $\|R_a(z, z')\| \leq Ca(\min\{\ell_1, \ell_2, \|z - z'\|\})^{-2}$, where

$$\begin{aligned} \mathfrak{g}_{\text{scal}}(z, z') &= \sum_{n \in \mathbb{Z}^2} (-1)^n \left\{ \mathfrak{g}_{\infty}^{\text{scal}}(z_1 - z'_1 + n_1 \ell_1, z_2 - z'_2 + 2n_2 \ell_2) \right. \\ &\quad \left. + \left[-g_1^{\text{scal}}(z_1 - z'_1 + n_1 \ell_1, z_2 + z'_2 + 2n_2 \ell_2) \quad g_2^{\text{scal}}(z_1 - z'_1 + n_1 \ell_1, z_2 + z'_2 + 2n_2 \ell_2) \right] \right\}, \end{aligned} \tag{2.3.3}$$

and where, letting

$$\begin{aligned} g^{\text{scal}}(z_1, z_2) &:= \frac{1}{t_2(1-t_2)} \iint_{\mathbb{R}^2} \frac{dk_1 dk_2}{(2\pi)^2} e^{-ik_1 z_1 - ik_2 z_2} \frac{-ik_1}{k_1^2 + k_2^2} \\ &= -\frac{1}{2\pi t_2(1-t_2)} \frac{z_1}{z_1^2 + z_2^2}, \end{aligned} \tag{2.3.4}$$

we denoted $g_1^{\text{scal}}(z_1, z_2) := g^{\text{scal}}(\frac{z_1}{1-t_2}, \frac{z_2}{1-t_1})$, $g_2^{\text{scal}}(z_1, z_2) := g^{\text{scal}}(\frac{z_2}{1-t_1}, \frac{z_1}{1-t_2})$, and

$$\mathfrak{g}_{\infty}^{\text{scal}}(z_1, z_2) := \begin{bmatrix} g_1^{\text{scal}}(z_1, z_2) & g_2^{\text{scal}}(z_1, z_2) \\ g_2^{\text{scal}}(z_1, z_2) & -g_1^{\text{scal}}(z_1, z_2) \end{bmatrix}. \tag{2.3.5}$$

The proof of Proposition 2.9 is given in Appendix C. It is easy to see from the definition of $\mathfrak{g}_c^{\text{scal}}$ that its entries vanish for $z_2 = 0$ and/or $z_2 = \ell_2$ and/or $z'_2 = 0$ and/or $z'_2 = \ell_2$ in a fashion analogous to the one discussed in Remark 2.2.

2.4. Symmetries of the Propagator

Note that the action $\mathcal{S}_{t_1, t_2}(\Phi)$ of Eq. (2.1.2) is unchanged by the substitutions

$$\overline{H}_z \rightarrow iH_{\theta_1 z}, \quad H_z \rightarrow i\overline{H}_{\theta_1 z}, \quad \overline{V}_z \rightarrow i\overline{V}_{\theta_1 z}, \quad V_z \rightarrow -iV_{\theta_1 z} \tag{2.4.1}$$

with $\theta_1(z_1, z_2) := (L + 1 - z_1, z_2)$, or

$$\overline{H}_z \rightarrow -i\overline{H}_{\theta_2 z}, \quad H_z \rightarrow iH_{\theta_2 z}, \quad \overline{V}_z \rightarrow iV_{\theta_2 z}, \quad V_z \rightarrow i\overline{V}_{\theta_2 z} \tag{2.4.2}$$

where $\theta_2(z_1, z_2) := (z_1, M + 1 - z_2)$. These transformations, of course, correspond to the reflection symmetries of the Ising model on a cylinder. In terms of the ϕ, ξ variables, it is easy to see from Eq. (2.1.10) that these substitutions are equivalent to

$$\phi_{\pm, z} \rightarrow \Theta_1 \phi_{\pm, z} := \pm i \phi_{\pm, \theta_1 z}, \quad \xi_{\pm, z} \rightarrow \Theta_1 \xi_{\pm, z} := i \xi_{\mp, \theta_1 z} \tag{2.4.3}$$

and

$$\phi_{\pm, z} \rightarrow \Theta_2 \phi_{\pm, z} := i \phi_{\mp, \theta_2 z}, \quad \xi_{\pm, z} \rightarrow \Theta_2 \xi_{\pm, z} := \mp i \xi_{\pm, \theta_2 z}. \tag{2.4.4}$$

With a little more notation we can write this more compactly: letting $\phi_{\omega, z}$ denote $\phi_{+,x}, \phi_{-,x}, \xi_{+,x}, \xi_{-,x}$ for $\omega = 1, -1, +i, -i$, respectively, and letting $\theta_j \omega := (-1)^{j+1} \overline{\omega}$ for $\omega \in \mathbb{C}$, Eqs. (2.4.3) and (2.4.4) can be combined into

$$\Theta_j \phi_{\omega, z} := i \alpha_{j, \omega} \phi_{\theta_j \omega, \theta_j z} \tag{2.4.5}$$

where $\alpha_{j, \omega}$ is -1 if $(j = 1 \text{ and } \omega = -1)$ or $(j = 2 \text{ and } \omega = i)$, and 1 otherwise. Since these transformations act on the vector Ψ as orthogonal matrices, this is equivalent to the symmetry of the coefficient matrix A (and therefore

its inverse) under the associated similarity transform, and since \mathfrak{g}_c is just a diagonal block of A^{-1} we have

$$\mathfrak{g}_c(z, z') = \begin{bmatrix} -g_{c; ++} & g_{c; +-} \\ g_{c; -+} & -g_{c; --} \end{bmatrix} (\theta_1 z, \theta_1 z') \tag{2.4.6}$$

and

$$\mathfrak{g}_c(z, z') = - \begin{bmatrix} g_{c; --} & g_{c; -+} \\ g_{c; +-} & g_{c; ++} \end{bmatrix} (\theta_2 z, \theta_2 z'). \tag{2.4.7}$$

These relationships can also be recovered from Eq. (2.1.19), using the observations on $\hat{\mathfrak{g}}_\infty$ in Remark 2.1. This latter point is helpful because, since f_η is even in both k_1 and k_2 , it also applies to $\mathfrak{g}_c^{[\eta]}$. Taking the appropriate $L, M \rightarrow \infty$ limit we also obtain

$$\mathfrak{g}_\infty^{[\eta]}(z_1, z_2) = \begin{bmatrix} -g_{\infty; ++}^{[\eta]} & g_{\infty; +-}^{[\eta]} \\ g_{\infty; -+}^{[\eta]} & -g_{\infty; --}^{[\eta]} \end{bmatrix} (-z_1, z_2) = - \begin{bmatrix} g_{\infty; --}^{[\eta]} & g_{\infty; -+}^{[\eta]} \\ g_{\infty; +-}^{[\eta]} & g_{\infty; ++}^{[\eta]} \end{bmatrix} (z_1, -z_2), \tag{2.4.8}$$

which also implies that $\mathfrak{g}_B^{[\eta]}$ has the symmetries (2.4.6) and (2.4.7). Applying the differences and integrals in the relevant definitions we see that

Lemma 2.10. $\mathfrak{g}_c^{(h)}$, $\mathfrak{g}_c^{(\leq h)}$, $\mathfrak{g}_B^{(h)}$, and $\mathfrak{g}_E^{(h)}$ all have the symmetries (2.4.6) and (2.4.7) for any $h_* \leq h \leq 0$, and $\mathfrak{g}_\infty^{(h)}$ has the symmetries (2.4.8) for any $h \leq 0$.

For $\mathfrak{g}_c^{(1)} \equiv \mathfrak{g}_m$ we have

$$\mathfrak{g}_m(z, z') = - \begin{bmatrix} g_{m; --} & g_{m; -+} \\ g_{m; +-} & g_{m; ++} \end{bmatrix} (\theta_1 z, \theta_1 z') = \begin{bmatrix} -g_{m; ++} & g_{m; +-} \\ g_{m; -+} & -g_{m; --} \end{bmatrix} (\theta_2 z, \theta_2 z') \tag{2.4.9}$$

which similarly extends to $\mathfrak{g}_\infty^{(1)}$, $\mathfrak{g}_B^{(1)}$, and $\mathfrak{g}_E^{(1)}$.

3. Grassmann Representation of the Generating Function

In this section we rewrite the generating function of the energy correlations for the Ising model (1.1) with finite range interactions as an *interacting* Grassmann integral, and we set the stage for the multiscale integration thereof, to be discussed in the following sections. The estimates in this and in the following section are uniform for $J_1/J_2, L/M \in K$ and $t_1, t_2 \in K'$, but may depend upon the choice of K, K' , with K, K' the compact sets introduced before the statement of Theorem 1.1. As anticipated there, we will think of K, K' as being fixed once and for all and, for simplicity, we will not track the dependence upon these sets in the constants $C, C', \dots, c, c', \dots, \kappa, \kappa', \dots$, appearing below. Unless otherwise stated, the values of these constants may change from line to line.

Our goal is to show that the generating function (1.6) of the energy correlations can be replaced, for the purpose of computing multipoint energy

correlations at distinct edges, by the Grassmann generating function (1.7). Our main representation result for the Taylor coefficients at $\mathbf{A} = \mathbf{0}$ of the logarithm of $Z_\Lambda(\mathbf{A})$, analogous to [18, Proposition 1], is the following.

Proposition 3.1. *For any translation invariant interaction V of finite range, there exists $\lambda_0 = \lambda_0(V)$ such that, for any $|\lambda| \leq \lambda_0(V)$,*

$$\frac{\partial}{\partial A_{x_1}} \dots \frac{\partial}{\partial A_{x_n}} \log Z_\Lambda(\mathbf{A}) \Big|_{\mathbf{A}=\mathbf{0}} = \frac{\partial}{\partial A_{x_1}} \dots \frac{\partial}{\partial A_{x_n}} \log \tilde{\Xi}_\Lambda(\mathbf{A}) \Big|_{\mathbf{A}=\mathbf{0}}$$

as long as $n \geq 2$ and the x_j are distinct, where

$$\tilde{\Xi}_\Lambda(\mathbf{A}) := e^{\mathcal{W}(\mathbf{A})} \int \mathcal{D}\Phi e^{\mathcal{S}_{t_1, t_2}(\Phi) + \mathcal{V}(\Phi, \mathbf{A})} \tag{3.1}$$

where \mathcal{S}_{t_1, t_2} was defined in (2.1.2) and, recalling that E_x is the Grassmann binomial defined after (2.1.2):

1.

$$\begin{aligned} \mathcal{V}(\Phi, \mathbf{A}) &= \sum_{x \in \mathfrak{B}_\Lambda} (1 - t_{j(x)}^2) E_x A_x + \sum_{\substack{X, Y \subset \mathfrak{B}_\Lambda \\ X \neq \emptyset}} W_\Lambda^{\text{int}}(X, Y) \prod_{x \in X} E_x \prod_{x \in Y} A_x \\ &\equiv \mathcal{B}^{\text{free}}(\Phi, \mathbf{A}) + \mathcal{V}^{\text{int}}(\Phi, \mathbf{A}) \end{aligned} \tag{3.2}$$

where, for any $n \in \mathbb{N}$, $m \in \mathbb{N}_0$, and suitable positive constants C, c, κ ,

$$\sup_{x_0 \in \mathfrak{B}_\Lambda} \sum_{\substack{X, Y \subset \mathfrak{B}_\Lambda: x_0 \ni X \\ |X|=n, |Y|=m}} |W_\Lambda^{\text{int}}(X, Y)| e^{c\delta(X \cup Y)} \leq C^{m+n} |\lambda|^{\max(1, \kappa(m+n))} \tag{3.3}$$

and $\delta(X)$, for $X \subset \mathfrak{B}_\Lambda$, denotes the size of the smallest $Z \supset X$ which is the edge set of a connected subgraph of \mathfrak{G}_Λ .

2.

$$\mathcal{W}(\mathbf{A}) = \sum_{\substack{Y \subset \mathfrak{B}_\Lambda \\ |Y| \geq 2}} w_\Lambda(Y) \prod_{x \in Y} A_x \tag{3.4}$$

where, for any $m \in \mathbb{N}$, and the same C, c, κ as above,

$$\sup_{x_0 \in \mathfrak{B}_\Lambda} \sum_{\substack{Y \subset \mathfrak{B}_\Lambda: \\ x_0 \ni Y, |Y|=m}} |w_\Lambda(Y)| e^{c\delta(Y)} \leq C^m |\lambda|^{\max(1, \kappa m)}. \tag{3.5}$$

3. $W_\Lambda^{\text{int}}, w_\Lambda$, considered as functions of λ, t_1 , and t_2 , can be analytically continued to any complex λ, t_1, t_2 such that $|\lambda| \leq \lambda_0$ and $|t_1|, |t_2| \in K'$, with K' the same compact set introduced before the statement of Theorem 1.1, and the analytic continuations satisfies the same bounds above.

Proof. The proof is basically the same as [18, Proposition 1], so we refer to that for the details. Note that the restriction in [18] to a pair interaction is unimportant, since any even interaction of the form $V(X)\sigma_X$ with $X \subset \Lambda$ can always be written as a product of factors of ϵ_x , in analogy with the rewriting $\sigma_z \sigma'_z = \frac{1}{2}(U_{z, z'} + D_{z, z'})$ discussed after [18, Equation (2.8)]; the only difference

in the current setting is that the ‘strings’ graphically associated with $U_{z,z'}$ and $D_{z,z'}$, see [18, Figure 3], are replaced by other figures, whose specific shape depends on $V(X)$ and that one should use t_1 or t_2 in place of t as appropriate⁵. Note that the set of strings associated with a pair interaction, or the set of more general figures associated with a generic even interaction, is, or can be chosen to be, invariant under the basic symmetries of the model, namely horizontal translations, and horizontal and vertical reflections; therefore, in the following, we shall assume that such a graphical representation is invariant under these symmetries.

By proceeding as in [18] we get the analogue of [18, Eq.(2.20)], namely

$$\mathcal{W}(\mathbf{A}) + \mathcal{V}(\Phi, \mathbf{A}) = \sum_{\Gamma \in C_\Lambda} \varphi^T(\Gamma) \prod_{\gamma \in \Gamma} \zeta_G(\gamma) = \sum_{X, Y \subset \mathfrak{B}_\Lambda} W_\Lambda(X, Y) \prod_{x \in X} E_x \prod_{x \in Y} A_x \tag{3.6}$$

where C_Λ is the set of multipolygons in Λ , φ^T is the Mayer’s coefficient, and ζ_G is the activity of the polygon γ , which is a polynomial in the E_x, A_x for edges x in γ (for more details about the notation and more precise definitions, we refer to [18]). The terms with $X = \emptyset$ contribute to $\mathcal{W}(\mathbf{A})$ (that is, we let $w_\Lambda(Y) := W_\Lambda(\emptyset, Y)$), while those with $X \neq \emptyset$ contribute to $\mathcal{V}(\Phi, \mathbf{A})$ (note that, for the purpose of computing the derivatives of $\log \tilde{\Xi}_\Lambda(\mathbf{A})$ of order 2 or more, the terms with $|X| = 0$ and $|Y| = 0, 1$ can be dropped from the definition of $\mathcal{W}(\mathbf{A})$, and we do so). The explicit computation of the term independent of λ , which has $X = Y \in \mathfrak{B}$, leads to the decomposition in Eq. (3.2). The bounds Eqs. (3.3) and (3.5) follow directly from the bounds in [18], see, e.g., [18, Eq.(2.25)] and following discussion.

Finally, the analyticity property (which was used implicitly in [18]) follows by noting that we have defined all of the quantities of interest as uniformly absolutely convergent sums of terms which are themselves analytic functions of λ, t_1, t_2 as long as the absolute values of these parameters belong to the appropriate intervals. □

With a view toward the analysis of finite size effects in [4] (and, in particular, toward the claims done in [4, Section 2.2] after the statement of [4, Proposition 2.5]), it is convenient to decompose the kernel W_Λ^{int} of $\mathcal{V}^{\text{int}}(\Phi, \mathbf{A})$ into a ‘bulk’ plus an ‘edge’ part. This requires a bit of notation. Note that any subset X of Λ with horizontal diameter smaller than $L/2$ can be identified (non-uniquely, of course) with a subset of \mathbb{Z}^2 with the same diameter and ‘shape’ as X ; we call $X_\infty \subset \mathbb{Z}^2$ one of these arbitrarily chosen representatives⁶

⁵We take the occasion to point out that [18, Figure 4] contains a mistake: the string S_2 depicted there is not allowed by the conventions of [18]: the shape S_2 can only be obtained as the union of two appropriate strings.

⁶For instance, given $X = \{z_1, \dots, z_n\}$, recalling that $(z_i)_1 \in \{1, 2, \dots, L\}$ and $(z_i)_2 \in \{1, 2, \dots, M\}$, we can let $X_\infty = \{y_1, \dots, y_n\}$ be the set of points such that: (1) the vertical coordinates are the same as those of z , i.e., $(y_i)_2 = (z_i)_2, \forall i = 1, \dots, n$; (2) the horizontal coordinate of y_1 is the same as z_1 , i.e., $(y_1)_1 = (z_1)_1$; (3) all the other horizontal coordinates are the same modulo L , i.e., $(y_i)_1 = (z_i)_1 \bmod L, \forall i = 2, \dots, n$; (4) the specific values of $(y_i)_2$

of X , and we shall use an analogous convention for the subsets of \mathfrak{B}_Λ with horizontal diameter smaller than $L/2$.

Lemma 3.2. *Under the same assumptions of Proposition 3.1, the kernel W_Λ^{int} of $\mathcal{V}^{\text{int}}(\Phi, \mathbf{A})$ can be decomposed as*

$$\begin{aligned} W_\Lambda^{\text{int}}(X, Y) &= W_\infty^{\text{int}}(X_\infty, Y_\infty) \mathbb{1}(\text{diam}_1(X \cup Y) \leq L/3) + W_E^{\text{int}}(X, Y) \\ &=: W_B^{\text{int}}(X, Y) + W_E^{\text{int}}(X, Y), \end{aligned} \tag{3.7}$$

where diam_1 is the horizontal diameter on the cylinder Λ ; $X_\infty, Y_\infty \subset \mathfrak{B} := \mathfrak{B}_{\mathbb{Z}^2}$ are two representatives of X, Y , respectively, such that $X_\infty \cup Y_\infty$ is a representative of $X \cup Y$, in the sense defined before the statement of the lemma; W_∞^{int} is a function, independent of L, M , invariant under translations and under reflections about either coordinate axis, which satisfies the same weighted L^1 bound (3.3) as W_Λ^{int} . Moreover, for any $n \in \mathbb{N}$ and $m \in \mathbb{N}_0$, W_E^{int} satisfies

$$\frac{1}{L} \sum_{\substack{X, Y \subset \mathfrak{B}_\Lambda \\ |X|=n, |Y|=m}} |W_E^{\text{int}}(X, Y)| e^{c\delta_E(X \cup Y)} \leq C^{m+n} |\lambda|^{\max(1, \kappa(m+n))}, \tag{3.8}$$

with the same C, c, κ as in Proposition 3.1, where $\delta_E(X)$ is the cardinality of the smallest connected subset of \mathfrak{B}_Λ which includes X and either touches the boundary of the cylinder⁷, or its horizontal diameter is larger than $L/3$.

Proof. In order to obtain the decomposition (3.7), let

$$W_\infty^{\text{int}}(X_\infty, Y_\infty) := \lim_{L, M \rightarrow \infty} W_\Lambda^{\text{int}}(X_\infty + z_{L, M}, Y_\infty + z_{L, M}), \tag{3.9}$$

where $z_{L, M} = (L/2, \lfloor M/2 \rfloor)$ and $X_\infty + z_{L, M}$ is the translate of X_∞ by $z_{L, M}$; note that this limit is well defined thanks to the fact that W_Λ^{int} can be expressed in terms of a sum like [18, Eq.(2.21)], which is exponentially convergent, see [18, Eq.(2.25)]. The kernel W_∞^{int} satisfies the analogue of (3.6), that is

$$\sum_{\Gamma \in C_\infty} \varphi^T(\Gamma) \left(\prod_{\gamma \in \Gamma} \zeta_G(\gamma) - \prod_{\gamma \in \Gamma} \zeta_G(\gamma)|_{\lambda=0} \right) = \sum_{X, Y \subset \mathfrak{B}} W_\infty^{\text{int}}(X, Y) \prod_{x \in X} E_x \prod_{x \in Y} A_x, \tag{3.10}$$

with C_∞ the set of multipolygons on \mathbb{Z}^2 , and the activity $\zeta_G(\gamma)$ the same as the one in (3.6), provided that γ is considered now as a polygon in \mathbb{Z}^2 , rather than in Λ (note that such identification is possible as long as γ does not wrap around the cylinder). Moreover, W_∞^{int} is translation invariant, and, letting

$$W_E^{\text{int}}(X, Y) := W_\Lambda^{\text{int}}(X, Y) - \mathbb{1}(\text{diam}_1(X \cup Y) \leq L/3) W_\infty^{\text{int}}(X_\infty, Y_\infty), \tag{3.11}$$

the contribution to the first term in the right side of all multipolygons in C_Λ with horizontal diameter $\leq L/3$ cancels with their counterparts in C_∞ from

for $i \geq 2$ are chosen in such a way that the horizontal distances between the corresponding pairs in X and X_∞ are the same, if measured on the cylinder Λ or on \mathbb{Z}^2 , respectively.

⁷We say that $X \subset \mathfrak{B}_\Lambda$ ‘touches the boundary of the cylinder Λ ’, if at least one of the edges in X has an endpoint whose vertical coordinate is either equal to 1 or to M .

the second term in the right side. Each of the remaining multipolygons either comes from the first term in the right side and involves a multipolygon in C_Λ whose support has horizontal diameter larger than $L/3$, or comes from the second term in the right side and involves a multipolygon in C_∞ whose support contains a set Z_∞ which is the representative (in the sense explained before the statement of the lemma) of a connected subset Z_Λ of \mathfrak{B}_Λ that contains $X \cup Y$ and touches the boundary of Λ ; in either case the number of edges in the support of such a multipolygon is at least $\delta_E(X \cup Y)$, from which the bound (3.8) follows. \square

In the following, we will wish to work in the ϕ, ξ variables introduced in Sect. 2.1.2; applying the change of variables (2.1.10), with some abuse of notation we rewrite $\mathcal{V}(\Phi, \mathbf{A})$ in (3.2) as:

$$\begin{aligned} \mathcal{V}(\phi, \xi, \mathbf{A}) &= \sum_{x \in \mathfrak{B}_\Lambda} \sum_{(\omega, z) \in \mathcal{O}^2 \times \Lambda^2} B_\Lambda^{\text{free}}((\omega, z), x) \phi(\omega, z) A_x \\ &\quad + \sum_{\substack{n \in 2\mathbb{N} \\ m \in \mathbb{N}_0}} \sum_{(\omega, z) \in \mathcal{O}^n \times \Lambda^n} \sum_{\mathbf{x} \in \mathfrak{B}_\Lambda^m} W_\Lambda^{\text{int}}((\omega, z), \mathbf{x}) \phi(\omega, z) \mathbf{A}(\mathbf{x}) \\ &=: \mathcal{B}^{\text{free}}(\phi, \xi, \mathbf{A}) + \mathcal{V}^{\text{int}}(\phi, \xi, \mathbf{A}), \end{aligned} \tag{3.12}$$

where \mathbb{N} and \mathbb{N}_0 are the sets of positive and nonnegative integers, respectively, $\mathcal{O} := \{1, -1, i, -i\}$, and, for $\omega \in \mathcal{O}^n, z \in \Lambda$, we denote

$$\phi(\omega, z) = \prod_{j=1}^n \phi_{\omega_j, z_j} \tag{3.13}$$

with $\phi_{\pm i, z} = \xi_{\pm, z}$, and similarly $\mathbf{A}(\mathbf{x}) := \prod_{j=1}^m A_{x_j}$ (for $\mathbf{x} = \emptyset$, we interpret $\mathbf{A}(\emptyset) = 1$). The decay properties of s_\pm noted after Eq. (2.1.11) together with Eq. (3.3) imply a similar decay property for the new coefficients:

$$\sup_{\omega \in \mathcal{O}^2} \sup_{x \in \mathfrak{B}_\Lambda} \sum_{z \in \Lambda^2} e^{c\delta(z, x)} |B_\Lambda^{\text{free}}((\omega, z), x)| \leq C, \tag{3.14}$$

and

$$\begin{aligned} \sup_{\omega \in \mathcal{O}^n} \sup_{z_1 \in \Lambda} \sum_{z_2, \dots, z_n \in \Lambda} e^{c\delta(z)} |W_\Lambda^{\text{int}}((\omega, z), \emptyset)| &\leq C^n |\lambda|^{\max(1, \kappa n)} \\ \sup_{\omega \in \mathcal{O}^n} \sup_{x_1 \in \mathfrak{B}_\Lambda} \sum_{x_2, \dots, x_m \in \mathfrak{B}_\Lambda} \sum_{z \in \Lambda^n} e^{c\delta(z, x)} |W_\Lambda^{\text{int}}((\omega, z), \mathbf{x})| &\leq C^{n+m} |\lambda|^{\max(1, \kappa(n+m))}. \end{aligned} \tag{3.15}$$

Note that, with this rewriting in terms of the ϕ, ξ variables, recalling that $\mathcal{S}_{t_1, t_2}(\Phi) = \mathcal{S}_m(\xi) + \mathcal{S}_c(\phi)$, see (2.1.12)–(2.1.13), and denoting $P_c(\mathcal{D}\phi) := \mathcal{D}\phi e^{\mathcal{S}_c(\phi)} / \text{Pf}(A_c), P_m(\mathcal{D}\xi) := \mathcal{D}\xi e^{\mathcal{S}_m(\xi)} / \text{Pf}(A_m)$ (here A_c and A_m are the

two $2|\Lambda| \times 2|\Lambda|$ anti-symmetric matrices associated with the Grassmann quadratic forms $\mathcal{S}_c(\phi)$ and $\mathcal{S}_m(\xi)$, respectively), the Grassmann generating function $\tilde{\Xi}_\Lambda(\mathbf{A})$ in (3.1) can be rewritten as

$$\tilde{\Xi}_\Lambda(\mathbf{A}) \propto e^{\mathcal{V}(\mathbf{A})} \int P_c(\mathcal{D}\phi) \int P_m(\mathcal{D}\xi) e^{\mathcal{V}(\phi, \xi, \mathbf{A})}, \tag{3.16}$$

where \propto means ‘up to a multiplicative constant independent of \mathbf{A} ’. In view of these rewritings, Proposition 3.1 implies [4, Proposition 2.5] as an immediate corollary.

Of course, the bulk–edge decomposition of Lemma 3.2 implies an analogous decomposition for the kernel of $\mathcal{V}^{\text{int}}(\phi, \xi, \mathbf{A})$, which reads as follows:

$$\begin{aligned} W_\Lambda^{\text{int}}((\omega, \mathbf{z}), \mathbf{x}) &= (-1)^{\alpha(\mathbf{z})} \mathbb{1}(\text{diam}_1(\mathbf{z}, \mathbf{x}) \leq L/3) W_\infty^{\text{int}}((\omega, \mathbf{z}_\infty), \mathbf{x}_\infty) \\ &\quad + W_E^{\text{int}}((\omega, \mathbf{z}), \mathbf{x}) \\ &=: W_B^{\text{int}}((\omega, \mathbf{z}), \mathbf{x}) + W_E^{\text{int}}((\omega, \mathbf{z}), \mathbf{x}), \end{aligned} \tag{3.17}$$

where, for any \mathbf{z} with $\text{diam}_1(\mathbf{z}) \leq L/3$,

$$\alpha(\mathbf{z}) = \begin{cases} \#\{z_i \in \mathbf{z} : (z_i)_1 \leq L/3\}, & \text{if } \max_{z_i, z_j \in \mathbf{z}} \{(z_i)_1 - (z_j)_1\} \geq 2L/3, \\ 0, & \text{otherwise.} \end{cases} \tag{3.18}$$

and

$$W_\infty^{\text{int}}((\omega, \mathbf{z}), \mathbf{x}) := \lim_{L, M \rightarrow \infty} W_\Lambda^{\text{int}}((\omega, \mathbf{z} + z_{L, M}), \mathbf{x} + z_{L, M}). \tag{3.19}$$

The factor $(-1)^{\alpha(\mathbf{z})}$ in front of the first term in the right side of (3.17), in light of the antiperiodicity of the ϕ, ξ fields, guarantees that W_∞^{int} is translation invariant (in both coordinate directions), and that both W_B^{int} and W_E^{int} are invariant under simultaneous translations of \mathbf{z} and \mathbf{x} in the horizontal direction, with anti-periodic and periodic boundary conditions in \mathbf{z} and \mathbf{x} , respectively. In terms of this new notation, Eq.(3.8) implies that, for any $n \in \mathbb{N}$ and $m \in \mathbb{N}_0$,

$$\frac{1}{L} \sup_{\omega \in \mathcal{O}^n} \sum_{\substack{\mathbf{z} \in \Lambda^n \\ \mathbf{x} \in \mathfrak{B}_\Lambda^m}} |W_E^{\text{int}}((\omega, \mathbf{z}), \mathbf{x})| e^{c\delta_E(\mathbf{z}, \mathbf{x})} \leq C^{m+n} |\lambda|^{\max(1, \kappa(m+n))}, \tag{3.20}$$

with $\delta_E(\mathbf{z}, \mathbf{x})$ is the ‘edge’ tree distance of (\mathbf{z}, \mathbf{x}) , i.e., the cardinality of the smallest connected subset of \mathfrak{B}_Λ that includes \mathbf{x} , touches the points of \mathbf{z} and either touches the boundary of the cylinder or it has horizontal diameter larger than $L/3$. Of course, B_Λ^{free} admits a similar bulk–edge decomposition: in analogy with (3.7), $B_\Lambda^{\text{free}} = B_B^{\text{free}} + B_E^{\text{free}}$ with

$$\frac{1}{L} \sup_{\omega \in \mathcal{O}^2} \sum_{\substack{\mathbf{z} \in \Lambda^2 \\ \mathbf{x} \in \mathfrak{B}_\Lambda}} |B_E^{\text{free}}(\omega, \mathbf{z}, \mathbf{x})| e^{c\delta_E(\mathbf{z}, \mathbf{x})} \leq C. \tag{3.21}$$

Before concluding this section, let us comment on the connection between (3.16) and (1.7). Fix once and for all a neighborhood $U \subset \mathbb{R}$ of 1 not containing 0; say, for definiteness, $U := \{z \in \mathbb{R} : |z - 1| \leq 1/2\}$. For any $Z \in U$ and

$t_1^* \in K'$, we let $t_2^* := (1 - t_1^*) / (1 + t_1^*)$ and let $\mathcal{S}_c^*(\phi) = \frac{1}{2}(\phi, A_c^* \phi)$ (resp. $\mathcal{S}_m^*(\xi) = \frac{1}{2}(\xi, A_m^* \xi)$) be obtained from \mathcal{S}_c (resp. \mathcal{S}_m) by replacing t_1, t_2 with t_1^*, t_2^* in Eq. (2.1.13) (resp. Eq. (2.1.12)). We also let $P_c^*(\mathcal{D}\phi) := \mathcal{D}\phi e^{\mathcal{S}_c^*(\phi)} / \text{Pf}(A_c^*)$, $P_m^*(\mathcal{D}\xi) := \mathcal{D}\xi e^{\mathcal{S}_m^*(\xi)} / \text{Pf}(A_m^*)$. Given these definitions, in (3.16) we first rescale the ϕ and ξ variables by $Z^{-1/2}$, then multiply and divide the Grassmann integrand by $e^{\mathcal{S}_c^*(\phi) + \mathcal{S}_m^*(\xi)}$, thus getting

$$\tilde{\Xi}(\mathbf{A}) \propto e^{\mathcal{W}(\mathbf{A})} \int P_c^*(\mathcal{D}\phi) P_m^*(\mathcal{D}\xi) e^{\mathcal{V}^{(1)}(\phi, \xi, \mathbf{A})} \tag{3.22}$$

with

$$\begin{aligned} \mathcal{V}^{(1)}(\phi, \xi, \mathbf{A}) &:= Z^{-1} \mathcal{S}_c(\phi) - \mathcal{S}_c^*(\phi) + Z^{-1} \mathcal{S}_m(\xi) - \mathcal{S}_m^*(\xi) + \mathcal{V}(Z^{-1/2} \phi, Z^{-1/2} \xi, \mathbf{A}) \\ &=: \mathcal{N}_c(\phi) + \mathcal{N}_m(\xi) + \mathcal{V}(Z^{-1/2} \phi, Z^{-1/2} \xi, \mathbf{A}). \end{aligned} \tag{3.23}$$

This proves (1.7) and puts us in the position of setting the multiscale computation of the sequence of effective potentials, whose infinite plane counterparts are constructed and bounded in the next section. For later reference, we note that, in light of (2.1.12), (2.1.13), (3.12), $\mathcal{V}^{(1)}(\phi, \xi, \mathbf{A})$ can be written as:

$$\mathcal{V}^{(1)}(\phi, \xi, \mathbf{A}) = \sum_{\substack{n \in 2\mathbb{N}, \\ m \in \mathbb{N}_0}} \sum_{(\omega, z) \in \mathcal{O}^n \times \Lambda^n} \sum_{x \in \mathfrak{B}_\Lambda^m} W_\Lambda^{(1)}((\omega, z), x) \phi(\omega, z) \mathbf{A}(x), \tag{3.24}$$

for an appropriate kernel, which inherits its properties from those of \mathcal{S}_c , \mathcal{S}_m and \mathcal{V} . With no loss of generality, we can assume that $W_\Lambda^{(1)}$ is anti-symmetric under simultaneous permutations of ω and z , symmetric under permutations of x , invariant under simultaneous translations of z and x in the horizontal direction (with anti-periodic and periodic boundary conditions in z and x , respectively), invariant under the reflection symmetries induced by the transformations $A_x \rightarrow A_{\theta_l x}$ and $\phi_{\omega, z} \rightarrow \Theta_l \phi_{\omega, z}$, see (2.4.3)–(2.4.4). From now on, with some abuse of notation, given $\omega = (\omega_1, \dots, \omega_n) \in \mathcal{O}^n$ and $z = (z_1, \dots, z_n) \in \Lambda^n$, we shall identify the pair (ω, z) with the n -ple $((\omega_1, z_1), \dots, (\omega_n, z_n)) \in (\mathcal{O} \times \Lambda)^n$.

4. The Renormalized Expansion in the Full-Plane Limit

In this section we construct the sequence of effective potentials (see the last part of Sect. 1) in the infinite volume limit and derive weighted L^1 bounds for their kernels, in a form appropriate for the subsequent generalization to the finite cylinder, discussed in [4, Section 3]. The construction of this section will allow us, in particular, to fix the free parameters β, Z, t_1^* , which the Grassmann integral in the right side of (3.22) depends on, in such a way that the sequence of running coupling constants goes to zero exponentially fast in the infrared limit; see Sect. 4.5 below.

As anticipated at the end of Sect. 1, here we limit ourselves to construct the sequence of effective potentials at $\mathbf{A} = \mathbf{0}$, so, for lightness of notation, we denote by $\mathcal{V}^{(h)}(\phi) := \mathcal{V}^{(h)}(\phi, \mathbf{0})$ the effective potentials with $h \leq 0$ at zero

external fields (similarly, we let $\mathcal{V}^{(1)}(\phi, \xi) = \mathcal{V}^{(1)}(\phi, \xi, \mathbf{A})$). In light of (1.8) and (1.10), these effective potentials are iteratively defined via

$$\mathcal{V}^{(h-1)}(\phi) = \text{const.} + \begin{cases} \log \int P_m^*(\mathcal{D}\xi) e^{\mathcal{V}^{(1)}(\phi, \xi)} & \text{if } h = 1, \\ \log \int P^{(h)}(\mathcal{D}\varphi) e^{\mathcal{L}\mathcal{V}^{(h)}(\phi + \varphi) + \mathcal{R}\mathcal{V}^{(h)}(\phi + \varphi)} & \text{if } h \leq 0, \end{cases} \tag{4.1}$$

where the const. is fixed so that $\mathcal{V}^{(h)}(0) = 0$, for all $h \leq 0$, and $\mathcal{L}\mathcal{V}^{(h)} + \mathcal{R}\mathcal{V}^{(h)}$ is an equivalent rewriting of $\mathcal{V}^{(h)}$, to be defined (in the full plane limit) below. Expanding the exponential and the logarithm in the right side of (4.1) allows us to rewrite

$$\mathcal{V}^{(0)}(\phi) = \text{const.} + \sum_{s \geq 1} \frac{1}{s!} \mathbb{E}_m^* \left(\underbrace{\mathcal{V}^{(1)}(\phi, \cdot); \dots; \mathcal{V}^{(1)}(\phi, \cdot)}_{s \text{ times}} \right) \tag{4.2}$$

and, for $h \leq 0$,

$$\mathcal{V}^{(h-1)}(\phi) = \text{const.} + \sum_{s \geq 1} \frac{1}{s!} \mathbb{E}^{(h)} \left(\underbrace{\mathcal{L}\mathcal{V}^{(h)}(\phi + \cdot) + \mathcal{R}\mathcal{V}^{(h)}(\phi + \cdot); \dots; \mathcal{L}\mathcal{V}^{(h)}(\phi + \cdot) + \mathcal{R}\mathcal{V}^{(h)}(\phi + \cdot)}_{s \text{ times}} \right), \tag{4.3}$$

where \mathbb{E}_m^* (resp. $\mathbb{E}^{(h)}$) denotes the *truncated expectation* [15, Eq.(4.13)] with respect to the Grassmann Gaussian integration P_m^* (resp. $P^{(h)}$). Expanding the effective potentials in terms of their kernels, in analogy with (3.24), Eqs. (4.2)–(4.3) allow us to iteratively compute the kernels of $\mathcal{V}^{(h)}$, for all $h \leq 0$. For instance, at the first step, using (3.24), (4.2) and the BBFK formula (for Battle, Brydges, Federbush, Kennedy) for the Grassmann truncated expectations [1, 7–9], we find that, denoting by $V_\Lambda^{(1)}(\Psi) = W_\Lambda^{(1)}(\Psi, \emptyset)$ with $\Psi = ((\omega_1, z_1), \dots, (\omega_n, z_n)) \in \cup_{n \in 2\mathbb{N}} (\mathcal{O} \times \Lambda)^n =: \mathcal{M}_{1, \Lambda}$ the kernel of $\mathcal{V}^{(1)}(\phi, \xi)$, the kernel $V_\Lambda^{(0)}$ of $\mathcal{V}^{(0)}(\phi)$ satisfies, for any $\Psi = ((\omega_1, z_1), \dots, (\omega_n, z_n)) \in \cup_{n \in 2\mathbb{N}} (\{+, -\} \times \Lambda)^n$,

$$\begin{aligned} V_\Lambda^{(0)}(\Psi) &= \sum_{s=1}^\infty \frac{1}{s!} \sum_{\Psi_1, \dots, \Psi_s \in \mathcal{M}_{1, \Lambda}}^{(\Psi)} \sum_{T \in \mathcal{S}(\Psi_1, \dots, \Psi_s)} \mathfrak{G}_T^{(1)}(\bar{\Psi}_1, \dots, \bar{\Psi}_s) \\ &\quad \alpha(\Psi; \Psi_1, \dots, \Psi_s) \left(\prod_{j=1}^s V_\Lambda^{(1)}(\Psi_j) \right), \end{aligned} \tag{4.4}$$

where

- the symbol (Ψ) on the second sum means that the sum runs over all ways of representing Ψ as an ordered sum of s (possibly empty) tuples, $\Psi'_1 \oplus \dots \oplus \Psi'_s = \Psi$, and over all tuples $\mathcal{M}_{1, \Lambda} \ni \Psi_j \supseteq \Psi'_j$ (here \oplus indicates concatenation of ordered tuples); for each such term in the second sum,

we denote by $\bar{\Psi}_j := \Psi_j \setminus \Psi'_j$ and by $\alpha(\Psi; \Psi_1, \dots, \Psi_s)$ the sign of the permutation from $\Psi_1 \oplus \dots \oplus \Psi_s$ to $\Psi \oplus \bar{\Psi}_1 \oplus \dots \oplus \bar{\Psi}_s$ (here \oplus indicates concatenation of ordered tuples);

- $\mathcal{S}(\bar{\Psi}_1, \dots, \bar{\Psi}_s)$ denotes the set of all the ‘spanning trees’ on $\bar{\Psi}_1, \dots, \bar{\Psi}_s$, that is, of all the sets T of ordered pairs (f, f') , with $f \in \bar{\Psi}_i, f' \in \bar{\Psi}_j$ and $i < j$, whose corresponding graph $G_T = (V, E_T)$, with vertex set $V = \{1, \dots, s\}$ and edge set $E_T = \{(i, j) \in V^2 : \exists (f, f') \in T \text{ with } f \in Q_i, f' \in Q_j\}$, is a tree graph (for $s = 1$, we let $\mathcal{S}(\bar{\Psi}_1) \equiv \{\emptyset\}$);
- $\mathfrak{G}_T^{(1)}(\bar{\Psi}_1, \dots, \bar{\Psi}_s)$ is different from zero only if $\bar{\Psi}_j \in \cup_{2\mathbb{N}_0}(\{+i, -i\} \times \Lambda)^n$ for all $j = 1, \dots, s$, and, if $s > 1$, only if $\bar{\Psi}_j \neq \emptyset$ for all $j = 1, \dots, s$; more precisely: if $s = 1$ and $\bar{\Psi}_1 = \emptyset$, then $\mathfrak{G}_\emptyset^{(1)}(\emptyset) = 1$; if $s = 1$ and $\bar{\Psi}_1 \neq \emptyset$, then $\mathfrak{G}_\emptyset^{(1)}(\bar{\Psi}_1) = \text{Pf}(G_{\bar{\Psi}_1}^{(1)})$, where, given a pair $\ell = ((\omega, z), (\omega', z'))$ of distinct elements of $\bar{\Psi}_1$, $(G_{\bar{\Psi}_1}^{(1)})_\ell = g_\ell^{(1)} := g_{-i\omega, -i\omega'}(z, z')$ ⁸; if $s > 1$ and $\bar{\Psi}_j \neq \emptyset$ for all $j = 1, \dots, s$, then

$$\mathfrak{G}_T^{(1)}(\bar{\Psi}_1, \dots, \bar{\Psi}_s) := \alpha_T(\bar{\Psi}_1, \dots, \bar{\Psi}_s) \left[\prod_{\ell \in T} g_\ell^{(1)} \right] \int P_{\bar{\Psi}_1, \dots, \bar{\Psi}_s, T}(\mathbf{dt}) \text{Pf}(G_{\bar{\Psi}_1, \dots, \bar{\Psi}_s, T}^{(1)}(\mathbf{t})), \tag{4.5}$$

where

- $\alpha_T(\bar{\Psi}_1, \dots, \bar{\Psi}_s)$ is the sign of the permutation from $\bar{\Psi}_1 \oplus \dots \oplus \bar{\Psi}_s$ to $T \oplus (\bar{\Psi}_1 \setminus T) \oplus \dots \oplus (\bar{\Psi}_s \setminus T)$;
- $\mathbf{t} = \{t_{i,j}\}_{1 \leq i, j \leq s}$, and $P_{\bar{\Psi}_1, \dots, \bar{\Psi}_s, T}(\mathbf{dt})$ is a probability measure with support on a set of \mathbf{t} such that $t_{i,j} = \mathbf{u}_i \cdot \mathbf{u}_j$ for some family of vectors $\mathbf{u}_i = \mathbf{u}_i(\mathbf{t}) \in \mathbb{R}^s$ of unit norm;
- letting $2q = \sum_{i=1}^s |\bar{\Psi}_i|$, $G_{\bar{\Psi}_1, \dots, \bar{\Psi}_s, T}^{(1)}(\mathbf{t})$ is an antisymmetric $(2q - 2s + 2) \times (2q - 2s + 2)$ matrix, whose off-diagonal elements are given by $(G_{\bar{\Psi}_1, \dots, \bar{\Psi}_s, T}^{(1)}(\mathbf{t}))_{f, f'} = t_{i(f), i(f')} g_{\ell(f, f')}^{(1)}$, where f, f' are elements of the tuple $(\bar{\Psi}_1 \setminus T) \oplus \dots \oplus (\bar{\Psi}_s \setminus T)$, and $i(f)$ is the integer in $\{1, \dots, s\}$ such that f is an element of $\bar{\Psi}_i \setminus T$.

Recalling that $g_\ell^{(1)}$ and $V_\Lambda^{(1)}$ admit infinite volume limits $g_{\ell, \infty}^{(1)}$ and $V_\infty^{(1)}$, respectively, in the sense of (2.2.9) and (3.19), from (4.4) it follows that $V_\Lambda^{(0)}$ admits an infinite volume limit as well, equal to the ‘obvious’ analogue of the right side of (4.4), namely, the expression obtained from that one by replacing: $\mathcal{M}_{1, \Lambda}$ by $\mathcal{M}_{1, \infty} := \cup_{n \in 2\mathbb{N}} (\mathcal{O} \times \mathbb{Z}^2)^n$; $\mathfrak{G}_T^{(1)}$ by $\mathfrak{G}_{T, \infty}^{(1)}$ (the latter being defined analogously to the former, with $g_\ell^{(1)}$ replaced by $g_{\ell, \infty}^{(1)}$ in all the involved expressions); and $V_\Lambda^{(1)}$ by $V_\infty^{(1)}$.

⁸Here $g_{\sigma, \sigma'}^{(1)}(z, z')$ with $\sigma, \sigma' \in \{+, -\}$ are the elements of the 2×2 matrix $\mathfrak{g}^{(1)}(z, z') \equiv \mathfrak{g}_m^*(z, z')$, where $\mathfrak{g}_m^*(z, z')$ the same as in (2.1.16) with t_1 replaced by t_1^* .

Proceeding inductively in $h \leq 0$, one finds that (4.3) implies a representation of the kernel $V_\Lambda^{(h)}$ of $\mathcal{V}^{(h)}$ analogous to (4.4). Also in that case, the resulting formula for $V_\Lambda^{(h)}$ admits a natural infinite volume limit. In this way, we obtain a recursive equation for the infinite plane kernels, denoted $V_\infty^{(h)}$, whose definition and solution is described below. Convergence of the finite volume kernels $V_\Lambda^{(h)}$ to their infinite volume counterparts, with optimal bounds on the norm of the finite size corrections, is deferred to [4, Section 3]. A key point in the derivation of bounds on the kernels that are uniform in the scale label h is the definition of an appropriate action of the \mathcal{L} and \mathcal{R} operators, as well as of their infinite volume counterparts, \mathcal{L}_∞ and \mathcal{R}_∞ . As anticipated above, these operators allow us to isolate the potentially divergent part of the kernels, $\mathcal{L}\mathcal{V}^{(h)}$ (the ‘local’ contributions, parametrized at any given scale, by a finite number of ‘running coupling constants’) from a remainder $\mathcal{R}\mathcal{V}^{(h)}$, which is ‘dimensionally better behaved’ than $\mathcal{L}\mathcal{V}^{(h)}$; in order for the remainder to be shown to satisfy ‘improved dimensional bounds’, it is necessary to rewrite it in an appropriate, interpolated, form, involving the action of discrete derivatives on the Grassmann fields.

The plan of the incoming subsections is the following: in Sect. 4.1 we describe the representation of the effective potentials in the infinite volume limit and introduce the notion of equivalent kernels; in Sect. 4.2 we define the operators \mathcal{L}_∞ and \mathcal{R}_∞ ; in Sect. 4.3 we derive the solution to the recursive equations for the infinite volume kernels in terms of a tree expansion; in Sect. 4.4, we use such a tree expansion to derive weighted L^1 bounds on the kernels; importantly, these bounds depend upon the sequence of running coupling constants, and they imply analyticity of the kernels provided such a sequence is uniformly bounded in the scale label; in Sect. 4.5, as a corollary of the weighted L^1 bounds of the previous subsection, we prove a fixed point theorem, which allows us to fix the free parameters Z, β, t_1^* in such a way that the flow of the running coupling constants is, in fact, uniformly bounded in h , as desired: even more, the running coupling constants go to zero exponentially fast as $h \rightarrow -\infty$, a consequence of the irrelevance of the quartic effective interaction in the theory at hand.

4.1. Effective Potentials and Kernels: Representation and Equivalence

In this subsection we define the effective potential in the full plane in terms of equivalence classes of *kernels* $V(\Psi)$, namely, of real-valued functions playing the same role as the coefficient functions $V_\infty^{(1)}(\Psi)$ and $V_\infty^{(0)}(\Psi)$ introduced above. This point of view avoids defining an infinite-dimensional Grassmann algebra. The equivalence relation among kernels, to be introduced momentarily, generalizes the relationships which hold between different ways of writing the coefficients of a given Grassmann polynomial.

As mentioned above, in order to obtain bounds on the kernels of the effective potentials which are uniform in the scale label, we will need to group some of the Grassmann fields into discrete derivatives; we will mainly use the directional derivative $\partial_j \phi_{\omega, z} := \phi_{\omega, z + \hat{e}_j} - \phi_{\omega, z}$ (note that this is the same

convention used in Sect. 2.2.2). We consequently consider kernels which specify when and how this is done, and in particular define the equivalence relationship with this in mind.

Let $\Lambda_\infty := \mathbb{Z}^2$, let \mathfrak{B} denote the set of nearest neighbor edges of Λ_∞ , and let $\mathcal{D} := \{D \in \{0, 1, 2\}^2 : \|D\|_1 \leq 2\}$. Let $\mathcal{M}_\infty = \cup_{n \in 2\mathbb{N}} (\{+, -\} \times \mathcal{D} \times \Lambda_\infty)^n$ be the set of *field multilabels*. for some $n \in 2\mathbb{N}$, such that $\|D_i\|_1 \leq 2$. We can think of any $\Psi = ((\omega_1, D_1, z_1), \dots, (\omega_n, D_n, z_n)) \in \mathcal{M}_\infty$ as indexing a formal Grassmann monomial $\phi(\Psi)$ given by

$$\phi(\Psi) = \partial^{D_1} \phi_{\omega_1, z_1} \cdots \partial^{D_n} \phi_{\omega_n, z_n}, \tag{4.1.1}$$

where, denoting $D_i = ((D_i)_1, (D_i)_2) \in \mathcal{D}$, we let

$$\partial^{D_i} \phi_{\omega_i, z_i} := \partial_1^{(D_i)_1} \partial_2^{(D_i)_2} \phi_{\omega_i, z_i},$$

with ∂_1 and ∂_2 the (right) discrete derivatives introduced above. In the following, with some abuse of notation, any element $\Psi \in \mathcal{M}_\infty$ of length $|\Psi| = n$ will be denoted indistinctly by $\Psi = ((\omega_1, D_1, z_1), \dots, (\omega_n, D_n, z_n))$ or $\Psi = (\omega, \mathbf{D}, \mathbf{z})$, with the understanding that $\omega = (\omega_1, \dots, \omega_n)$, etc.

We will call a function $V : \mathcal{M}_\infty \rightarrow \mathbb{R}$ a *kernel function*, let V_n denote its restriction to field multilabels of length n , and let $V_{n,p}$ be the restriction of V_n to field multilabels with $\|\mathbf{D}\|_1 = p$. Thinking of such a V as the coefficient function of a formal Grassmann polynomial

$$\mathcal{V}(\phi) = \sum_{\Psi \in \mathcal{M}_\infty} V(\Psi) \phi(\Psi) \tag{4.1.2}$$

suggests an equivalence relationship corresponding to manipulations allowed by the anticommutativity of the Grassmann variables and by the definition of discrete derivative.

More precisely, we say that V is equivalent to V' , and write $V \sim V'$, if either:

1. V' is obtained from V by permuting the arguments and changing the sign according to the parity of the permutation;
2. V' is obtained from V by writing out the action of a derivative: that is, there exist $n \in 2\mathbb{N}$, $i \in \{1, \dots, n\}$ and $j \in \{1, 2\}$ such that, letting $\mathbf{D}_{i,j}^+ = (D_1, \dots, D_{i-1}, D_i + \hat{e}_j, D_{i+1}, \dots, D_n)$ and $\mathbf{z}_{i,j}^- = (z_1, \dots, z_{i-1}, z_i - \hat{e}_j, z_{i+1}, \dots, z_n)$,

$$V'_{n,p}(\omega, \mathbf{D}, \mathbf{z}) = \begin{cases} 0 & \text{if } (D_i)_j = 2, \\ V_n(\omega, \mathbf{D}_{i,j}^+, \mathbf{z}_{i,j}^-) - V_n(\omega, \mathbf{D}_{i,j}^+, \mathbf{z}) & \text{if } (D_i)_j = 1, \\ V_n(\omega, \mathbf{D}, \mathbf{z}) + V_n(\omega, \mathbf{D}_{i,j}^+, \mathbf{z}_{i,j}^-) - V_n(\omega, \mathbf{D}_{i,j}^+, \mathbf{z}) & \text{if } (D_i)_j = 0, \end{cases} \tag{4.1.3}$$

while $V'_m = V_m$ for all $m \in 2\mathbb{N} \setminus \{n\}$;

3. V' is obtained from V by adding an arbitrary kernel V^* that is different from zero only for arguments with common repetition, that is, $V^*(\omega, \mathbf{D}, \mathbf{z}) = 0$ unless there is some $i \neq j$ such that $(\omega_i, D_i, z_i) = (\omega_j, D_j, z_j)$;

or V' is obtained from V by a countable sequence of such elementary operations and of convex combinations thereof. Moreover, we assume that the equivalence relation \sim is preserved by linear combinations, i.e., if $V'_\alpha \sim V_\alpha$ for all α in the countable index set \mathcal{I} , then $\sum_{\alpha \in \mathcal{I}} V'_\alpha \sim \sum_{\alpha \in \mathcal{I}} V_\alpha$. We will call the equivalence classes generated by \sim *potentials* and often specify them by formal sums like (4.1.2).

Remark 4.1. The operation in item 2 can be thought of as a form of ‘integration by parts’. The kernels equivalent to zero, $V \sim 0$, correspond to what are known as ‘null fields’ in the literature on conformal field theories.

4.2. Localization and Interpolation

In this section we define the operators \mathcal{L}_∞ and \mathcal{R}_∞ acting on kernels indexed by field multilabels in \mathcal{M}_∞ and show several estimates related to \mathcal{R}_∞ . We recall that, given a kernel V , the symbol $V_{n,p}$ denotes its restriction to field multilabels of length n , such that $\|\mathbf{D}\|_1 = p$.

The operator \mathcal{L}_∞ . First of all, we let

$$\mathcal{L}_\infty(V_{n,p}) := 0, \quad \text{if } 2 - \frac{n}{2} - p < 0. \tag{4.2.1}$$

In the RG jargon, the combination $2 - \frac{n}{2} - p$ is the *scaling dimension* of $V_{n,p}$ and will reappear below, for example in Lemma 4.7; in this sense, (4.2.1) says that the local part of the terms with negative scaling dimension (the irrelevant terms, in the RG jargon) is set equal to zero.

There are only a few cases for which $2 - \frac{n}{2} - p \geq 0$, namely $(n, p) = (2, 0), (2, 1), (4, 0)$. In these cases, the action of \mathcal{L}_∞ on $V_{n,p}$ is non-trivial and will be defined in terms of other basic operators, the first of which is $\tilde{\mathcal{L}}$, which is defined as: $\tilde{\mathcal{L}}(V_{n,p}) = (\tilde{\mathcal{L}}V)_{n,p} \equiv \tilde{\mathcal{L}}V_{n,p}$, with

$$\tilde{\mathcal{L}}V_{n,p}(\omega, \mathbf{D}, (z_1, \dots, z_n)) := \left(\prod_{j=2}^n \delta_{z_j, z_1} \right) \sum_{y_2, \dots, y_n \in \Lambda_\infty} V_{n,p}(\omega, \mathbf{D}, (z_1, y_2, \dots, y_n)). \tag{4.2.2}$$

If $(n, p) = (2, 1), (4, 0)$, we let

$$\mathcal{L}_\infty(V_{2,1}) := \mathcal{A}(\tilde{\mathcal{L}}V_{2,1}), \quad \mathcal{L}_\infty V_{4,0} := \mathcal{A}(\tilde{\mathcal{L}}V_{4,0}), \tag{4.2.3}$$

where \mathcal{A} is the operator that antisymmetrizes with respect to permutations *and* symmetrizes with respect to reflections in the horizontal and vertical directions⁹. A first important remark, related to the definitions (4.2.3), is that, if $V_{2,1}$ is invariant under translations and under the action of \mathcal{A} , then

$$\mathcal{L}_\infty(V_{2,1}) = c_1 F_{\zeta, \infty} + c_2 F_{\eta, \infty}, \tag{4.2.4}$$

⁹The reflection transformations in the infinite plane act on the Grassmann fields in the same way as (2.4.5), with the difference that the reflections $\theta_1(x_1, x_2)$ and $\theta_2(x_1, x_2)$ are replaced by their infinite-plane analogues, namely by $\tilde{\theta}_1(x_1, x_2) = (-x_1, x_2)$ and $\tilde{\theta}_2(x_1, x_2) = (x_1, -x_2)$, respectively.

for some real numbers c_1, c_2 and $F_{\zeta, \infty}, F_{\eta, \infty}$ the \mathcal{A} -invariant kernels associated with the potentials

$$\mathcal{F}_{\zeta, \infty}(\phi) := \sum_{\omega=\pm} \sum_{z \in \Lambda_\infty} \omega \phi_{\omega, z} d_1 \phi_{\omega, z}, \quad \mathcal{F}_{\eta, \infty}(\phi) := \sum_{\omega=\pm} \sum_{z \in \Lambda_\infty} \phi_{\omega, z} d_2 \phi_{-\omega, z},$$

with d_j the symmetric discrete derivatives, acting on the Grassmann fields as $d_j \phi_{\omega, z} = \frac{1}{2}(\partial_j \phi_{\omega, z} + \partial_j \phi_{\omega, z - \hat{e}_j})$. A second and even more important remark is that, due to the fact that ω only assumes two values and that $\tilde{\mathcal{L}}V_{4,0}$ is supported on \mathbf{z} such that $z_1 = z_2 = z_3 = z_4$, one has

$$\mathcal{L}_\infty(V_{4,0}) = 0, \tag{4.2.5}$$

a cancellation that will play an important role in the following.

In order to define the action of \mathcal{L}_∞ on $V_{2,0}$, we want to obtain a kernel function equivalent to $V_{2,0} - \tilde{\mathcal{L}}V_{2,0}$, denoted by $(\tilde{\mathcal{R}}V)_{2,1}$, which is supported on arguments with an additional derivative. As we will see, the definition of $\tilde{\mathcal{R}}$ will also play a central role in the definition of the operator \mathcal{R}_∞ below. We rewrite

$$\begin{aligned} & \sum_{z \in \Lambda_\infty^2} [V_{2,0}(\boldsymbol{\omega}, \mathbf{0}, \mathbf{z}) - \tilde{\mathcal{L}}V_{2,0}(\boldsymbol{\omega}, \mathbf{0}, \mathbf{z})] \phi(\boldsymbol{\omega}, \mathbf{0}, \mathbf{z}) \\ &= \sum_{z \in \Lambda_\infty^2} V_{2,0}(\boldsymbol{\omega}, \mathbf{0}, \mathbf{z}) [\phi(\boldsymbol{\omega}, \mathbf{0}, \mathbf{z}) - \phi(\boldsymbol{\omega}, \mathbf{0}, (z_1, z_1))] \\ &= \sum_{z \in \Lambda_\infty^2} V_{2,0}(\boldsymbol{\omega}, \mathbf{0}, \mathbf{z}) \phi_{\omega_1, z_1} (\phi_{\omega_2, z_2} - \phi_{\omega_2, z_1}). \end{aligned} \tag{4.2.6}$$

We now intend to write the difference $\phi_{\omega_1, z_1} (\phi_{\omega_2, z_2} - \phi_{\omega_2, z_1})$ as a sum of terms of the form $\phi_{\omega_1, z_1} \partial^{D'} \phi_{\omega_2, y}$, with $\|D'\|_1 = 1$, over the sites y on a path from z_1 to z_2 . To do this we must specify which path is to be used.

For each $z, z' \in \Lambda_\infty$, let $\gamma(z, z') = (z, z_1, z_2, \dots, z_n, z')$ be the shortest path obtained by going first horizontally and then vertically from z to z' . Note that γ is covariant under the symmetries of the model on the infinite plane, i.e.,

$$S\gamma(z, z') = \gamma(Sz, Sz') \tag{4.2.7}$$

where $S : \Lambda_\infty \rightarrow \Lambda_\infty$ is some composition of translations and reflections parallel to the coordinate axes. Given z, z' two distinct sites in Λ_∞ , let $\text{INT}(z, z')$ be the set of $(\sigma, (D_1, D_2), (y_1, y_2)) \equiv (\sigma, \mathbf{D}, \mathbf{y}) \in \{\pm\} \times \{0, \hat{e}_1, \hat{e}_2\}^2 \times \Lambda_\infty^2$ such that: (1) $y_1 = z$, (2) $D_1 = 0$, (3) $y_2, y_2 + D_2 \in \gamma(z, z')$, (4) $\sigma = +$ if y_2 precedes $y_2 + D_2$ in the sequence defining $\gamma(z, z')$, and $\sigma = -$ otherwise. In terms of this definition, one can easily check that (4.2.6) can be rewritten as

$$\begin{aligned} (4.2.6) &= \sum_{z \in \Lambda_\infty^2} V_{2,0}(\boldsymbol{\omega}, \mathbf{0}, \mathbf{z}) \sum_{(\sigma, \mathbf{D}, \mathbf{y}) \in \text{INT}(z)} \sigma \phi(\boldsymbol{\omega}, \mathbf{D}, \mathbf{y}) \\ &\equiv \sum_{\mathbf{y} \in \Lambda_\infty^2} \sum_{\mathbf{D}}^{(1)} (\tilde{\mathcal{R}}V)_{2,1}(\boldsymbol{\omega}, \mathbf{D}, \mathbf{y}) \phi(\boldsymbol{\omega}, \mathbf{D}, \mathbf{y}), \end{aligned} \tag{4.2.8}$$

where, if $z = (z_1, z_1)$, the sum over $(\sigma, \mathbf{D}, \mathbf{y})$ in the first line should be interpreted as being equal to zero (in this case, we let $\text{INT}(z_1, z_1)$ be the empty set). In going from the first to the second line, we exchanged the order of summations over z and \mathbf{y} ; moreover, $\sum_D^{(p)}$ denotes the sum over the pairs $\mathbf{D} = (D_1, D_2)$ such that $\|\mathbf{D}\|_1 = p$, and

$$(\tilde{\mathcal{R}}V)_{2,1}(\omega, \mathbf{D}, \mathbf{y}) := \sum_{\substack{\sigma, z: \\ (\sigma, \mathbf{D}, \mathbf{y}) \in \text{INT}(z)}} \sigma V_{2,0}(\omega, \mathbf{0}, z). \tag{4.2.9}$$

From the previous manipulations, it is clear that $(\tilde{\mathcal{R}}V)_{2,1} \sim V_{2,0} - \tilde{\mathcal{L}}V_{2,0}$. We are finally ready to define the action of \mathcal{L}_∞ on $V_{2,0}$:

$$\mathcal{L}_\infty(V_{2,0}) := \mathcal{A}(\tilde{\mathcal{L}}V_{2,0} + \tilde{\mathcal{L}}(\tilde{\mathcal{R}}V)_{2,1}). \tag{4.2.10}$$

Note that, if $V_{2,0}$ is invariant under translations and under the action of \mathcal{A} , then

$$\mathcal{L}_\infty(V_{2,0}) = c_0 F_{\nu, \infty} + c_1 F_{\zeta, \infty} + c_2 F_{\eta, \infty}, \tag{4.2.11}$$

for some real numbers c_0, c_1, c_2 , and $F_{\nu, \infty}$ the \mathcal{A} -invariant kernel associated with the potential

$$\mathcal{F}_{\nu, \infty}(\phi) := \frac{1}{2} \sum_{\omega = \pm} \sum_{z \in \Lambda_\infty} \omega \phi_{\omega, z} \phi_{-\omega, z}, \tag{4.2.12}$$

while we recall that $F_{\zeta, \infty}, F_{\eta, \infty}$ were defined right after (4.2.4). Summarizing, in view of Eq. (4.2.5),

$$(\mathcal{L}_\infty V)_{n,p} = \begin{cases} \mathcal{A}(\tilde{\mathcal{L}}V_{2,0}) & \text{if } (n, p) = (2, 0), \\ \mathcal{A}(\tilde{\mathcal{L}}V_{2,1} + \tilde{\mathcal{L}}(\tilde{\mathcal{R}}V)_{2,1}) & \text{if } (n, p) = (2, 1), \\ 0 & \text{otherwise.} \end{cases} \tag{4.2.13}$$

The operator \mathcal{R}_∞ . We now want to define an operator \mathcal{R}_∞ such that $\mathcal{R}_\infty V \sim V - \mathcal{L}_\infty V$ for kernels V that are invariant under translations and under the action of \mathcal{A} . First of all, we let

$$\begin{aligned} \mathcal{R}_\infty(V_{n,p}) = (\mathcal{R}_\infty V)_{n,p} &:= V_{n,p}, & \text{if: } n \geq 6, & \text{ or } n = 4 \text{ and } p \geq 2, \\ & & \text{or } n = 2 \text{ and } p \geq 3. & \end{aligned} \tag{4.2.14}$$

Moreover, we let

$$(\mathcal{R}_\infty V)_{2,0} = (\mathcal{R}_\infty V)_{2,1} = (\mathcal{R}_\infty V)_{4,0} := 0. \tag{4.2.15}$$

The only remaining cases are $(n, p) = (2, 2), (4, 1)$. For these values of (n, p) , $(\mathcal{R}_\infty V)_{n,p}$ is defined in terms of an interpolation generalizing the definition of $(\tilde{\mathcal{R}}V)_{2,1}$ in (4.2.9). As a preparation for the definition, we first introduce $(\tilde{\mathcal{R}}V)_{n,p}$ for $(n, p) = (2, 2), (4, 1)$. For this purpose, we start from the analogues of (4.2.6), (4.2.8) in the case that $(2, 0)$ is replaced by $(n, p) = (2, 1), (4, 0)$: for such values of (n, p) we write

$$\begin{aligned}
 & \sum_{z \in \Lambda_\infty^n} [V_{n,p}(\omega, \mathbf{D}, z) - \tilde{\mathcal{L}}V_{n,p}(\omega, \mathbf{D}, z)]\phi(\omega, \mathbf{D}, z) \\
 &= \sum_{z \in \Lambda_\infty^n} V_{n,p}(\omega, \mathbf{D}, z)[\phi(\omega, \mathbf{D}, z) - \phi(\omega, \mathbf{D}, (z_1, z_1, \dots, z_1))] \\
 &= \sum_{z \in \Lambda_\infty^n} V_{n,p}(\omega, \mathbf{D}, z) \sum_{(\sigma, \mathbf{D}', \mathbf{y}) \in \text{INT}(z)} \sigma \phi(\omega, \mathbf{D} + \mathbf{D}', \mathbf{y}). \tag{4.2.16}
 \end{aligned}$$

In the last expression, if $n = 2$, then $\text{INT}(z)$ is the same defined after (4.2.7); if $n = 4$, then $\text{INT}(z)$ is the set of $(\sigma, (D_1, \dots, D_4), (y_1, \dots, y_4)) \equiv (\sigma, \mathbf{D}, \mathbf{y}) \in \{\pm\} \times \{0, \hat{e}_1, \hat{e}_2\}^4 \times \Lambda_\infty^4$ such that: either $y_1 = y_2 = y_3 = z_1, D_1 = D_2 = D_3 = 0$, and $(\sigma, (0, D_4), (z_1, y_4)) \in \text{INT}(z_1, z_4)$; or $y_1 = y_2 = z_1, y_4 = z_4, D_1 = D_2 = D_4 = 0$, and $(\sigma, (0, D_3), (z_1, y_3)) \in \text{INT}(z_1, z_3)$; or $y_1 = z_1, y_3 = z_3, y_4 = z_4, D_1 = D_3 = D_4 = 0$, and $(\sigma, (0, D_2), (z_1, y_2)) \in \text{INT}(z_1, z_2)$. By summing (4.2.16) over \mathbf{D} and exchanging the order of summations over z and \mathbf{y} , we find

$$\begin{aligned}
 & \sum_{z \in \Lambda_\infty^n} \sum_{\mathbf{D}}^{(p)} V_{n,p}(\omega, \mathbf{D}, z) \sum_{(\sigma, \mathbf{D}', \mathbf{y}) \in \text{INT}(z)} \sigma \phi(\omega, \mathbf{D} + \mathbf{D}', \mathbf{y}) \\
 &= \sum_{\mathbf{y} \in \Lambda_\infty^n} \sum_{\mathbf{D}}^{(p+1)} (\tilde{\mathcal{R}}V)_{n,p+1}(\omega, \mathbf{D}, \mathbf{y})\phi(\omega, \mathbf{D}, \mathbf{y}). \tag{4.2.17}
 \end{aligned}$$

with

$$(\tilde{\mathcal{R}}V)_{n,p+1}(\omega, \mathbf{D}, \mathbf{y}) := \sum_{\substack{\sigma, z, \mathbf{D}' \\ (\sigma, \mathbf{D}', \mathbf{y}) \in \text{INT}(z)}} \sigma V_{n,p}(\omega, \mathbf{D} - \mathbf{D}', z). \tag{4.2.18}$$

We are now ready to define:

$$(\mathcal{R}_\infty V)_{2,2} := \mathcal{A}(V_{2,2} + (\tilde{\mathcal{R}}V)_{2,2} + (\tilde{\mathcal{R}}(\tilde{\mathcal{R}}V))_{2,2}), \quad (\mathcal{R}_\infty V)_{4,1} := \mathcal{A}(V_{4,1} + (\tilde{\mathcal{R}}V)_{4,1}). \tag{4.2.19}$$

Summarizing,

$$(\mathcal{R}_\infty V)_{n,p} = \begin{cases} 0 & \text{if } (n, p) = (2, 0), (2, 1), (4, 0), \\ \mathcal{A}(V_{2,2} + (\tilde{\mathcal{R}}V)_{2,2} + (\tilde{\mathcal{R}}(\tilde{\mathcal{R}}V))_{2,2}) & \text{if } (n, p) = (2, 2), \\ \mathcal{A}(V_{4,1} + (\tilde{\mathcal{R}}V)_{4,1}) & \text{if } (n, p) = (4, 1), \\ V_{n,p} & \text{otherwise} \end{cases} \tag{4.2.20}$$

From the previous manipulations and definitions, it is clear that, if V is invariant under translations and under the action of \mathcal{A} , then $V - \mathcal{L}_\infty V \sim \mathcal{R}_\infty V$. For later use, given $\mathbf{D} = (D_1, \dots, D_n)$ with $\|\mathbf{D}\|_1 = p$, we let $\mathcal{R}_\infty V|_{\mathbf{D}}$ be the restriction of $(\mathcal{R}V)_{n,p}$ to that specific choice of derivative label.

Remark 4.2. From the definitions of \mathcal{L}_∞ and \mathcal{R}_∞ , it also follows that, if V is invariant under translations and under the action of \mathcal{A} , then $\mathcal{L}_\infty(\mathcal{L}_\infty V) = \mathcal{L}_\infty V$ and $\mathcal{R}_\infty(\mathcal{L}_\infty V) = 0$, two properties that will play a role in the following.

Norm bounds. Let us conclude this section by a couple of technical estimates, which relate a suitable weighted norm of $\mathcal{R}_\infty V$ to that of V , and will be useful in the following. Suppose that V is translationally invariant. Let, for any $\kappa \geq 0$,

$$\|V_{n,p}\|_{(\kappa)} := \sup_{\omega} \sum_{\substack{\mathbf{z} \in \Lambda_\infty^n: \\ z_1 \text{ fixed}}} e^{\kappa\delta(\mathbf{z})} \sup_{\mathbf{D}}^{(p)} |V_{n,p}(\omega, \mathbf{D}, \mathbf{z})|, \tag{4.2.21}$$

where the label (p) on the sup over \mathbf{D} indicates the constraint that $\|\mathbf{D}\|_1 = p$, and $\delta(\mathbf{z})$ is the tree distance of \mathbf{z} . With these definitions,

Lemma 4.3. *For any positive ϵ ,*

$$\|(\tilde{\mathcal{R}}V)_{n,p}\|_{(\kappa)} \leq (n-1)\epsilon^{-1}\|V_{n,p-1}\|_{(\kappa+\epsilon)} \text{ if } (n,p) = (2,2), (4,1), \tag{4.2.22}$$

$$\|(\tilde{\mathcal{R}}(\tilde{\mathcal{R}}V))_{2,2}\|_{(\kappa)} \leq \epsilon^{-2}\|V_{2,0}\|_{(\kappa+2\epsilon)}. \tag{4.2.23}$$

As a consequence, noting that $\|\mathcal{A}V_{n,p}\|_{(\kappa)} \leq \|V_{n,p}\|_{(\kappa)}$ and recalling the definitions (4.2.20), we find that

$$\|(\mathcal{R}_\infty V)_{2,2}\|_{(\kappa)} \leq \|V_{2,2}\|_{(\kappa)} + \epsilon^{-1}\|V_{2,1}\|_{(\kappa+\epsilon)} + \epsilon^{-2}\|V_{2,0}\|_{(\kappa+2\epsilon)}, \tag{4.2.24}$$

$$\|(\mathcal{R}_\infty V)_{4,1}\|_{(\kappa)} \leq \|V_{4,1}\|_{(\kappa)} + 3\epsilon^{-1}\|V_{4,0}\|_{(\kappa+\epsilon)}. \tag{4.2.25}$$

In the following, we will use bounds of this kind in order to evaluate the size of the renormalized part of the effective potential on scale h . In such a case, both κ and ϵ will be chosen of the order 2^h .

Proof. In order to prove (4.2.22), note that it follows directly from the definition of $\tilde{\mathcal{R}}$ that

$$\|(\tilde{\mathcal{R}}V)_{n,p}\|_{(\kappa)} \leq \sup_{\omega} \sum_{\substack{\mathbf{z} \in \Lambda_\infty^n: \\ z_1 \text{ fixed}}} \sup_{\mathbf{D}}^{(p)} \sum_{\substack{\sigma, \mathbf{y}, \mathbf{D}': \\ (\sigma, \mathbf{D}', \mathbf{z}) \in \text{INT}(\mathbf{y})}} e^{\kappa\delta(\mathbf{z})} |V_{n,p-1}(\omega, \mathbf{D} - \mathbf{D}', \mathbf{y})|. \tag{4.2.26}$$

If we now exchange the order of summations over \mathbf{z} and \mathbf{y} , we find

$$\|(\tilde{\mathcal{R}}V)_{n,p}\|_{(\kappa)} \leq \sup_{\omega} \sum_{\substack{\mathbf{y} \in \Lambda_\infty^n: \\ y_1 \text{ fixed}}} \sum_{(\sigma, \mathbf{D}', \mathbf{z}) \in \text{INT}(\mathbf{y})} e^{\kappa\delta(\mathbf{z})} \sup_{\mathbf{D}}^{(p-1)} |V_{n,p-1}(\omega, \mathbf{D}, \mathbf{y})|. \tag{4.2.27}$$

Now note that $\delta(\mathbf{z}) \leq \delta(\mathbf{y})$ and that $|\text{INT}(\mathbf{y})| \leq (n-1)\delta(\mathbf{y})$, so that

$$\begin{aligned} \|(\tilde{\mathcal{R}}V)_{n,p}\|_{(\kappa)} &\leq (n-1) \sup_{\omega} \sum_{\substack{\mathbf{y} \in \Lambda_\infty^n: \\ y_1 \text{ fixed}}} e^{\kappa\delta(\mathbf{y})} \delta(\mathbf{y}) \sup_{\mathbf{D}}^{(p-1)} |V_{n,p-1}(\omega, \mathbf{D}, \mathbf{y})| \\ &\leq \frac{n-1}{\epsilon} \|V_{n,p-1}\|_{\kappa+\epsilon}, \end{aligned} \tag{4.2.28}$$

where in the last step we used the fact that $\delta \leq e^{\epsilon\delta}/\epsilon$, for any $\epsilon > 0$. A two-step iteration of the bound (4.2.22) proves (4.2.23). \square

Similar estimates are valid for more general values of (n, p) , but Lemma 4.3 includes all the cases which are relevant to the present work.

The running coupling constants. At each scale $h \leq 0$ we represent the infinite volume effective potential, as arising from the iterative application of (4.3) in the infinite volume limit, in the form (4.1.2), namely:

$$\mathcal{V}_\infty^{(h)}(\phi) = \sum_{\Psi \in \mathcal{M}_\infty} V_\infty^{(h)}(\Psi)\phi(\Psi). \tag{4.2.29}$$

For each $h \leq 0$, in order to compute $\mathcal{V}_\infty^{(h-1)}$ from $\mathcal{V}_\infty^{(h)}$ via the infinite volume limit of (4.3), we decompose $V_\infty^{(h)} \sim \mathcal{L}_\infty V_\infty^{(h)} + \mathcal{R}_\infty V_\infty^{(h)}$. Note that the kernel $\mathcal{L}_\infty V_\infty^{(h)}$, in light of (4.2.4), (4.2.5) and (4.2.4), takes the form

$$\mathcal{L}_\infty V_\infty^{(h)} = 2^h \nu_h F_{\nu,\infty} + \zeta_h F_{\zeta,\infty} + \eta_h F_{\eta,\infty} =: \nu_h \cdot F_\infty, \tag{4.2.30}$$

for three real constants ν_h, ζ_h, η_h , called the running coupling constants at scale h . The factor 2^h in front of ν_h is motivated by the fact that the $F_{\nu,\infty} = (F_{\nu,\infty})_{2,0}$ has scaling dimension $2 - \frac{n}{2} - p|_{(n,p)=(2,0)} = 1$, see (4.2.1), see also (4.3.10) below.

4.3. Trees and Tree Expansions

In this section, we describe the expansion for the kernels of the effective potentials, as it arises from the iterative application of Eq. (4.3). As anticipated above, it is convenient to graphically represent the result of the expansion in terms of GN trees. At the first step, recalling (3.23) and denoting by $N_c(\Psi)$ and $N_m(\Psi)$ the kernels of $\mathcal{N}_c(\phi)$ and $\mathcal{N}_m(\xi)$, respectively, we reorganize the expression for $V_\infty^{(0)}$ obtained by taking the infinite volume of Eq. (4.4) to obtain, for any $\Psi = (\omega, \mathbf{D}, \mathbf{z}) \in \mathcal{M}_\infty$ such that $\mathbf{D} = \mathbf{0}$ (which we identify with the corresponding element (ω, \mathbf{z}) of $\cup_{n \in 2\mathbb{N}}(\{+, -\} \times \Lambda_\infty)^n$),

$$\begin{aligned} V_\infty^{(0)}(\Psi) &= N_c(\Psi) + Z^{-|\Psi|/2} V_\infty^{\text{int}}(\Psi) + \sum_{\substack{\Psi_1 \in \mathcal{M}_{1,\infty}: \\ \Psi_1 \supset \Psi}} Z^{-|\Psi_1|/2} V_\infty^{\text{int}}(\Psi_1) \mathfrak{G}_{\emptyset,\infty}^{(1)}(\Psi_1 \setminus \Psi) \\ &+ \sum_{s=2}^{\infty} \frac{1}{s!} \sum_{\Psi_1, \dots, \Psi_s \in \mathcal{M}_{1,\infty}} \sum_{\Psi} \left(\prod_{j=1}^s [N_m(\Psi_j) + Z^{-|\Psi_j|/2} V_\infty^{\text{int}}(\Psi_j)] \right) \alpha(\Psi; \Psi_1, \dots, \Psi_s) \\ &\times \sum_{T \in \mathcal{S}(\bar{\Psi}_1, \dots, \bar{\Psi}_s)} \mathfrak{G}_{T,\infty}^{(1)}(\bar{\Psi}_1, \dots, \bar{\Psi}_s), \end{aligned} \tag{4.3.1}$$

where V_∞^{int} is the infinite volume limit of the kernel of $\mathcal{V}^{\text{int}}(\phi, \xi) := \mathcal{V}^{\text{int}}(\phi, \xi, \mathbf{0})$, and we recall that in the first (resp. last) line, the factor $\mathfrak{G}_{T,\infty}^{(1)}(\Psi_1 \setminus \Psi)$ (resp. $\mathfrak{G}_{T,\infty}^{(1)}(\bar{\Psi}_1, \dots, \bar{\Psi}_s)$) is different from zero only if $\Psi_1 \setminus \Psi$ is an element (resp. $\bar{\Psi}_1, \dots, \bar{\Psi}_s$ are elements) of $\cup_{n \in 2\mathbb{N}}(\{+i, -i\} \times \Lambda_\infty)^n$. For the definition of $\bar{\Psi}_j$, see the first item in the list after Eq. (4.4).

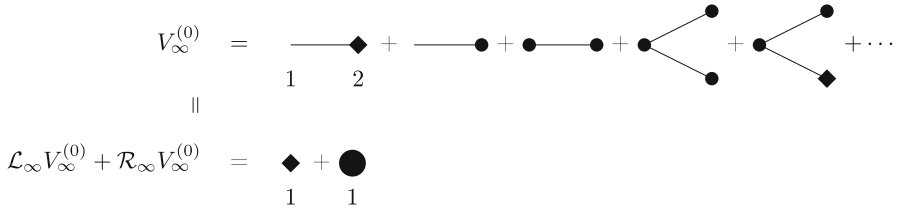


FIGURE 1. Graphical interpretation of (4.3.1)

Similarly, at the following steps, for any $h \leq 0$ and any $\Psi \in \mathcal{M}_\infty$, we can write $V_\infty^{(h-1)}(\Psi)$, as computed via the infinite volume analogue of (4.3), as follows:

$$\begin{aligned}
 V_\infty^{(h-1)}(\Psi) &\sim v_h \cdot F_\infty(\Psi) + \mathcal{R}_\infty V_\infty^{(h)}(\Psi) + \sum_{\substack{\Psi_1 \in \mathcal{M}_{1,\infty}: \\ \Psi_1 \supset \Psi}} \mathcal{R}_\infty V_\infty^{(h)}(\Psi_1) \mathfrak{G}_{\emptyset,\infty}^{(h)}(\Psi_1 \setminus \Psi) \\
 &+ \sum_{s=2}^\infty \frac{1}{s!} \sum_{\Psi_1, \dots, \Psi_s \in \mathcal{M}_{1,\infty}}^{(\Psi)} \left(\prod_{j=1}^s [v_h \cdot F_\infty(\Psi_j) + \mathcal{R}_\infty V_\infty^{(h)}(\Psi_j)] \right) \alpha(\Psi; \Psi_1, \dots, \Psi_s) \\
 &\times \sum_{T \in \mathcal{S}(\bar{\Psi}_1, \dots, \bar{\Psi}_s)} \mathfrak{G}_{T,\infty}^{(h)}(\bar{\Psi}_1, \dots, \bar{\Psi}_s), \tag{4.3.2}
 \end{aligned}$$

where $\mathfrak{G}_{T,\infty}^{(h)}(\bar{\Psi}_1, \dots, \bar{\Psi}_s)$ is the infinite volume analogue of the function defined in (4.5), differing from it for an important feature (besides the ‘obvious’ one that $\mathfrak{G}_{T,\infty}^{(h)}$ is defined in terms of the infinite plane propagators $g_{\ell,\infty}^{(h)}$ rather than those on the cylinder): since now the field multilabels Ψ_i have the form $(\omega_i, \mathbf{D}_i, z_i)$, with \mathbf{D}_i different from $\mathbf{0}$, in general, the infinite plane propagators $g_{\ell,\infty}^{(h)}$, with $\ell = ((\omega_i, D_i, z_i), (\omega_j, D_j, z_j))$, entering the definition of $\mathfrak{G}_{T,\infty}^{(h)}$ should now be interpreted as $\partial_{z_i}^{D_i} \partial_{z_j}^{D_j} g_{\omega_i \omega_j}^{(h)}(z_i, z_j)$.

We graphically interpret (4.3.1) as in Fig. 1. On the right-hand side of the first line, we have drawn a series of diagrams consisting of a root at scale 1, which we will usually denote v_0 , connected to s other vertices (which we will call *endpoints*) at scale 2 which are of two different types: \blacklozenge , which we call *counterterm endpoints* and which represent N_c or N_m , and \bullet , called *interaction endpoints* and representing V_∞^{int} . In the first two terms (in which there is no factor $\mathfrak{G}_{T,\infty}^{(1)}$, because it is ‘trivial’, i.e., it equals $\mathfrak{G}_{\emptyset,\infty}^{(1)}(\emptyset) = 1$) the root is drawn simply as the end of a line segment (we will say it is *undotted*); in the other terms (including all those with $s > 1$) we draw a dot \bullet to indicate the presence of a non-trivial $\mathfrak{G}_{T,\infty}^{(1)}$ and α factors and additional sums.

In order to iterate the scheme, we decompose $V_\infty^{(0)}$ as $V_\infty^{(0)} \sim \mathcal{L}_\infty V_\infty^{(0)} + \mathcal{R}_\infty V_\infty^{(0)}$ and graphically represent $\mathcal{L}_\infty V_\infty^{(0)} = v_0 \cdot F_\infty$ by a counterterm vertex \blacklozenge at scale 1, and $\mathcal{R}_\infty V_\infty^{(0)}$ by a big vertex \bullet at scale 1, as indicated in the second line of Fig. 1. Next, using the conventions of Fig. 1, we graphically represent $V_\infty^{(-1)}$, computed by (4.3.2) with $h = 0$, as described in Fig. 2.

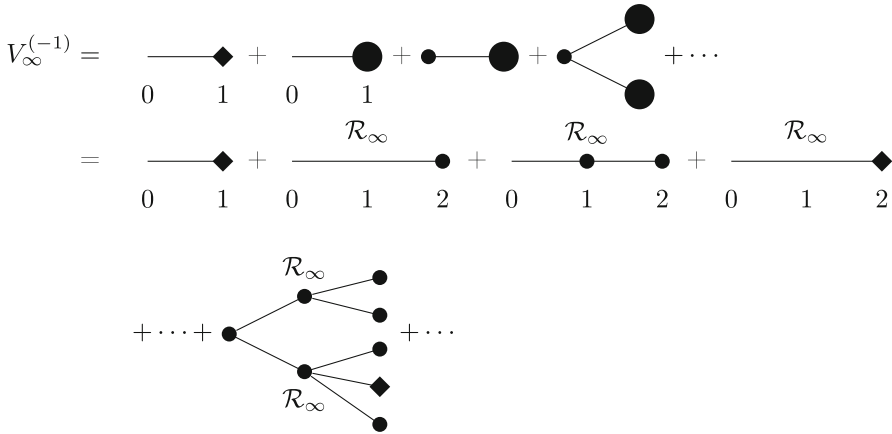


FIGURE 2. Graphical interpretation of (4.3.2) with $h = 0$

In passing from the first to the second line of Fig. 2 we expanded the big vertex on scale 1, which represents $\mathcal{R}_\infty V_\infty^{(0)}$, by using the first line of Fig. 1, with an additional label \mathcal{R}_∞ on the vertices on scale 1, to represent the action of \mathcal{R}_∞ .

The graphical equations in Figs. 1 and 2 are the analogues of the graphical equations in [15, Figures 6-7], which contains a more detailed discussion of some aspects of this construction. By iterating the same kind of graphical equations on lower scales, expanding the big vertices \bullet until we are left with endpoints all of type \blacklozenge or \bullet , we find that $V_\infty^{(h)}$ can be graphically expanded in terms of trees of the kind depicted in Fig. 3, with the understanding that in principle there should be a label \mathcal{R}_∞ at all the intersections of the branches with the vertical lines, with the sole exception of v_0 ; however, by convention, in order not to overwhelm the figures, we prefer not to indicate these labels explicitly. We call such trees ‘GN trees’ and denote by $\mathcal{T}_\infty^{(h)}$, with $h \leq 0$, the set of GN trees with root v_0 on scale $h + 1$. We call ‘vertices’ of a GN tree the root v_0 , all its dotted nodes, and its endpoints.

We introduce some conventions and observations about the set of GN trees:

- The root v_0 is the unique leftmost vertex of the tree. Its degree (number of incident edges) must be at least 1, i.e., v_0 cannot be an endpoint. It may or may not be dotted; in order for v_0 not to be dotted, its degree must be 1.
- Vertices, other than the root, with exactly one successor, are called ‘trivial’.
- Interaction endpoints \bullet can only be on scale 2. Counterterm endpoints \blacklozenge can be on all scales ≤ 2 ; if such an endpoint is on a scale $h < 2$, then it must be connected to a non-trivial vertex on scale $h - 1$. [The reason is the following: if this were not case, then there would be an \mathcal{R}_∞ operator

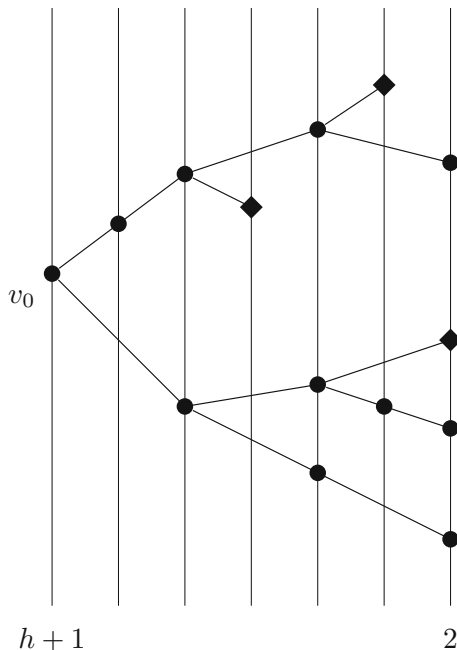


FIGURE 3. Example of a tree in $\mathcal{T}_\infty^{(h)}$. As explained in the text, one should imagine that a label \mathcal{R} , indicating an action of the \mathcal{R} operator, is present at all the intersections of the branches with the vertical lines, with the only exception of v_0 . In order not to overwhelm the figures, these labels are left implicit

acting on the value of the endpoint, but this would annihilate it, because a \blacklozenge endpoint on scale $h < 2$ corresponds to $\mathcal{L}_\infty V_\infty^{(h)}$, and the definitions of $\mathcal{L}_\infty, \mathcal{R}_\infty$ are such that $\mathcal{R}_\infty(\mathcal{L}_\infty V_\infty^{(h)}) \sim 0$, see Remark 4.2.]

In terms of these trees, we shall write the expansion for $V_\infty^{(h)} = V_\infty^{(h)}[\underline{v}]$, thought of as a function of $\underline{v} := \{(\nu_h, \zeta_h, \eta_h)\}_{h \leq 0}$, as

$$V_\infty^{(h)}[\underline{v}] \sim \sum_{\tau \in \mathcal{T}_\infty^{(h)}} W_\infty[\underline{v}; \tau]. \tag{4.3.3}$$

In order to write $W_\infty[\underline{v}; \tau]$ more explicitly, we need to specify some additional notations and conventions about GN trees. Let $\tau \in \mathcal{T}_\infty := \cup_{h \leq 0} \mathcal{T}_\infty^{(h)}$ and $v_0 = v_0(\tau)$ be its root. Then:

- We let $V(\tau)$ be the set of vertices, $V_e(\tau) \subset V(\tau)$ the set of endpoints, and $V_0(\tau) := V(\tau) \setminus V_e(\tau)$. We also let $V'(\tau) := V(\tau) \setminus \{v_0\}$ and $V'_0(\tau) := V_0(\tau) \setminus \{v_0\}$.
- Given $v \in V(\tau)$, we let h_v be its scale.

- $v \geq w$ or ‘ v is a successor of w ’ means that the (unique) path from v to v_0 passes through w . Obviously, $v > w$ means that v is a successor of w and $v \neq w$.
- ‘ v is an immediate successor of w ’, denoted $v \triangleright w$, means that $v \geq w$, $v \neq w$, and v and w are directly connected. For any $v \in \tau$, S_v is the set of $w \in \tau$ such that $w \triangleright v$.
- For any $v > v_0$, we denote by v' the unique vertex such that $v \triangleright v'$.
- Subtrees: for each $v \in V_0(\tau)$, let $\tau_v \in \mathcal{T}_\infty^{(h_v-1)}$ denote the subtree consisting of the vertices with $w \geq v$.

Next, we need to attach labels to their vertices, in order to distinguish the various contributions to the kernels arising from the different choices of the sets Ψ_i , etc., in (4.3.1), (4.3.2), also keeping track of the order in which they appear. In particular, with each $v \in V(\tau)$ we associate a set P_v of *field labels*, sometimes called the set of *external fields*, whose elements carry two informations: their position within an ordered list which they belong to, and their ω index; more precisely, the family $\underline{P} = \{P_v\}_{v \in V(\tau)}$ is characterized by the following properties, which correspond to properties of the iteration of the kernel:

- $|P_v|$ is always even and positive. If v is a \blacklozenge endpoint, then $|P_v| = 2$.
- If v is an endpoint of τ , then P_v has the form $\{(j, 1, \omega_1), \dots, (j, 2n, \omega_{2n})\}$, where j is the position of v in the ordered list of endpoints, and $\omega_l \in \{+, -, i, -i\}$, if $h_v = 2$, while $\omega_l \in \{+, -\}$, if $h_v < 2$. Given $f = (j, l, \omega_l)$, we let $o(f) = (j, l)$ and $\omega(f) := \omega_l$.
- If v is not an endpoint, $P_v \subset \bigcup_{w \in S_v} P_w$.
- If $v \in V_0(\tau)$, we let $Q_v := (\bigcup_{w \in S_v} P_w) \setminus P_v$ be the set of *contracted fields*. If v is dotted, then we require $|Q_v| \geq 2$ and $|Q_v| \geq 2(|S_v| - 1)$; and conversely Q_v is empty if and only if $v = v_0$ and v_0 is not dotted.
- If $h_v = 1$ and v is not an endpoint, then $Q_v = \bigcup_{w \in S_v} \{f \in P_w \mid \omega(f) \in \{+i, -i\}\}$ (all and only massive fields are integrated on scale 1).

For $\tau \in \mathcal{T}_\infty$, we denote by $\mathcal{P}(\tau)$ be the set of allowed $\underline{P} = \{P_v\}_{v \in V(\tau)}$. We also denote by ω_v the tuple of components $\omega(f)$, with $f \in P_v$, and by $\omega_v|_Q$ the restriction of ω_v to any subset $Q \subseteq P_v$. Note that the definitions imply that for $v, w \in \tau$ such that neither $v \geq w$ or $v \leq w$ (for example when $v' = w'$ but $v \neq w$), P_v and P_w are disjoint, as are Q_v and Q_w .

Next, given $\underline{P} \in \mathcal{P}(\tau)$, for all $v \in V_0(\tau)$ we define sets T_v ,

$$T_v = \{(f_1, f_2), \dots, (f_{2|S_v|-3}, f_{2|S_v|-2})\} \subset Q_v^2,$$

called *spanning trees* associated with v , characterized by the following properties: if $w(f)$ denote the (unique) $w \in S_v$ for which $f \in P_w$, then $(f, f') \in T_v \Rightarrow w(f) \neq w(f')$ and $o(f) < o(f')$; moreover, $\{\{w(f_1), w(f_2)\}, \dots, \{w(f_{2|S_v|-3}), w(f_{2|S_v|-2})\}\}$ is the edge set of a tree with vertex set S_v . We denote by $\mathcal{S}(\tau, \underline{P})$ the set of allowed $\underline{T} = \{T_v\}_{v \in V_0(\tau)}$.

Finally, for each $v \in V(\tau)$, we denote by D_v a map $D_v : P_v \rightarrow \mathcal{D} = \{D \in \{0, 1, 2\}^2 : \|D\|_1 \leq 2\}$; the reader should think that a derivative operator $\partial^{D_v(f)}$ acts on the field labeled f . We denote by $\mathcal{D}(\tau, \underline{P})$ the set of families

of maps $\underline{D} = \{D_v\}_{v \in V(\tau)}$. We also denote by \mathbf{D}_v the tuple of components $D_v(f)$, with $f \in P_v$, and by $\mathbf{D}_v|_Q$ the restriction of \mathbf{D}_v to any subset $Q \subseteq P_v$. Additionally, if a map $z : P_v \rightarrow \Lambda_\infty$ is assigned, we denote by \mathbf{z}_v the tuple of components $z(f)$, with $f \in P_v$, and by $\Psi_v = \Psi(P_v) := (\omega_v, \mathbf{D}_v, \mathbf{z}_v)$ the field multilabel associated with $\omega_v, \mathbf{D}_v, \mathbf{z}_v$; moreover, if $v \in V_0(\tau)$ and also the maps $z : P_w \rightarrow \Lambda_\infty$, for all $w \in S_v$, are assigned, for each $w \in S_v$ we denote by $\bar{\Psi}_w = \Psi(P_w \setminus P_v) = (\omega_w|_{P_w \setminus P_v}, \mathbf{D}_w|_{P_w \setminus P_v}, \mathbf{z}_w|_{P_w \setminus P_v})$ the restriction of Ψ_w to $P_w \setminus P_v$ (here $\mathbf{z}_w|_Q$ is the restriction of \mathbf{z}_w to the subset $Q \subset P_w$).

In terms of these definitions, we write $W_\infty[\underline{v}; \tau]$ in the right side of (4.3.3) as

$$W_\infty[\underline{v}; \tau] = \sum_{P \in \mathcal{P}(\tau)} \sum_{T \in \mathcal{S}(\tau, P)} \sum_{D \in \mathcal{D}(\tau, P)} W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}], \tag{4.3.4}$$

where $W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]$ is the translationally invariant kernel inductively defined as follows: letting $\mathbf{D}'_{v_0} := \bigoplus_{v \in S_{v_0}} \mathbf{D}_v|_{P_{v_0}}$ for $h_{v_0} < 1$ and $\mathbf{D}'_{v_0} := \mathbf{0}$ for $h_{v_0} = 1$,

$$W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}](\omega_0, \mathbf{D}_0, \mathbf{z}_0) = \mathbb{1}(\omega_0 = \omega_{v_0}) \mathbb{1}(\mathbf{D}_0 = \mathbf{D}_{v_0} = \mathbf{D}'_{v_0}) \frac{\alpha_{v_0}}{|S_{v_0}|!} \times \sum_{\substack{z: P_{v_0} \cup Q_{v_0} \rightarrow \Lambda_\infty \\ \mathbf{z}_0 = \mathbf{z}_{v_0}}} \mathfrak{G}_{T_{v_0}, \infty}^{(h_{v_0})}(\bar{\Psi}_{v_1}, \dots, \bar{\Psi}_{v_{s_{v_0}}}) \prod_{v \in S_{v_0}} K_{v, \infty}(\Psi_v), \tag{4.3.5}$$

where $\alpha_{v_0} = \alpha(\Psi_{v_0}; \Psi_{v_1}, \dots, \Psi_{v_{s_{v_0}}})$, cf. (4.3.1), and we recall that, if $|S_{v_0}| = 1$, then $T_{v_0} = \emptyset$. In this case, if $\Psi_{v_1} = \Psi_{v_0}$, then $\mathfrak{G}_{\emptyset, \infty}^{(h_{v_0})}(\emptyset)$ should be interpreted as being equal to 1; this latter case is the one in which, graphically, v_0 is not dotted. In the second line of (4.3.5), if $h_{v_0} = 1$,

$$K_{v, \infty}(\Psi_v) = \begin{cases} N_c(\Psi) & \text{if } v \text{ is of type } \blacklozenge \text{ and } v_0 \text{ is undotted,} \\ N_m(\Psi_v) & \text{if } v \text{ is of type } \blacklozenge \text{ and } v_0 \text{ is dotted,} \\ Z^{-|\Psi_v|/2} V_\infty^{\text{int}}(\Psi_v) & \text{if } v \text{ is of type } \bullet, \end{cases} \tag{4.3.6}$$

while, if $h_{v_0} < 1$,

$$K_{v, \infty}(\Psi_v) := \begin{cases} v_{h_{v_0}} \cdot F_\infty(\Psi_v) & \text{if } v \in V_e(\tau) \text{ is of type } \blacklozenge \text{ and } h_v = h_{v_0} + 1 \\ \mathcal{R}_\infty N_c(\Psi_v) & \text{if } v \in V_e(\tau) \text{ is of type } \blacklozenge \text{ and } h_v = 2 \\ Z^{-|\Psi_v|/2} \mathcal{R}_\infty V_\infty^{\text{int}}(\Psi_v) & \text{if } v \in V_e(\tau) \text{ is of type } \bullet \\ \bar{W}_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}_v](\Psi_v) & \text{if } v \in V_0(\tau), \end{cases} \tag{4.3.7}$$

where, in the last line of (4.3.7), letting \underline{P}_v (resp. \underline{T}_v , resp. \underline{D}_v) be the restriction of \underline{P} (resp. \underline{T} , resp. \underline{D}) to the subtree τ_v , and $\underline{D}'_v := \{D'_w\} \cup \{D_w\}_{w \in V(\tau): w > v_0}$ (here D'_v is the map such that $\mathbf{D}'_v := \bigoplus_{w \in S_v} \mathbf{D}_w|_{P_v}$), we denoted

$$\bar{W}_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}_v] := \mathcal{R}_\infty W_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}'_v] \Big|_{\mathbf{D}_v}, \tag{4.3.8}$$

and we recall that the definition of $\mathcal{R}_\infty V|_{\underline{D}}$ was given a few lines after (4.2.20). The inductive proof that $V_\infty^{(h)} = V_\infty^{(h)}[\underline{v}]$, as iteratively computed by (4.3.1)–(4.3.2), is equivalent to $\sum_{\tau \in \mathcal{T}_\infty^{(h)}} W_\infty[\underline{v}; \tau]$, with $W_\infty[\underline{v}; \tau]$ as in (4.3.4), (4.3.5), is straightforward and left to the reader.

Remark 4.4. Given $\tau \in \mathcal{T}_\infty$ and $\underline{P} \in \mathcal{P}(\tau)$, we say that \underline{D} is ‘allowed’ if $W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}] \not\approx 0$. With some abuse of notation, from now on we will redefine $\mathcal{D}(\tau, \underline{P})$ to be the set of allowed \underline{D} for a given τ and \underline{P} . Of course, such a redefinition has no impact on the validity of (4.3.4). If \underline{D} is allowed, then it must satisfy a number of constraints. In particular, if $v \in V_0(\tau)$, $w \in V_e(\tau)$ and $f \in P_v \cap P_w$, then $D_v(f) \geq D_w(f)$. Moreover, letting, for any $v \in V_0(\tau)$, $R_v := \|\mathbf{D}_v\|_1 - \sum_{w \in S_v} \|\mathbf{D}_w|_{P_v}\|_1 \equiv \|\mathbf{D}_v\|_1 - \|\mathbf{D}'_v\|_1$, one has

$$R_v = \begin{cases} 2, & |P_v| = 2 \text{ and } \|\mathbf{D}'_v\|_1 = 0 \\ 1, & |P_v| = 2 \text{ and } \|\mathbf{D}'_v\|_1 = 1 \text{ or } |P_v| = 4 \text{ and } \|\mathbf{D}'_v\|_1 = 0 \\ 0, & \text{otherwise,} \end{cases} \tag{4.3.9}$$

with the exception of v_0 , for which $R_{v_0} \equiv 0$ (in other words $\mathbf{D}_{v_0} = \mathbf{D}'_{v_0}$, see (4.3.5)). Finally, the combination

$$d(P_v, \mathbf{D}_v) := 2 - \frac{|P_v|}{2} - \|\mathbf{D}_v\|_1, \tag{4.3.10}$$

known as the **scaling dimension** of v , is ≤ -1 for all $v \in V'_0(\tau)$, and for all $v \in V_e(\tau)$ such that $h_v = 2$ and $h_{v'} < 1$. As we shall see below, see in particular the statement of Lemma 4.7, the fact that $d(P_v, \mathbf{D}_v) \leq -1$ for all such vertices guarantees that the expansion in GN trees is convergent uniformly in h_{v_0} .

Remark 4.5. With the other arguments fixed, the number of choices of \underline{D} for which $W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]$ does not vanish is no more than $10^{|\mathcal{V}(\tau)|}$: there is a choice of at most 10 possible values for each endpoint¹⁰, and then the other values are fixed except for a choice of up to 10 possibilities each time that \mathcal{R}_∞ acts non-trivially on a vertex $v \in V'_0(\tau)$, i.e., each time that, for such a vertex, $R_v > 0$, see (4.3.9).

4.4. Bounds on the Kernels of the Full Plane Effective Potentials

In this subsection we show that the norm of the kernels $W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]$ is summable over $\tau, \underline{P}, \underline{T}, \underline{D}$, provided that the elements of the sequence \underline{v} are bounded and sufficiently small. We measure the size of $W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]$ in terms of the weighted norm (4.2.21), with $\kappa = \frac{c_0}{2} 2^{h_{v_0}}$, where h_{v_0} is the scale of

¹⁰The value 10 bounds the number of different terms that the operator \mathcal{R}_∞ produces when it acts non-trivially on an interaction endpoints. In fact, the cases in which \mathcal{R}_∞ acts non-trivially are those listed in the right side of (4.2.20) with $(n, p) = (2, 2), (4, 1)$. If $(n, p) = (2, 2)$, the number of possible values taken by D_v is 10 (one derivative in direction $i \in \{1, 2\}$ on the first Grassmann field and one derivative in direction $j \in \{1, 2\}$ on the second Grassmann field, etc.); if $(n, p) = (4, 1)$, the number of possible values taken by D_v is 8, which is smaller than 10 (one derivative in direction $i \in \{1, 2\}$ on the k -the Grassmann field, with $k \in \{1, 2, 3, 4\}$).

the root of τ and c_0 is the minimum between the constant c in Proposition 2.3 and half the constant c in (3.14), (3.15), (3.20), (3.21). With some abuse of notation, we let

$$\|W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]\|_{h_{v_0}} := \|W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]\|_{(\frac{c_0}{2} 2^{h_{v_0}})}. \tag{4.4.1}$$

The first, basic, bound on the kernels $W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]$ is provided by the following proposition. We recall that we assumed once and for all that $Z \in U$, with $U = \{z \in \mathbb{R} : |z - 1| \leq 1/2\}$, and that $t_1^*, t_1, t_2 \in K'$, with K' the compact set defined before the statement of Theorem 1.1.

Proposition 4.6. *Let $W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]$ be inductively defined as in (4.3.5). There exist $C, \kappa, \lambda_0 > 0$ such that, for any $\tau \in \mathcal{T}_\infty$, $\underline{P} \in \mathcal{P}(\tau)$, $\underline{T} \in \mathcal{S}(\tau, \underline{P})$, $\underline{D} \in \mathcal{D}(\tau, \underline{P})$, and $|\lambda| \leq \lambda_0$,*

$$\begin{aligned} & \|W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]\|_{h_{v_0}} \\ & \leq C^{\sum_{v \in V_c(\tau)} |P_v|} \left(\prod_{v \in V_0(\tau)} \frac{2^{(\frac{1}{2}|Q_v| + \sum_{w \in S_v} \|D_w|_{Q_v}\|_{1-R_v+2-2|S_v|)h_v}}}{|S_v|!} \right) \\ & \times \prod_{v \in V_c(\tau)} \begin{cases} 2^{(h_v-1)(2-\frac{1}{2}|P_v|-\|D_v\|_1)} \epsilon_{h_v-1} & \text{if } v \text{ is of type } \blacklozenge \\ |\lambda|^{\max\{1, \kappa|P_v\}} & \text{if } v \text{ is of type } \bullet \end{cases} \end{aligned} \tag{4.4.2}$$

where $\epsilon_h := \max\{|\nu_h|, |\zeta_h|, |\eta_h|\}$ if $h \leq 0$ and $\epsilon_1 := \max\{\|N_c\|_2, \|N_m\|_2\}$.

Proof. Let us first consider the case $h_{v_0} = 1$, in which case, using (4.3.5) and (4.3.6), we find

$$\begin{aligned} & \|W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{0}]\|_{h_{v_0}} \leq \frac{1}{|S_{v_0}|!} \sum_{\substack{z: \cup_{v \in S_{v_0}} P_v \rightarrow \Lambda_\infty \\ z(f_1) \text{ fixed}}} \\ & \times e^{c_0 \delta(z_{v_0})} |\mathfrak{G}_{T_{v_0}, \infty}^{(1)}(\bar{\Psi}_{v_1}, \dots, \bar{\Psi}_{v_{s_{v_0}}})| \cdot \\ & \cdot \prod_{v \in S_{v_0}} \begin{cases} |N_c(\Psi_v)| & \text{if } v \text{ is of type } \blacklozenge \text{ and } v_0 \text{ is undotted,} \\ |N_m(\Psi_v)| & \text{if } v \text{ is of type } \blacklozenge \text{ and } v_0 \text{ is dotted,} \\ |Z|^{-|\Psi_v|/2} |V_\infty^{\text{int}}(\Psi_v)| & \text{if } v \text{ is of type } \bullet, \end{cases} \end{aligned} \tag{4.4.3}$$

where f_1 is the first element of P_{v_0} . By using the definition of $\mathfrak{G}_{T, \infty}^{(1)}$ and the property $(\text{Pf}M)^2 = \det M$, valid for any antisymmetric matrix M , we find

$$|\mathfrak{G}_{T, \infty}^{(1)}(Q_1, \dots, Q_s)| \leq \left(\prod_{\ell \in T} |g_\ell^{(1)}| \right) \sup_{\mathbf{t}} \sqrt{|\det G_{Q_1, \dots, Q_s, T, \infty}^{(1)}(\mathbf{t})|}, \tag{4.4.4}$$

so that, thanks to items 1,3,4 of Proposition 2.3 (which apply to $\mathfrak{g}_\infty^{(1)}$ by Remark 2.6), and to the Gram-Hadamard inequality [15, Appendix A.3], which allows one to bound the determinant of any matrix M with elements

$M_{i,j} = (\gamma_i, \tilde{\gamma}_j)$ as $|\det M| \leq \prod_i |\gamma_i| |\tilde{\gamma}_i|$,

$$\begin{aligned}
 |\mathfrak{G}_{T,\infty}^{(1)}(Q_1, \dots, Q_s)| &\leq C^{s-1} \left(\prod_{(f,f') \in T} e^{-2c_0 \|z(f) - z(f')\|_1} \right) \\
 &\quad \times \left(\prod_{f \in \cup_i Q_i \setminus T} |\gamma_{\omega(f),0,z(f)}^{(1)}| \cdot |\tilde{\gamma}_{\omega(f),0,z(f)}^{(1)}| \right)^{1/2} \\
 &\leq (C')^{\sum_i |Q_i|} \left(\prod_{(f,f') \in T} e^{-2c_0 \|z(f) - z(f')\|_1} \right). \tag{4.4.5}
 \end{aligned}$$

If we now note that

$$\delta(\mathbf{z}_0) \leq \sum_{(f,f') \in T_{v_0}} \|z(f) - z(f')\|_1 + \sum_{v \in S_{v_0}} \delta(\mathbf{z}_v), \tag{4.4.6}$$

and plug Eq. (4.4.5) into (4.4.3), we find

$$\begin{aligned}
 \|W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{0}]\|_{h_{v_0}} &\leq \frac{C^{\sum_{v \in S_{v_0}} |P_v|}}{|S_{v_0}|!} \left(\sum_{z \in \Lambda_\infty} e^{-c_0 \|z\|_1} \right)^{|S_{v_0}|-1} \\
 &\quad \times \prod_{v \in S_{v_0}} \begin{cases} \|(N_c)_2\|_{(c_0)} & \text{if } v \text{ is of type } \blacklozenge \text{ and } v_0 \text{ is undotted,} \\ \|(N_m)_2\|_{(c_0)} & \text{if } v \text{ is of type } \blacklozenge \text{ and } v_0 \text{ is dotted,} \\ \|(V_\infty^{\text{int}})_{|P_v|}\|_{(c_0)} & \text{if } v \text{ is of type } \bullet, \end{cases} \tag{4.4.7}
 \end{aligned}$$

which immediately implies the desired bound for $h_{v_0} = 1$, because of the definition of ϵ_1 and $\|(V_\infty^{(1)})_{|P_v|}\|_{(c_0)} \leq C^{|P_v|} |\lambda|^{\max\{1, \kappa|P_v|\}}$, see (3.15) and recall the definition of c_0 at the beginning of this subsection.

Next, we consider the case $h_{v_0} < 1$, in which case $K_{v,\infty}$ is defined by (4.3.7). We proceed similarly: we start from (4.3.5) and use the analogue of (4.4.4)–(4.4.5), namely

$$\begin{aligned}
 |\mathfrak{G}_{T,\infty}^{(h_{v_0})}(Q_1, \dots, Q_s)| &\leq \left(\prod_{(f,f') \in T} |g_{\ell,\infty}^{(h_{v_0})}| \right) \\
 &\quad \left(\prod_{f \in \cup_i Q_i \setminus T} |\gamma_{\omega(f),D(f),z(f)}^{(h_{v_0})}| \cdot |\tilde{\gamma}_{\omega(f),D(f),z(f)}^{(h_{v_0})}| \right)^{1/2} \\
 &\leq (C2^{h_{v_0}})^{\frac{1}{2} \sum_i |Q_i|} 2^{h_{v_0} \sum_{f \in \cup_i Q_i} \|D(f)\|_1} \left(\prod_{(f,f') \in T} e^{-c_0 2^{h_{v_0}} \|z(f) - z(f')\|_1} \right), \tag{4.4.8}
 \end{aligned}$$

where in the second inequality we used again the bounds in items 1,4 of Prop.2.3. Using (4.4.8) and, again, (4.4.6), we obtain the analogue of (4.4.7),

$$\|W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]\|_{h_{v_0}}$$

$$\leq \frac{1}{|S_{v_0}|!} (C2^{h_{v_0}})^{\frac{|Q_{v_0}|}{2}} 2^{h_{v_0} \sum_{v \in S_{v_0}} \|D_v\|_{Q_{v_0}} \|1\|} \cdot \left(\sum_{z \in \Lambda_\infty} e^{-\frac{c_0}{2} 2^{h_{v_0}} \|z\|_1} \right)^{|S_{v_0}|-1} \cdot \prod_{v \in S_{v_0}} \begin{cases} \|(v_{h_{v_0}} \cdot F_\infty)_{2, \|D_v\|_1}\|_{h_{v_0}} & \text{if } v \in V_e(\tau) \text{ is of type } \blacklozenge \text{ and } h_v = h_{v_0} + 1 \\ \|(\mathcal{R}_\infty N_c)_{2, \|D_v\|_1}\|_{h_{v_0}} & \text{if } v \in V_e(\tau) \text{ is of type } \blacklozenge \text{ and } h_v \neq h_{v_0} + 1 \\ |Z|^{-|\Psi_v|/2} \|(\mathcal{R}_\infty V_\infty^{\text{int}})_{|P_v|, \|D_v\|_1}\|_{h_{v_0}} & \text{if } v \in V_e(\tau) \text{ is of type } \bullet \\ \|(\overline{W}_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}_v])_{|P_v|, \|D_v\|_1}\|_{h_{v_0}} & \text{if } v \in V_0(\tau), \end{cases} \tag{4.4.9}$$

The terms in the second line can be bounded as follows:

- If $v \in V_e(\tau)$ is of type \blacklozenge and $h_v = h_{v_0} + 1$,

$$\begin{aligned} \|(v_{h_{v_0}} \cdot F_\infty)_{2, \|D_v\|_1}\|_{h_{v_0}} &\leq C \begin{cases} 2^{h_{v_0}} |\nu_{h_{v_0}}| & \text{if } \|D_v\|_1 = 0 \\ \max\{|\zeta_{h_{v_0}}|, |\eta_{h_{v_0}}|\} & \text{if } \|D_v\|_1 = 1 \end{cases} \\ &\leq C' 2^{(h_v-1)(2-\frac{|P_v|}{2}-\|D_v\|_1)} \epsilon_{h_v-1} \end{aligned} \tag{4.4.10}$$

where we recall that $\epsilon_h = \max\{|\nu_h|, |\zeta_h|, |\eta_h|\}$ and, in passing from the first to the second line, we used the fact that $|P_v| = 2$ and $h_{v_0} = h_v - 1$, so that $2^{(h_v-1)(2-\frac{|P_v|}{2}-\|D_v\|_1)}$ is equal to $2^{h_{v_0}}$, if $\|D_v\|_1 = 0$, and is equal to 1, if $\|D_v\|_1 = 1$.

- If $v \in V_e(\tau)$ is of type \blacklozenge and $h_v \neq h_{v_0} + 1$, using Lemma 4.3 and the definition of ϵ_1 we have

$$\begin{aligned} \|(\mathcal{R}_\infty N_c)_{2, \|D_v\|_1}\|_{h_{v_0}} &\leq \|(\mathcal{R}_\infty N_c)_{2, \|D_v\|_1}\|_0 \leq 2C \|N_c\|_1 \leq 2C \epsilon_1 \\ &= C 2^{(h_v-1)(2-\frac{|P_v|}{2}-\|D_v\|_1)} \epsilon_1, \end{aligned}$$

where the last equality holds trivially since we necessarily have $h_v = 2$ and $|P_v| = 2$, and $(\mathcal{R}_\infty N_c)_{2, \|D_v\|_1}$ vanishes unless $\|D_v\|_1 = 2$.

- If $v \in V_e(\tau)$ is of type \bullet (and, therefore, $h_v = 2$), then

$$\|(\mathcal{R}_\infty V_\infty^{\text{int}})_{|P_v|, \|D_v\|_1}\|_{h_{v_0}} \leq \|(\mathcal{R}_\infty V_\infty^{\text{int}})_{|P_v|, \|D_v\|_1}\|_0 \leq C^{|P_v|} |\lambda|^{\max\{1, \kappa|P_v|\}},$$

thanks to Lemma 4.3 and Eq.(3.15).

- If $v \in V_0(\tau)$, recalling the definition of $\overline{W}_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}_v]$, see Eq. (4.3.8), and the bounds on the norm of \mathcal{R}_∞ , see (4.2.22)–(4.2.25), we find

$$\begin{aligned} &\|(\overline{W}_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}_v])_{|P_v|, \|D_v\|_1}\|_{h_{v_0}} \\ &\leq C 2^{-h_v(\|D_v\|_1 - \|D'_v\|_1)} \| (W_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}'_v])_{|P_v|, \|D'_v\|_1} \|_{h_v}. \end{aligned} \tag{4.4.11}$$

Recalling the definition of D'_v , that is $D'_v = \bigoplus_{w \in S_v} D_w|_{P_v}$, as well as the one of R_v , see the line before (4.3.9), we recognize that $\|D_v\|_1 - \|D'_v\|_1 = R_v$. Note also that $W_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}'_v]$ coincides with its restriction $(W_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}'_v])_{|P_v|, \|D'_v\|_1}$, so that the second line of (4.4.11) can be rewritten more compactly as $C 2^{-h_v R_v} \|W_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}'_v]\|_{h_v}$.

Plugging these bounds in (4.4.9), noting that $\sum_{z \in \Lambda_\infty} e^{-\frac{c_0}{2} 2^{h_{v_0}} \|z\|_1} \leq C 2^{-2h_{v_0}}$, we find

$$\begin{aligned} \|W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]\|_{h_{v_0}} &\leq \frac{1}{|S_{v_0}|} C^{\lfloor \frac{|Q_{v_0}|}{2} \rfloor + |S_{v_0}|} 2^{h_{v_0} (\lfloor \frac{|Q_{v_0}|}{2} \rfloor + \sum_{v \in S_{v_0}} \|D_v\|_{Q_{v_0}} \|1 + 2^{-2|S_{v_0}|})}. \\ \cdot \prod_{v \in S_{v_0}} &\begin{cases} 2^{(h_v - 1)(2 - \lfloor \frac{|P_v|}{2} \rfloor - \|D_v\|_1)} \epsilon_{h_v - 1} & \text{if } v \in V_e(\tau) \text{ is of type } \blacklozenge \\ C^{|P_v|} |\lambda|^{\max\{1, \kappa|P_v|\}} & \text{if } v \in V_e(\tau) \text{ is of type } \bullet \\ 2^{-h_v R_v} \|W_\infty[\underline{v}; \tau_v, \underline{P}_v, \underline{T}_v, \underline{D}'_v]\|_{h_v} & \text{if } v \in V_0(\tau), \end{cases} \end{aligned} \tag{4.4.12}$$

Now, in the last line, for $v \in S_{v_0} \cap V_0(\tau)$, we iterate the bound, and we continue to do so until we reach all the endpoints. By doing so, recalling that $R_{v_0} = 0$, we obtain the desired bound, (4.4.2), provided that

$$C^{\sum_{v \in V_0(\tau)} (\lfloor \frac{|Q_v|}{2} \rfloor + |S_v|)} \leq (C')^{\sum_{v \in V_e(\tau)} |P_v|}. \tag{4.4.13}$$

In order to prove this, note that, for any dotted $v \in V_0(\tau)$, $|S_v| \leq 1 + \lfloor \frac{|Q_v|}{2} \rfloor \leq |Q_v|$, because $|Q_v| \geq \max\{2, 2(|S_v| - 1)\}$; moreover, if v_0 is not dotted, $|S_{v_0}| = 1$ and $|Q_{v_0}| = 0$. Therefore, recalling also that $|Q_v| = \sum_{w \in S_v} |P_w| - |P_v|$, we find

$$C^{\sum_{v \in V_0(\tau)} (\lfloor \frac{|Q_v|}{2} \rfloor + |S_v|)} \leq C^{1 + \frac{3}{2} \sum_{v \in V_0} |Q_v|} \leq C^{1 + \frac{3}{2} \sum_{v \in V_e(\tau)} |P_v|},$$

which implies (4.4.13). □

Next, we rearrange (4.4.2) in a different form, more suitable for summing over GN trees and their labels.

Lemma 4.7. *Under the same assumptions as Proposition 4.6,*

$$\|W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]\|_{h_{v_0}} \leq C^{\sum_{v \in V_e(\tau)} |P_v|} \frac{1}{|S_{v_0}|} 2^{h_{v_0} d(P_{v_0}, D_{v_0})} \tag{4.4.14}$$

$$\times \left(\prod_{v \in V'(\tau)} \frac{1}{|S_v|} 2^{(h_v - h_{v'}) d(P_v, D_v)} \right) \prod_{v \in V_e(\tau)} \begin{cases} \epsilon_{h_v - 1} & \text{if } v \text{ is of type } \blacklozenge \\ |\lambda|^{\max\{1, \kappa|P_v|\}} & \text{if } v \text{ is of type } \bullet \end{cases}$$

where $d(P_v, D_v) = 2 - \lfloor \frac{|P_v|}{2} \rfloor - \|D_v\|_1$ is the scaling dimension of v , see (4.3.10).

Note that, as observed in Remark 4.4, the scaling dimensions appearing at exponent in the product over $v \in V'(\tau)$ are all negative, with the exception of the case that v is an endpoint such that $h_{v'} = h_v - 1$. Note, however, that in such a case $2^{(h_v - h_{v'}) d(P_v, D_v)} \leq 2$, which is a constant that can be reabsorbed in $C^{\sum_{v \in V_e(\tau)} |P_v|}$, up to a redefinition of the constant C .

Proof. First note that, for all $v \in V_0(\tau)$, $|Q_v| = \sum_{w \in S_v} |P_w| - |P_v|$, so that, recalling that $R_v = \|D_v\|_1 - \sum_{w \in S_v} \|D_w\|_{P_v}$, we can rewrite the factor

$$\begin{aligned} \prod_{v \in V_0(\tau)} &2^{\left(\frac{1}{2}|Q_v| + \sum_{w \in S_v} \|D_w\|_{Q_v} \|1 - R_v + 2^{-2|S_v|}\right) h_v} \text{ in (4.4.2) as} \\ &2^{h_{v_0} \sum_{v \in V_0(\tau)} \left(\frac{1}{2} \sum_{w \in S_v} |P_w| - \frac{1}{2} |P_v| + \sum_{w \in S_v} \|D_w\|_1 - \|D_v\|_1 - 2(|S_v| - 1)\right)} \\ &\cdot 2^{\sum_{v \in V_0(\tau)} (h_v - h_{v_0}) \left(\frac{1}{2} \sum_{w \in S_v} |P_w| - \frac{1}{2} |P_v| + \sum_{w \in S_v} \|D_w\|_1 - \|D_v\|_1 - 2(|S_v| - 1)\right)}. \end{aligned} \tag{4.4.15}$$

Now, the factor in the first line can be further rewritten by noting that:

$$\sum_{v \in V_0(\tau)} \left(\sum_{w \in S_v} |P_w| - |P_v| \right) = \sum_{v \in V_e(\tau)} |P_v| - |P_{v_0}|, \tag{4.4.16}$$

$$\sum_{v \in V_0(\tau)} \left(\sum_{w \in S_v} \|D_w\|_1 - \|D_v\|_1 \right) = \sum_{v \in V_e(\tau)} \|D_v\|_1 - \|D_{v_0}\|_1, \tag{4.4.17}$$

$$\sum_{v \in V_0(\tau)} (|S_v| - 1) = |V_e(\tau)| - 1. \tag{4.4.18}$$

(The first two identities are ‘obvious’, due to the telescopic structure of the summand; the latter identity can be easily proved by induction.) Therefore,

$$\begin{aligned} & 2^{h_{v_0} \sum_{v \in V_0(\tau)} (\frac{1}{2} \sum_{w \in S_v} |P_w| - \frac{|P_v|}{2} + \sum_{w \in S_v} \|D_w\|_1 - \|D_v\|_1 - 2(|S_{v_0}| - 1))} \\ &= 2^{h_{v_0} (2 - \frac{|P_{v_0}|}{2} - \|D_{v_0}\|_1)} \left(\prod_{v \in V_e(\tau)} 2^{-h_{v_0} (2 - \frac{|P_v|}{2} - \|D_v\|_1)} \right). \end{aligned} \tag{4.4.19}$$

Similarly, the exponent of the factor in the second line of (4.4.15) can be rewritten as

$$\begin{aligned} & \sum_{v \in V_0(\tau)} \left(\sum_{\substack{w \in V_0(\tau) \\ v_0 < w \leq v}} (h_w - h_{w'}) \right) \left(\frac{1}{2} \sum_{w \in S_v} |P_w| - \frac{1}{2} |P_v| \right. \\ & \quad \left. + \sum_{w \in S_v} \|D_w\|_1 - \|D_v\|_1 - 2(|S_{v_0}| - 1) \right) \\ &= \sum_{\substack{w \in V_0(\tau) \\ w > v_0}} (h_w - h_{w'}) \sum_{v \in V_0(\tau_w)} \left(\frac{1}{2} \sum_{w \in S_v} |P_w| - \frac{1}{2} |P_v| \right. \\ & \quad \left. + \sum_{w \in S_v} \|D_w\|_1 - \|D_v\|_1 - 2(|S_{v_0}| - 1) \right) \\ &= \sum_{\substack{w \in V_0(\tau) \\ w > v_0}} (h_w - h_{w'}) \left(2 - \frac{|P_w|}{2} - \|D_w\|_1 - \sum_{v \in V_e(\tau_w)} \left(2 - \frac{|P_v|}{2} - \|D_v\|_1 \right) \right). \end{aligned} \tag{4.4.20}$$

Using (4.4.19) and (4.4.20) and recalling that $2 - \frac{|P_v|}{2} - \|D_v\|_1 \equiv d(P_v, D_v)$, we rewrite (4.4.15) as

$$(4.4.15) = 2^{h_{v_0} d(P_{v_0}, D_{v_0})} \left(\prod_{v \in V'_0(\tau)} 2^{(h_v - h_{v'}) d(P_v, D_v)} \right)$$

$$\begin{aligned} & \left(\prod_{v \in V_e(\tau)} 2^{-(h_{v_0} + \sum_{w \in V_0^<v} (h_w - h_{w'}))d(P_v, \mathbf{D}_v)} \right) \\ &= 2^{h_{v_0}d(P_{v_0}, \mathbf{D}_{v_0})} \left(\prod_{v \in V_0^<(\tau)} 2^{(h_v - h_{v'})d(P_v, \mathbf{D}_v)} \right) \left(\prod_{v \in V_e(\tau)} 2^{-h_{v'}d(P_v, \mathbf{D}_v)} \right). \end{aligned} \tag{4.4.21}$$

By using this rewriting in (4.4.2) and noting that

$$\begin{aligned} & \left(\prod_{v \in V_e(\tau)} 2^{-h_{v'}d(P_v, \mathbf{D}_v)} \cdot \begin{cases} 2^{(h_v - 1)(2 - \frac{1}{2}|P_v| - \|\mathbf{D}_v\|_1)} \epsilon_{h_v - 1} & \text{if } v \text{ is of type } \blacklozenge \\ |\lambda|^{\max\{1, \kappa|P_v|\}} & \text{if } v \text{ is of type } \bullet \end{cases} \right) \\ & \leq \left(\prod_{v \in V_e(\tau)} 2^{(h_v - h_{v'})d(P_v, \mathbf{D}_v)} \cdot \begin{cases} \epsilon_{h_v - 1} & \text{if } v \text{ is of type } \blacklozenge \\ 2^{|P_v|} |\lambda|^{\max\{1, \kappa|P_v|\}} & \text{if } v \text{ is of type } \bullet \end{cases} \right) \end{aligned} \tag{4.4.22}$$

(since if v is an endpoint of type \bullet , then $2^{h_v d(P_v, \mathbf{D}_v)} \geq 2^{-|P_v|}$), we readily obtain the desired estimate, (4.4.14). \square

As announced above, the bound (4.4.14) is written in a form suitable for summing over the GN trees and their labels, as summarized in the following lemma.

Lemma 4.8. *Under the same assumptions as Proposition 4.6, for any $\vartheta \in (0, 1)$, there exists $C_\vartheta > 0$ such that, letting $\mathcal{T}_{\infty; (N, M)}^{(h)}$ denote the subset of $\mathcal{T}_\infty^{(h)}$ whose trees have N endpoints of type \bullet and M endpoints of type \blacklozenge ,*

$$\begin{aligned} & 2^{-\vartheta h} \sum_{\tau \in \mathcal{T}_{\infty; (N, M)}^{(h)}} \sum_{\substack{P \in \mathcal{P}(\tau) \\ |P_{v_0}| = n}} \sum_{\underline{T} \in \mathcal{S}(\tau, P)} \sum_{\substack{D \in \mathcal{D}(\tau, P) \\ \|\mathbf{D}_{v_0}\|_1 = p}} \|W_\infty[\underline{v}; \tau, \underline{P}, \underline{T}, \underline{D}]\|_{h+1} \\ & \leq C_\vartheta^{N+M} |\lambda|^N \left(\max_{h < h' \leq 1} 2^{-\vartheta h'} \epsilon_{h'} \right)^M 2^{h \cdot d(n, p)}. \end{aligned} \tag{4.4.23}$$

Remark 4.9. (Short memory property). The fact that this estimate holds with the factors of $2^{-\vartheta h}$ included indicates that the contribution of trees covering a large range of scales are exponentially suppressed, a behavior referred to in previous works (e.g., [18]) as the ‘short memory property’. As we shall see below, taking advantage of this there is a way to choose the free parameters Z, t_1^*, β such that $|\epsilon_h| \leq K_\vartheta |\lambda| 2^{\vartheta h}$, see Propositions 4.10 and 4.11 below. Under this condition, Lemma 4.8 implies that

$$\begin{aligned}
 & \sum_{\tau \in \mathcal{T}_{\infty; (N, M)}^{(h)}} \sum_{\substack{P \in \mathcal{P}(\tau) \\ |P_{v_0}| = n}} \sum_{\mathcal{I} \in \mathcal{S}(\tau, P)} \sum_{\substack{D \in \mathcal{D}(\tau, P) \\ \|D_{v_0}\|_1 = p}} \|W_{\infty}[\underline{v}; \tau, \underline{P}, \underline{I}, \underline{D}]\|_{h+1} \\
 & \leq C_{\vartheta}^{N+M} |\lambda|^{N+M} 2^{hd(n,p)} 2^{\vartheta h}, \tag{4.4.24}
 \end{aligned}$$

for all $\vartheta < 1$ and $N + M > 0$. The factor $2^{\vartheta h}$ in the right side is called the ‘short-memory factor’.

Summing over $N + M \geq 1$, this immediately implies

$$\left\| (W_{\infty})_{n,p} \right\|_{h+1} \leq C_{\vartheta} |\lambda| 2^{hd(n,p)} 2^{\vartheta h}. \tag{4.4.25}$$

Proof. Thanks to (4.4.14), the left side of (4.4.23) can be bounded from above by

$$\begin{aligned}
 & 2^{h \cdot d(n,p)} |\lambda|^N \left(\max_{h < h' \leq 1} 2^{-\vartheta h'} \epsilon_{h'} \right)^M 2^{-\vartheta h} \sum_{\tau \in \mathcal{T}_{\infty; (N, M)}^{(h)}} \sum_{\substack{P \in \mathcal{P}(\tau) \\ |P_{v_0}| = n}} C^{\sum_{v \in V_{ep}(\tau)} |P_v|} \sum_{\mathcal{I} \in \mathcal{S}(\tau, P)} \\
 & \sum_{\substack{D \in \mathcal{D}(\tau, P) \\ \|D_{v_0}\|_1 = p}} \frac{1}{|S_{v_0}|!} \left(\prod_{v \in V'(\tau)} \frac{1}{|S_v|!} 2^{(h_v - h_{v'})d(P_v, D_v)} \right) \tag{4.4.26} \\
 & \prod_{\substack{v \in V_e(\tau) \\ v \text{ of type } \blacklozenge}} 2^{\vartheta h_v} \prod_{\substack{v \in V_e(\tau) \\ v \text{ of type } \bullet}} |\lambda|^{[\kappa |P_v| - 1]_+},
 \end{aligned}$$

where, in the last factor, $[\cdot]_+$ indicates the positive part. Note that, if either $v \in V'_0(\tau)$ or v is an endpoint such that $h_{v'} < h_v - 1$, then the scaling dimension of v can be bounded uniformly in D_v , i.e., $d(P_v, D_v) \leq \min\{-1, 2 - |P_v|/2\}$; if v is an endpoint such that $h_{v'} = h_v - 1$, then the factor $2^{(h_v - h_{v'})d(P_v, D_v)}$ is smaller than 2 (and, therefore, it can be reabsorbed in $C^{\sum_{v \in V_{ep}(\tau)} |P_v|}$ up to a redefinition of the constant C). Moreover, recall that the number of elements of $\mathcal{D}(\tau, \underline{P})$ is bounded by $10^{|\mathcal{V}(\tau)|}$, see Remark 4.5. Finally, the number of elements of $\mathcal{S}(\tau, \underline{P})$ is bounded by

$$|\mathcal{S}(\tau, \underline{P})| \leq C^{\sum_{v \in V_{ep}(\tau)} |P_v|} \prod_{v \in V'_0(\tau)} |S_v|!, \tag{4.4.27}$$

see, e.g., [15, Lemma A.5]. Therefore, putting these observations together, we see that (4.4.26) can be bounded from above by

$$\begin{aligned}
 & 2^{h \cdot d(n,p)} |\lambda|^N \left(\max_{h < h' \leq 1} 2^{-\vartheta h'} \epsilon_{h'} \right)^M \sum_{\tau \in \mathcal{T}_{\infty; (N, M)}^{(h)}} 2^{\vartheta(h_{v_0}^{max} - h)} \sum_{P \in \mathcal{P}(\tau)} (C')^{\sum_{v \in V_{ep}(\tau)} |P_v|} \times \\
 & \times \left(\prod_{v \in V'(\tau)} 2^{(h_v - h_{v'}) \min\{-1, 2 - \frac{|P_v|}{2}\}} \right) \prod_{\substack{v \in V_e(\tau) \\ v \text{ of type } \blacklozenge}} 2^{\vartheta h_v} \prod_{\substack{v \in V_e(\tau) \\ v \text{ of type } \bullet}} |\lambda|^{[\kappa |P_v| - 1]_+}. \tag{4.4.28}
 \end{aligned}$$

We can cancel the factor of $2^{-\vartheta h}$ with a product of factors leading to an endpoint, and simplify the remaining bounds, to get

$$2^{-\vartheta h} \left(\prod_{v \in V'(\tau)} 2^{(h_v - h_{v'}) \min\{-1, 2 - \frac{|P_v|}{2}\}} \right) \prod_{\substack{v \in V_e(\tau) \\ v \text{ of type } \blacklozenge}} 2^{\vartheta h_v} \leq \left(\prod_{v \in V'(\tau)} 2^{(h_v - h_{v'}) (\vartheta - 1) \frac{|P_v|}{6}} \right).$$

This leaves a variety of exponential factors which make it possible to control the sum over \underline{P} , giving

$$\sum_{\tau \in \mathcal{T}_{\infty; (N, M)}^{(h)}} \sum_{\underline{P} \in \mathcal{P}(\tau)} (C')^{\sum_{v \in V_{eP}(\tau)} |P_v|} \left(\prod_{v \in V'(\tau)} 2^{(h_v - h_{v'}) (\vartheta - 1) \frac{|P_v|}{6}} \right) \prod_{\substack{v \in V_e(\tau) \\ v \text{ of type } \bullet}} |\lambda|^{[\kappa |P_v| - 1]_+} \leq (C'')^{N+M},$$

see, e.g., [15, App.A.6.1], from which (4.4.23) follows. □

We conclude this subsection by noting that, in order for the right side of the bound (4.4.23) to be summable over N, M uniformly in h , we need that ϵ_h is bounded and small, uniformly in h . In view of Lemma 4.8, this condition is sufficient for the whole sequence of kernels $V_{\infty}^{(h)}$, $h \leq 1$, to be well defined. In the next subsection, we study the iterative definition of the running coupling constants and prove that they in fact remain bounded and small, uniformly in h , provided that the counterterms ν_1 are properly fixed.

4.5. Beta Function Equation and Choice of the Counterterms

The definition of the running coupling constants, (4.2.30), combined with the GN tree expansion for the effective potentials implies that the running coupling constants $\underline{\nu} = \{(\nu_h, \zeta_h, \eta_h)\}_{h \leq 0}$ satisfy the following equation, for all $h \leq 0$:

$$2^h \nu_h F_{\nu, \infty} + \zeta_h F_{\zeta, \infty} + \eta_h F_{\eta, \infty} = \sum_{\tau \in \mathcal{T}_{\infty}^{(h)}} \mathcal{L}_{\infty} W_{\infty}[\underline{\nu}; \tau]. \tag{4.5.1}$$

More explicitly, using the definitions of $F_{\nu, \infty}$, $F_{\zeta, \infty}$, $F_{\eta, \infty}$ (see the lines after (4.2.4) and after (4.2.11)), for any $z \in \Lambda_{\infty}$,

$$\begin{aligned} \nu_h &= 2^{-h} (2\omega) \sum_{\tau \in \mathcal{T}_{\infty}^{(h)}} \mathcal{L}_{\infty} W_{\infty}[\underline{\nu}; \tau]((\omega, 0, z), (-\omega, 0, z)), \\ \zeta_h &= 4\omega \sum_{\tau \in \mathcal{T}_{\infty}^{(h)}} \mathcal{L}_{\infty} W_{\infty}[\underline{\nu}; \tau]((\omega, 0, z), (\omega, \hat{e}_1, z)), \\ \eta_h &= 4 \sum_{\tau \in \mathcal{T}_{\infty}^{(h)}} \mathcal{L}_{\infty} W_{\infty}[\underline{\nu}; \tau]((\omega, 0, z), (-\omega, \hat{e}_2, z)). \end{aligned} \tag{4.5.2}$$

In view of the iterative definition of the kernels $W_\infty[\underline{v}; \tau]$, see Eqs. (4.3.4) and (4.3.5), the right sides of these three equations can be naturally thought of as functions of \underline{v} or, better yet, of the restriction of \underline{v} to the scales larger than h . Therefore, these are recursive equations for the components of \underline{v} : given $V_\infty^{(1)}$ (and, in particular, given t_1^*, β, Z), one can in principle construct the whole sequence \underline{v} . More precisely, since the definition of the right sides of Eq. (4.5.2) requires the summations over τ to be well defined, in view of Lemma 4.8, these recursive equations allow one to construct the running coupling constants only for the scales h such that $\max_{h' > h} \epsilon_{h'}$ is small enough. As we shall soon see, given a sufficiently small λ , the quantity $\max_{h' > h} \epsilon_{h'}$ stays small, uniformly in h , only for a special choice of the free parameters t_1^*, β, Z ; to understand the appropriate choice of these parameters, it is helpful first to isolate those trees with only counterterm vertices, whose contribution in (4.5.2) is especially simple:

1. For $h < 0$, $\mathcal{T}_{\infty;(0,1)}^{(h)}$ consists of two trees, one with a single counterterm endpoint v on scale $h_v = 2 > h_{v_0} + 1$, for which $\mathcal{L}_\infty W_\infty[\underline{v}, \tau] = \mathcal{L}_\infty \mathcal{R}_\infty N_c = 0$, and one with a single counterterm endpoint on scale $h_v = h_{v_0} + 1$, for which we have

$$\mathcal{L}_\infty W_\infty[\underline{v}, \tau] = 2^h \nu_h F_{\nu, \infty} + \zeta_h F_{\zeta, \infty} + \eta_h F_{\eta, \infty},$$

2. For $h = 0$, $\mathcal{T}_{\infty;(0,1)}^{(h)}$ consists of only a single tree τ . This tree gives a contribution involving N_c which can be calculated quite explicitly from Eq. (3.23) and Eq. (2.1.13) taking advantage of the fact that the latter is written in terms of the Fourier transformed fields:

$$\begin{aligned} \mathcal{L}_\infty W_\infty[\underline{v}, \tau] &= \mathcal{L}_\infty N_c \\ &= \frac{1}{Z} \left(t_2 - \frac{1 - t_1}{1 + t_1} \right) F_{\nu, \infty} + \left(\frac{1}{Z} \frac{t_1}{(1 + t_1)^2} - \frac{t_1^*}{(1 + t_1^*)^2} \right) F_{\zeta, \infty} \\ &\quad + \left(\frac{t_2}{2Z} - \frac{t_2^*}{2} \right) F_{\eta, \infty} \end{aligned} \tag{4.5.3}$$

(recall $t_2^* = (1 - t_1^*) / (1 + t_1^*)$) which we can put in the same form as the other scales by letting

$$\begin{aligned} 2\nu_1 &:= \frac{1}{Z} \left(t_2 - \frac{1 - t_1}{1 + t_1} \right), \quad \zeta_1 := \frac{1}{Z} \frac{t_1}{(1 + t_1)^2} - \frac{t_1^*}{(1 + t_1^*)^2}, \\ \eta_1 &:= \frac{t_2}{2Z} - \frac{t_2^*}{2}. \end{aligned} \tag{4.5.4}$$

3. The contributions from the trees in $\mathcal{T}_{\infty,(0,M)}^{(h)}$ (i.e., those with no interaction endpoints and M counterterm endpoints) vanish for all $h \leq 0$ and $M \geq 2$, as can be seen as follows. In this case, each vertex must be assigned exactly 2 field labels, and for any endpoint v the components of ω_v must be both imaginary (corresponding to N_m , i.e., to $h_v = 2, h_{v'} = 1$) or both real (in all the other cases). If any endpoints of the first type appear there is no allowed assignment of field labels to the corresponding tree (since the ξ fields this case represents can neither be

external fields nor be contracted with ϕ fields to form a connected diagram). In any case, all of the internal field labels are included in spanning trees, and the local part of the contribution can be made quite explicit. For example, denoting by $[\mathfrak{g}_\infty^{(h)}(z - z')]_{\omega, \omega'}$ the elements of the translation invariant infinite volume propagator $\mathfrak{g}_\infty^{(h)}(z - z')$, see (2.2.4), (2.2.5), (2.2.9), the local part of the term with $\mathbf{D}_v = \mathbf{0}$ for all vertices (which contributes to ν_h) is a sum of terms proportional to

$$\begin{aligned} & \sum_{z_2, \dots, z_M \in \mathbb{Z}^2} [\mathfrak{g}_\infty^{(h_1)}(z_1 - z_2)]_{-\omega_1, \omega_2} [\mathfrak{g}_\infty^{(h_2)}(z_2 - z_3)]_{-\omega_2, \omega_3} \\ & \dots [\mathfrak{g}_\infty^{(h_{M-1})}(z_{M-1} - z_M)]_{-\omega_{M-1}, \omega_1} \\ & = [\hat{\mathfrak{g}}_\infty^{(h_1)}]_{-\omega_1, \omega_2}(0) [\hat{\mathfrak{g}}_\infty^{(h_2)}]_{-\omega_2, \omega_3}(0) \dots [\hat{\mathfrak{g}}_\infty^{(h_{M-1})}]_{-\omega_{M-1}, \omega_1}(0), \end{aligned} \tag{4.5.5}$$

for some tuple $(\omega_1, \dots, \omega_{M-1}) \in \{+, -\}^{M-1}$; from Eqs. (2.2.9), (2.2.4) and (2.2.5), the Fourier transform $\hat{\mathfrak{g}}^{(h)}$ in the right side satisfies

$$\hat{\mathfrak{g}}^{(h)}(k) = \int_{2^{-2h-2}}^{2^{-2h}} \hat{\mathfrak{g}}^{[\eta]}(k) d\eta, \quad \hat{\mathfrak{g}}^{[\eta]}(k) = e^{-\eta D(k)} D(k) \hat{\mathfrak{g}}(k) \tag{4.5.6}$$

(except for $h = 0$, where the lower limit of integration is 0), with D defined in Eq. (2.1.21), from which it is evident that $D(0) = 0$, and in consequence that the right-hand side of Eq. (4.5.5) is likewise zero. The same argument holds, *mutatis mutandis*, when \mathbf{D}_v does not vanish for all vertices of the tree, with some discrete derivatives appearing in Eq. (4.5.5).

In view of these properties, we can write

$$\nu_h = 2\nu_{h+1} + B_{h+1}^\nu[\underline{\nu}], \quad \zeta_h = \zeta_{h+1} + B_{h+1}^\zeta[\underline{\nu}], \quad \eta_h = \eta_{h+1} + B_{h+1}^\eta[\underline{\nu}], \tag{4.5.7}$$

for all $h \leq 0$, where the functions $B_{h+1}^\sharp[\underline{\nu}]$ are given by the restrictions of the sums on the right-hand sides of Eq. (4.5.2) to trees in some $\mathcal{T}_{\infty, (M, N)}^{(h)}$ with $N \geq 1$. The functions $B_{h+1}^\sharp[\underline{\nu}]$, with $\sharp \in \{\nu, \zeta, \eta\}$, are called the components of the *beta function*, and (4.5.7) are called the *beta function flow equations* for the running coupling constants; note that, even if not explicitly indicated, the functions $B_{h+1}^\sharp[\underline{\nu}]$, in addition to $\underline{\nu}$, depend analytically upon λ, t_1^*, β, Z . For later reference, we let $B_h^\sharp[\underline{\nu}; \tau]$ be the contribution to $B_{h+1}^\sharp[\underline{\nu}]$ associated with the GN tree τ . In view of Lemma 4.8, we find that, for any $|Z - 1| \leq 1/2$, any $|t_1^*|, |t_1|, |t_2| \in K'$, any $\vartheta \in (0, 1)$ there exists $C_\vartheta > 0$ such that, if $|\lambda|$ and $\max_{h' > h} \{\epsilon_{h'}\}$ are small enough, then

$$\max_{\sharp \in \{\nu, \zeta, \eta\}} |B_h^\sharp[\underline{\nu}]| \leq C_\vartheta |\lambda| 2^{\vartheta h}, \quad \forall h \leq 0. \tag{4.5.8}$$

Proposition 4.10. *For any $\vartheta \in (0, 1)$, there exist $K_\vartheta, \lambda_0(\vartheta) > 0$ and functions $\nu_1(\lambda; t_1^*, \beta, Z), \zeta_1(\lambda; t_1^*, \beta, Z), \eta_1(\lambda; t_1^*, \beta, Z), \underline{\nu}(\lambda; t_1^*, \beta, Z) := \{\nu_h(\lambda; t_1^*, \beta, Z), \zeta_h(\lambda; t_1^*, \beta, Z), \eta_h(\lambda; t_1^*, \beta, Z)\}_{h \leq 0}$, analytic in $|\lambda| \leq \lambda_0(\vartheta), |Z - 1| \leq 1/2$,*

$|t_1^*| \in K'$ and β such that $|t_1|, |t_2| \in K'$ (recall that $t_i = \tanh(\beta J_i)$), for all λ, t_1^*, β, Z in such an analyticity domain:

1. $\nu_1(\lambda; t_1^*, \beta, Z), \zeta_1(\lambda; t_1^*, \beta, Z), \eta_1(\lambda; t_1^*, \beta, Z)$ and the components of $\underline{\nu}(\lambda; t_1^*, \beta, Z)$ satisfy (4.5.7), for all $h \leq 0$;
2. For all $h \leq 1$,

$$\epsilon_h(\lambda; t_1^*, \beta, Z) := \max\{|\nu_h(\lambda; t_1^*, \beta, Z)|, |\zeta_h(\lambda; t_1^*, \beta, Z)|, |\eta_h(\lambda; t_1^*, \beta, Z)|\} \leq K_\vartheta |\lambda| 2^{\vartheta h}. \tag{4.5.9}$$

Note that, at this point, we do not prove that $\nu_1(\lambda; t_1^*, \beta, Z), \zeta_1(\lambda; t_1^*, \beta, Z), \eta_1(\lambda; t_1^*, \beta, Z)$ satisfy Eq. (4.5.4); this comes later.

Proof. For simplicity, we do not track the dependence of the constants and of the norms upon ϑ which, within the course of this proof, we assume to be a fixed constant in $(0, 1)$. In order to construct $\tilde{\nu}(\lambda; t_1^*, \beta, Z) := \{(\nu_h(\lambda; t_1^*, \beta, Z), \zeta_h(\lambda; t_1^*, \beta, Z), \eta_h(\lambda; t_1^*, \beta, Z))\}_{h \leq 1}$, we first note that the equations for ν_h, ζ_h, η_h in (4.5.7) imply that, for $k < h \leq 1$, $\nu_h = 2^{k-h} \nu_k - \sum_{k < j \leq h} 2^{j-h-1} B_j^\nu[\underline{\nu}]$, $\zeta_h = \zeta_k - \sum_{k < j \leq h} B_j^\zeta[\underline{\nu}]$, and $\eta_h = \eta_k - \sum_{k < j \leq h} B_j^\eta[\underline{\nu}]$. If we send $k \rightarrow -\infty$ and impose that $\epsilon_k := \max\{|\nu_k|, |\zeta_k|, |\eta_k|\} \rightarrow 0$ as $k \rightarrow -\infty$, we get

$$\begin{cases} \nu_h = - \sum_{j \leq h} 2^{j-h-1} B_j^\nu[\underline{\nu}], \\ \zeta_h = - \sum_{j \leq h} B_j^\zeta[\underline{\nu}], \\ \eta_h = - \sum_{j \leq h} B_j^\eta[\underline{\nu}], \end{cases} \tag{4.5.10}$$

which we regard as a fixed point equation $\tilde{\nu} = \mathbf{T}[\tilde{\nu}]$ for a map \mathbf{T} on the space of sequences $X_\epsilon := \{\tilde{\nu} = \{(\nu_h, \zeta_h, \eta_h)\}_{h \leq 1} : \|\tilde{\nu}\| \leq \epsilon\}$, with $\|\tilde{\nu}\| = \sup_{h \leq 1} \{2^{-\vartheta h} \epsilon_h\}$ and ϵ a sufficiently small constant.

We now intend to prove that \mathbf{T} is a contraction on X_ϵ and, more precisely, that: (1) the image of X_ϵ under the action of \mathbf{T} is contained in X_ϵ ; (2) $\|\mathbf{T}[\tilde{\nu}] - \mathbf{T}[\tilde{\nu}']\| \leq (1/2) \|\tilde{\nu} - \tilde{\nu}'\|$ for all $\tilde{\nu}, \tilde{\nu}' \in X_\epsilon$. Once \mathbf{T} is proved to be a contraction, it follows that it admits a unique fixed point in X_ϵ , which corresponds to the desired sequence $\tilde{\nu}(\lambda; t_1^*, \beta, Z)$. The analyticity of $\tilde{\nu}(\lambda; t_1^*, \beta, Z)$ follows from the analyticity of the components of the beta function with respect to λ, t_1^*, β, Z and $\underline{\nu}$ that, in turn, follows from the absolute summability of its tree expansion, which is a power series in λ , with coefficients analytically depending upon t_1^*, t_1, t_2, Z , with $t_i = \tanh(\beta J_i)$.

The fact that the image of X_ϵ under the action of \mathbf{T} is contained in X_ϵ follows immediately from Eq. (4.5.8). In order to prove that $\|\mathbf{T}[\tilde{\nu}] - \mathbf{T}[\tilde{\nu}']\| \leq (1/2) \|\tilde{\nu} - \tilde{\nu}'\|$, we rewrite the ν -component of $\mathbf{T}[\tilde{\nu}] - \mathbf{T}[\tilde{\nu}']$ at scale h as a linear interpolation

$$\nu_h - \nu_{h'} = - \sum_{j \leq h} \sum_{\tau \in \mathcal{T}_\infty^{(j-1)}} 2^{j-h-1} \int_0^1 \frac{d}{dt} B_j^\nu[\tilde{\nu}(t); \tau] dt, \tag{4.5.11}$$

where $\tilde{\nu}(t) = \tilde{\nu}' + t(\tilde{\nu} - \tilde{\nu}')$, and similarly for the ζ - and η -components. When the derivative with respect to t acts on the tree value $B_j^\nu[\underline{\nu}(t); \tau]$, it has the effect of replacing one of the factors $\nu_h(t)$, or $\zeta_h(t)$, or $\eta_h(t)$, associated with

one of the counterterm endpoints, by $\frac{d}{dt}\nu_h(t) = \nu_h - \nu'_h$, or $\frac{d}{dt}\zeta_h(t) = \zeta_h - \zeta'_h$, or $\frac{d}{dt}\eta_h(t) = \eta_h - \eta'_h$, respectively. Therefore, we get the analogue of Eq. (4.5.8): if $|\lambda|$ and $\|\tilde{\nu}(t)\|$ are sufficiently small, then

$$\begin{aligned} & \max_{\#\in\{\nu,\zeta,\eta\}} \sum_{\tau\in\mathcal{T}_\infty^{(h-1)}} \left| \frac{d}{dt} B_h^\#[\underline{\nu}(t); \tau] \right| \\ & \leq C|\lambda|2^{\vartheta h} \max_{h'\geq h} \{|\nu_{h'} - \nu'_{h'}|, |\zeta_{h'} - \zeta'_{h'}|, |\eta_{h'} - \eta'_{h'}|\} \end{aligned} \tag{4.5.12}$$

for all $h \leq 1$. Plugging this estimate into Eq. (4.5.11) and its analogues for $\zeta_h - \zeta'_h$ and $\eta_h - \eta'_h$, we readily obtain the desired estimate, $\|\mathbf{T}[\tilde{\nu}] - \mathbf{T}[\tilde{\nu}']\| \leq (1/2) \|\tilde{\nu} - \tilde{\nu}'\|$, for λ_0 and ε sufficiently small. \square

We now need to show that it is possible to choose the free parameters β, Z, t_1^* in such a way that the functions $\nu_1(\lambda; t_1^*, \beta, Z)$, $\zeta_1(\lambda; t_1^*, \beta, Z)$, $\eta_1(\lambda; t_1^*, \beta, Z)$ constructed in Proposition 4.10 satisfy Eq. (4.5.4); that is, for given J_1, J_2 and λ , there exists a critical value of the inverse temperature β for which the above expansion for the kernels of the infinite plane effective potentials is convergent, with dressed parameters t_1^* and Z . The desired result is summarized in the following proposition.

Proposition 4.11. *For any J_1, J_2 satisfying the conditions of Theorem 1.1, and any $\vartheta \in (0, 1)$, there exist $\lambda_0(\vartheta) > 0$ and functions $t_1^*(\lambda)$, $\beta_c(\lambda)$, $Z(\lambda)$, analytic in $|\lambda| \leq \lambda_0(\vartheta)$, such that Eq. (4.5.4) holds, with $t_1 = \tanh(\beta_c(\lambda)J_1)$, $t_2 = \tanh(\beta_c(\lambda)J_2)$, and $(\nu_1, \zeta_1, \eta_1) = ((\tilde{\nu}_1(\lambda), \tilde{\zeta}_1(\lambda), \tilde{\eta}_1(\lambda)))$, with $\tilde{\nu}_1(\lambda) = \nu_1(\lambda; t_1^*(\lambda), \beta_c(\lambda), Z(\lambda))$ (here $\nu_1(\lambda; t_1^*, \beta, Z)$ is the same as in Proposition 4.10), and similarly for $\tilde{\zeta}_1(\lambda)$ and $\tilde{\eta}_1(\lambda)$. Correspondingly, the flow of running coupling constants with initial datum $(\tilde{\nu}_1(\lambda), \tilde{\zeta}_1(\lambda), \tilde{\eta}_1(\lambda))$, generated by the flow equations (4.5.7), is well defined for all $h \leq 0$ and satisfies (4.5.9).*

Proof. The result is a direct consequence of the analytic implicit function theorem: we intend to fix $t_1^* = t_1^*(\lambda)$, $\beta = \beta_c(\lambda)$, $Z = Z(\lambda)$ in such a way that Eq. (4.5.4) holds, with $\nu_1 = \nu_1(\lambda; t_1^*, \beta, Z)$, $\zeta_1 = \zeta_1(\lambda; t_1^*, \beta, Z)$, $\eta_1 = \eta_1(\lambda; t_1^*, \beta, Z)$, the same functions as in Proposition 4.10. With this in mind, we recast the system of equations (4.5.4) in the following form (recall once more that $t_i = \tanh(\beta J_i)$ and $t_2^* = (1 - t_1^*)/(1 + t_1^*)$):

$$2Z\nu_1(\lambda; t_1^*, \beta, Z) - \tanh \beta J_2 + e^{-2\beta J_1} = 0, \tag{4.5.13}$$

$$4Z\zeta_1(\lambda; t_1^*, \beta, Z) + \frac{4Zt_1^*}{(1 + t_1^*)^2} + e^{-4\beta J_1} - 1 = 0, \tag{4.5.14}$$

$$2Z\eta_1(\lambda; t_1^*, \beta, Z) + Z\frac{1 - t_1^*}{1 + t_1^*} - \tanh \beta J_2 = 0. \tag{4.5.15}$$

Note that, by Proposition 4.10, everything appearing in Eqs. (4.5.13) to (4.5.15) is analytic, so if we succeed in showing that the implicit function theorem applies, the resulting solution will be analytic in λ , as desired. Note that, at $\lambda = 0$, we have $\nu_1(0; t_1^*, \beta, Z) = \zeta(0; t_1^*, \beta, Z) = \eta_1(0; t_1^*, \beta, Z) = 0$, so in this case the system (4.5.13)–(4.5.15) is solved by $\beta = \beta_c(J_1, J_2)$ (with

$\beta_c(J_1, J_2)$ the critical temperature of the nearest neighbor model, see the lines after (1.3), $t_1^* = \tanh(\beta_c(J_1, J_2)J_1)$ and $Z = 1$. Note also that the partial derivatives of $\nu_1(\lambda; t_1^*, \beta, Z), \zeta_1(\lambda; t_1^*, \beta, Z), \eta_1(\lambda; t_1^*, \beta, Z)$ with respect to t_1^*, β and Z vanish at $\lambda = 0$, so the determinant of the Jacobian of the system (4.5.13)–(4.5.15) with respect to t_1^*, β, Z , computed at $\lambda = 0$ and $(t_1^*, \beta, Z) = (\tanh(\beta_c(J_1, J_2)J_1), \beta_c(J_1, J_2), 1)$, equals

$$\begin{aligned} & \begin{vmatrix} 0 & -J_2 \operatorname{sech}^2(\beta_c J_2) - 2J_1 e^{-2\beta_c J_1} & 0 \\ 4 \frac{1-t_1^*}{(1+t_1^*)^3} & -4J_1 e^{-4\beta_c J_2} & \frac{4t_1^*}{(1+t_1^*)^2} \\ -\frac{2}{(1+t_1^*)^2} & -J_2 \operatorname{sech}^2(\beta_c J_2) & \frac{1-t_1^*}{1+t_1^*} \end{vmatrix} \\ & = -4(J_2 \operatorname{sech}^2(\beta_c J_2) + 2J_1 e^{-2\beta_c J_1}) \frac{1 + (t_1^*)^2}{(1 + t_1^*)^4} \end{aligned} \tag{4.5.16}$$

with $\beta_c = \beta_c(J_1, J_2)$ and $t_1^* = \tanh(\beta_c(J_1, J_2)J_1)$; the right-hand side is evidently nonzero, and, therefore, the analytic implicit function theorem applies, implying the desired claim. \square

This concludes the construction of the sequence of effective potentials in the infinite volume limit, uniformly in the scale label, with optimal bounds on the speed at which, after proper rescaling, such effective potentials go to zero as $h \rightarrow -\infty$. This result, and the methods introduced to prove it, is a key ingredient in the proof of Theorem 1.1, for whose completion we refer the reader to [4].

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A. Diagonalization of the Matrix A_c

In this section we compute the propagator of ϕ and the Gaussian integral associated with \mathcal{S}_c . For this purpose, we first need to block diagonalize the coefficient matrix, which we do by using a transformation which can be thought of as a Fourier sine transformation with modified frequencies. We write

$$\begin{aligned} \mathcal{S}_c(\phi) &= \frac{1}{2L} \sum_{k_1 \in \mathcal{D}_L} \left(\begin{bmatrix} \phi_+(-k_1) \\ \phi_-(-k_1) \end{bmatrix}, \begin{bmatrix} -i\Delta(k_1) & -b(k_1) + t_2\tau \\ b(k_1) - t_2\tau^T & i\Delta(k_1) \end{bmatrix} \begin{bmatrix} \phi_+(k_1) \\ \phi_-(k_1) \end{bmatrix} \right) \\ &=: \frac{1}{2L} \sum_{k_1 \in \mathcal{D}_L} \left(\begin{bmatrix} \phi_+(-k_1) \\ \phi_-(-k_1) \end{bmatrix}, \tilde{A}_c(k_1) \begin{bmatrix} \phi_+(k_1) \\ \phi_-(k_1) \end{bmatrix} \right), \end{aligned} \tag{A.1}$$

where $\phi_\omega(k_1)$, with $\omega = \pm$, is the column vector whose components are $\phi_{\omega, z_2}(k_1)$ with $z_2 = 1, \dots, M$, and τ is the $M \times M$ shift matrix $\tau_{z_2, z'_2} := \delta_{z_2+1, z'_2}$, that is,

$$\tau = \begin{bmatrix} 0 & 1 & 0 & & & & & & \\ 0 & 0 & 1 & \ddots & & & & & \\ 0 & \ddots & \ddots & \ddots & \ddots & & & & \\ & \ddots & \ddots & \ddots & 1 & 0 & & & \\ & & \ddots & 0 & 0 & 1 & & & \\ & & & 0 & 0 & 0 & & & \end{bmatrix}.$$

For brevity, we will write $\tilde{A}_c = \tilde{A}_c(k_1)$, $\Delta = \Delta(k_1)$, $b = b(k_1)$ since dependence on k_1 plays no role in the next several pages. It is helpful to begin by diagonalizing the real symmetric matrix

$$\tilde{A}_c^2 = \begin{bmatrix} \tilde{B}_c^+ & 0 \\ 0 & \tilde{B}_c^- \end{bmatrix},$$

where \tilde{B}_c^+ is the $M \times M$ tri-diagonal matrix

$$\tilde{B}_c^+ = \begin{bmatrix} -\Delta^2 - b^2 - t_2^2 & bt_2 & & & & & & & \\ & bt_2 & -\Delta^2 - b^2 - t_2^2 & bt_2 & & & & & \\ & & bt_2 & -\Delta^2 - b^2 - t_2^2 & \ddots & & & & \\ & & & \ddots & \ddots & \ddots & & & \\ & & & & \ddots & \ddots & \ddots & & \\ & & & & & bt_2 & -\Delta^2 - b^2 - t_2^2 & bt_2 & \\ & & & & & & bt_2 & -\Delta^2 - b^2 \end{bmatrix}.$$

Note that all the diagonal entries are equal to $-\Delta^2 - b^2 - t_2^2$ apart from the last one, which equals $-\Delta^2 - b^2$. The block \tilde{B}_c^- is obtained from \tilde{B}_c^+ by reversing

the order of the rows and columns. \tilde{B}_c^ω , with $\omega = \pm$, can each be thought of as a discrete Laplacian with mixed boundary conditions, which suggests the ansatz

$$v_{k_2} = \begin{bmatrix} \alpha_{k_2} e^{ik_2} + \beta_{k_2} e^{-ik_2} \\ \alpha_{k_2} e^{i2k_2} + \beta_{k_2} e^{-i2k_2} \\ \vdots \end{bmatrix}, \tag{A.2}$$

for their eigenvectors. In fact, we see that v_{k_2} is an eigenvector of \tilde{B}_c^+ iff the system of equations

$$\begin{aligned} (-\Delta^2 - b^2 - t_2^2)(\alpha_{k_2} e^{ik_2} + \beta_{k_2} e^{-ik_2}) + bt_2(\alpha_{k_2} e^{i2k_2} + \beta_{k_2} e^{-i2k_2}) \\ = \lambda_{k_2}(\alpha_{k_2} e^{ik_2} + \beta_{k_2} e^{-ik_2}) \end{aligned} \tag{A.3}$$

$$\begin{aligned} bt_2(\alpha_{k_2} e^{i(z_2-1)k_2} + \beta_{k_2} e^{-ik_2(z_2-1)}) + (-\Delta^2 - b^2 - t_2^2)(\alpha_{k_2} e^{ik_2 z_2} + \beta_{k_2} e^{-ik_2 z_2}) \\ + bt_2(\alpha_{k_2} e^{ik_2(z_2+1)} + \beta_{k_2} e^{-ik_2(z_2+1)}) \\ = \lambda_{k_2}(\alpha_{k_2} e^{ik_2 z_2} + \beta_{k_2} e^{-ik_2 z_2}), \quad 1 < z_2 < M \end{aligned} \tag{A.4}$$

$$\begin{aligned} bt_2(\alpha_{k_2} e^{ik_2(M-1)} + \beta_{k_2} e^{-ik_2(M-1)}) + (-\Delta^2 - b^2)(\alpha_{k_2} e^{ik_2 M} + \beta_{k_2} e^{-ik_2 M}) \\ = \lambda_{k_2}(\alpha_{k_2} e^{ik_2 M} + \beta_{k_2} e^{-ik_2 M}) \end{aligned} \tag{A.5}$$

are all satisfied. Equation (A.4) is solved by choosing

$$\lambda_{k_2} = bt_2(e^{ik_2} + e^{-ik_2}) + (-\Delta^2 - b^2 - t_2^2),$$

which reduces the other two conditions to

$$bt_2(\alpha_{k_2} + \beta_{k_2}) = 0, \tag{A.6}$$

$$bt_2(\alpha_{k_2} e^{ik_2(M+1)} + \beta_{k_2} e^{-ik_2(M+1)}) - t_2^2(\alpha_{k_2} e^{ik_2 M} + \beta_{k_2} e^{-ik_2 M}) = 0. \tag{A.7}$$

The first condition implies $\beta_{k_2} = -\alpha_{k_2}$, which can be used to rewrite Eq. (A.7) as

$$\sin k_2(M + 1) = B(k_1) \sin k_2 M \tag{A.8}$$

where for brevity we have introduced

$$B(k_1) := \frac{t_2}{b(k_1)} = t_2 \frac{|1 + t_1 e^{ik_1}|^2}{1 - t_1^2}, \tag{A.9}$$

(cf. Eq. (2.1.22)). We call $\mathcal{Q}_M^+(k_1)$ the set of the solutions of (2.1.23) with $\Re k_2 \in [0, \pi]$.

Restricting to the critical case (2.1.17),

$$B(k_1) = 1 - \frac{2t_1 t_2}{1 - t_1^2} (1 - \cos k_1) =: 1 - \kappa (1 - \cos k_1), \tag{A.10}$$

so that $0 < B(k_1) < 1$ for $k_1 \in \mathcal{D}_L$. Equation (2.1.23) is equivalent to

$$\tan k_2(M + 1) = \frac{B(k_1) \sin k_2}{B(k_1) \cos k_2 - 1}, \tag{A.11}$$

which for $0 < B(k_1) < 1$ has a unique real solution in each interval $I_n := \frac{\pi}{M+1}(n + \frac{1}{2}, n + 1)$, $n = 0, \dots, M - 1$, since the left-hand side increases monotonically from $-\infty$ to 0, while the right-hand side is negative and decreasing. Thus, all M eigenvectors of \tilde{B}_c^+ (and, as a consequence, of \tilde{B}_c^-) correspond to real solutions of this form by

$$u_{k_2}^+ = \sqrt{\frac{2}{N_M(k_1, k_2)}} \begin{bmatrix} \sin k_2 \\ \sin 2k_2 \\ \vdots \\ \sin k_2 M \end{bmatrix} \left(\text{and, respectively, } u_{k_2}^- = \sqrt{\frac{2}{N_M(k_1, k_2)}} \begin{bmatrix} \sin k_2 M \\ \sin k_2 (M - 1) \\ \vdots \\ \sin k_2 \end{bmatrix} \right), \tag{A.12}$$

where

$$N_M(k_1, k_2) := 2 \sum_{x=1}^M \sin^2 k_2 x = M + \frac{1}{2} - \frac{1}{2} \frac{\sin(2M + 1)k_2}{\sin k_2} \tag{A.13}$$

so that the eigenvectors are normalized.

To obtain Eq. (2.1.24), we note that since q satisfies (A.8), we have

$$\begin{aligned} \sin [k_2(M + 1) \pm k_2 M] &= \sin k_2(M + 1) \cos k_2 M \pm \sin k_2 M \cos k_2(M + 1) \\ &= [B(k_1) \cos k_2 M \pm \cos k_2(M + 1)] \sin k_2 M, \end{aligned}$$

and so we can rewrite $N_M(k_1, k_2)$ as

$$\begin{aligned} N_M(k_1, k_2) &= \frac{B(k_1)M \cos k_2 M - (M + 1) \cos k_2(M + 1)}{B(k_1) \cos k_2 M - \cos k_2(M + 1)} \\ &= \frac{\frac{d}{dk_2} (B(k_1) \sin k_2 M - \sin k_2(M + 1))}{B(k_1) \cos k_2 M - \cos k_2(M + 1)}. \end{aligned} \tag{A.14}$$

We now return to \tilde{A}_c . Equation (2.1.23) is equivalent to

$$be^{ik_2(M+1)} - t_2 e^{ik_2 M} = be^{-ik_2(M+1)} - t_2 e^{-ik_2 M}, \tag{A.15}$$

and therefore

$$\begin{aligned} (b - t_2 \tau^T) u_{k_2}^+ &= -(be^{ik_2(M+1)} - t_2 e^{ik_2 M}) u_{k_2}^-, \\ \text{and } (-b + t_2 \tau) u_{k_2}^- &= (be^{ik_2(M+1)} - t_2 e^{ik_2 M}) u_{k_2}^+ \end{aligned}$$

whenever $k_2 \in \mathcal{Q}_M(k_1)$. Combining this with the definition of $\tilde{A}_c(k_1)$ in Eq. (A.1), we see that

$$\begin{aligned} \tilde{A}_c \begin{bmatrix} u_{k_2}^+ \\ 0 \end{bmatrix} &= \begin{bmatrix} -i\Delta u_{k_2}^+ \\ -(be^{ik_2(M+1)} - t_2 e^{ik_2 M}) u_{k_2}^- \end{bmatrix}, \\ \text{and } \tilde{A}_c \begin{bmatrix} 0 \\ u_{k_2}^- \end{bmatrix} &= \begin{bmatrix} (be^{ik_2(M+1)} - t_2 e^{ik_2 M}) u_{k_2}^+ \\ i\Delta u_{k_2}^- \end{bmatrix}, \end{aligned} \tag{A.16}$$

or in other words, the change of variables induced by $\begin{bmatrix} u_{k_2}^+ \\ 0 \end{bmatrix}$ and $\begin{bmatrix} 0 \\ u_{\bar{k}_2} \end{bmatrix}$ puts \tilde{A}_c in block-diagonal form, with 2×2 blocks

$$\begin{aligned} \tilde{\mathfrak{g}}^{-1}(k_1, k_2) &:= \begin{bmatrix} -i\Delta & e^{-ik_2(M+1)}(b - t_2 e^{ik_2}) \\ -e^{ik_2(M+1)}(b - t_2 e^{-ik_2}) & i\Delta \end{bmatrix} \\ &= \frac{1}{|1 + t_1 e^{ik_1}|^2} \begin{bmatrix} -2it_1 \sin k_1 & e^{-ik_2(M+1)}(1 - t_1^2)(1 - B(k_1)e^{ik_2}) \\ -e^{ik_2(M+1)}(1 - t_1^2)(1 - B(k_1)e^{-ik_2}) & 2it_1 \sin k_1 \end{bmatrix}, \end{aligned} \tag{A.17}$$

recalling the definitions (2.1.14) and (2.1.22). This block-diagonalization implies that

$$\text{Pf} A_c = \prod_{k_1 \in \mathcal{D}_L} \prod_{k_2 \in \mathcal{Q}_M^+(k_1)} \sqrt{\det \tilde{\mathfrak{g}}^{-1}(k_1, k_2)}, \tag{A.18}$$

where explicitly, using the criticality condition (2.1.17),

$$\det \tilde{\mathfrak{g}}^{-1}(k_1, k_2) = \frac{2(1 - t_2)^2(1 - \cos k_1) + 2(1 - t_1)^2(1 - \cos k_2)}{|1 + t_1 e^{ik_1}|^2}. \tag{A.19}$$

Note that this determinant vanishes iff $k_1 = k_2 = 0 \pmod{2\pi}$ (in particular, it is positive if $k_1 \in \mathcal{D}_L$).

Concerning the propagator, denoting the inverse of (A.17) by

$$\begin{aligned} \tilde{\mathfrak{g}}(k_1, k_2) &:= \begin{bmatrix} \tilde{g}_{++}(k_1, k_2) & \tilde{g}_{+-}(k_1, k_2) \\ \tilde{g}_{-+}(k_1, k_2) & \tilde{g}_{--}(k_1, k_2) \end{bmatrix} \\ &:= \frac{1}{D(k_1, k_2)} \\ &\begin{bmatrix} 2it_1 \sin k_1 & -e^{-ik_2(M+1)}(1 - t_1^2)(1 - B(k_1)e^{ik_2}) \\ e^{ik_2(M+1)}(1 - t_1^2)(1 - B(k_1)e^{-ik_2}) & -2it_1 \sin k_1 \end{bmatrix}, \end{aligned} \tag{A.20}$$

where $D(k_1, k_2)$ is defined as in (2.1.21) and, letting

$$\tilde{\phi}_{k_2, \omega}(k_1) := \sum_{z_2=1}^M \phi_{\omega, x_2}(k_1) u_{k_2}^\omega(z_2),$$

we have, for $k_1, k'_1 \in \mathcal{D}_L$, $k_2, k'_2 \in \mathcal{Q}_M^+$, and $\omega, \omega' \in \{\pm\}$,

$$\langle \tilde{\phi}_{k_2, \omega}(k_1) \tilde{\phi}_{k'_2, \omega'}(k'_1) \rangle = -L \delta_{k_1, -k'_1} \delta_{k_2, k'_2} \tilde{g}_{\omega\omega'}(k_1, k_2), \tag{A.21}$$

so that, in terms of $\phi_{\omega, z} = \frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} e^{-ik_1 z_1} \phi_{\omega, z_2}(k_1)$,

$$\begin{aligned} \langle \phi_{\omega, z} \phi_{\omega', z'} \rangle &= -\frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} \sum_{k_2 \in \mathcal{Q}_M^+(k_1)} \tilde{g}_{\omega\omega'}(k_1, k_2) e^{-ik_1(z_1 - z'_1)} u_{k_2}^\omega(z_2) u_{k_2}^{\omega'}(z'_2) \\ &\equiv g_{\omega\omega'}(z, z'). \end{aligned} \tag{A.22}$$

Then recalling the definition of $u_{k_2}^\pm(z_2)$ and the identity (A.15), Eqs. (2.1.18) and (2.1.19) follow by writing out the sines in terms of complex exponentials and relabeling the sum in terms of $\mathcal{Q}_M(k_1) := \mathcal{Q}_M^+(k_1) \sqcup (-\mathcal{Q}_M^+(k_1))$.

B. Proof of Proposition 2.3

For the proof of items 1 and 2 it is convenient to start by proving their analogues for the infinite volume limit propagators.

Decay bounds on $\mathfrak{g}_\infty^{[\eta]}$, $\mathfrak{g}_\infty^{(h)}$. Recall that $\mathfrak{g}_\infty^{[\eta]}$ was defined in (2.2.9). We intend to prove that, for all $x \in \mathbb{Z}^2$,

$$\|\partial_1^r \partial_2^s \mathfrak{g}_\infty^{[\eta]}(z)\| \leq C^{1+r+s} \times \begin{cases} r!s!\eta^{-\frac{3+r+s}{2}} e^{-\eta^{-1/2}|z|_1} & \text{if } \eta \geq 1, \\ e^{-|z|_1} & \text{if } 0 \leq \eta \leq 1, \end{cases} \tag{B.1}$$

where ∂_j is the discrete derivative with respect to the j -th coordinate. By using (B.1) in the definition of $\mathfrak{g}_\infty^{(h)}$, namely

$$\mathfrak{g}_\infty^{(h)}(z) = \begin{cases} \int_0^1 \mathfrak{g}_\infty^{[\eta]}(z) d\eta, & \text{if } h = 0, \\ \int_{2^{-2h-2}}^{2^{-2h}} \mathfrak{g}_\infty^{[\eta]}(z) d\eta, & \text{if } h < 0, \end{cases} \tag{B.2}$$

we obtain the analogue of (2.2.15),

$$\|\partial_1^r \partial_2^s \mathfrak{g}_\infty^{(h)}(z)\| \leq C^{1+r+s} r!s! 2^{(1+r+s)h} e^{-c2^h|z|_1}, \tag{B.3}$$

for all $z \in \mathbb{Z}^2$ and $h \leq 0$.

In order to prove (B.1), we start from the explicit expression of the function in the left side,

$$\begin{aligned} & \partial_1^r \partial_2^s \mathfrak{g}_\infty^{[\eta]}(z) \\ &= \int_{[-\pi, \pi]^2} e^{-i(k_1 z_1 + k_2 z_2)} (e^{-ik_1} - 1)^r (e^{-ik_2} - 1)^s e^{-\eta D(k_1, k_2)} M(k_1, k_2) \frac{dk_1 dk_2}{(2\pi)^2}, \end{aligned} \tag{B.4}$$

where $D(k_1, k_2)$ is as in Eq. (2.1.21) (cf. Eq. (2.2.1)) and

$$\begin{aligned} M(k_1, k_2) &:= D(k_1, k_2) \hat{\mathfrak{g}}(k_1, k_2) \\ &= \begin{bmatrix} -2it_1 \sin k_1 & -(1-t_1^2)[1 - e^{-ik_2} B(k_1)] \\ (1-t_1^2)[1 - e^{ik_2} B(k_1)] & 2it_1 \sin k_1 \end{bmatrix}. \end{aligned} \tag{B.5}$$

Note that the integrand in Eq. (B.4) is periodic with period 2π both in k_1 and in k_2 , and is entire in both arguments. Therefore, we can shift both variables in the complex plane, $k_1 \rightarrow k_1 - ia$ and $k_2 \rightarrow k_2 - ib$, to obtain

$$\begin{aligned} \partial_1^r \partial_2^s \mathfrak{g}_\infty^{[\eta]}(z) &= e^{-az_1 - bz_2} \int_{[-\pi, \pi]^2} e^{-i(k_1 x_1 + k_2 x_2)} (e^{-ik_1 - a} - 1)^r (e^{-ik_2 - b} - 1)^s \\ &\quad \times e^{-\eta D(k_1 - ia, k_2 - ib)} M(k_1 - ia, k_2 - ib) \frac{dk_1 dk_2}{(2\pi)^2}. \end{aligned} \tag{B.6}$$

We now pick $a = \alpha \operatorname{sign}(z_1)$ and $b = \alpha \operatorname{sign}(z_2)$, with $\alpha = \min\{\eta^{-1/2}, 1\}$, and take the absolute value, thus getting

$$\begin{aligned} \|\partial_1^r \partial_2^s \mathfrak{g}_\infty^{[\eta]}(z)\| &\leq e^{-\alpha|z|_1} e^{-2\eta[(1-t_1)^2+(1-t_2)^2](\cosh \alpha-1)} \\ &\quad \times \int_{[-\pi,\pi]^2} e^{-\eta D(k_1,k_2)} (\alpha e^\alpha + |e^{-ik_1} - 1|)^r \\ &\quad (\alpha e^\alpha + |e^{-ik_2} - 1|)^s M_0(k_1, k_2, \alpha) \frac{dk_1 dk_2}{(2\pi)^2}, \end{aligned} \tag{B.7}$$

where $M_0(k_1, k_2, \alpha) = \max_{|a|=|b|=\alpha} \|M(k_1 - ia, k_2 - ib)\|$ and we used the fact that, if $|a| = |b| = \alpha$, then

$$|D(k_1 - ia, k_2 - ib)| \geq D(k_1, k_2) + 2[(1 - t_1)^2 + (1 - t_2)^2](\cosh \alpha - 1) \tag{B.8}$$

and

$$|e^{-ik_1-a} - 1| \leq \alpha e^\alpha + |e^{-ik_1} - 1|, \quad |e^{-ik_2-b} - 1| \leq \alpha e^\alpha + |e^{-ik_2} - 1|. \tag{B.9}$$

Now, if $\eta \leq 1$ and, therefore, $\alpha = 1$, then (B.7) immediately implies that $\|\partial_1^r \partial_2^s \mathfrak{g}_\infty^{[\eta]}(z)\| \leq C^{1+r+s} e^{-|z|_1}$, as desired. If $\eta \geq 1$ and, therefore, $\alpha = \eta^{-1/2}$, we make the following observations: (i) the factor $e^{-2\eta[(1-t_1)^2+(1-t_2)^2](\cosh \eta^{-1/2}-1)}$ is bounded from above uniformly in η ; (ii) if $-\pi \leq k_1, k_2 \leq \pi$, then $D(k_1, k_2) \geq c(k_1^2 + k_2^2)$, $|e^{-ik_1} - 1| \leq C|k_1|$, $|e^{-ik_2} - 1| \leq C|k_2|$ and $M_0(k_1, k_2, \eta^{-1/2}) \leq C(\eta^{-1/2} + |k_1| + |k_2|)$. By using these inequalities in (B.7), we find

$$\begin{aligned} \|\partial_1^r \partial_2^s \mathfrak{g}_\infty^{[\eta]}(z)\| &\leq C^{1+r+s} e^{-\eta^{-1/2}|z|_1} \\ &\quad \times \int_{\mathbb{R}^2} e^{-c\eta(k_1^2+k_2^2)} (\eta^{-1/2} + |k_1|)^r (\eta^{-1/2} + |k_2|)^s \\ &\quad (\eta^{-1/2} + |k_1| + |k_2|) dk_1 dk_2, \end{aligned} \tag{B.10}$$

and expanding the powers in the integrand we obtain a sum of Gaussian integrals which reduce to (B.1) for this case as well.

Decay bounds on $\mathfrak{g}_E^{[\eta]}$, $\mathfrak{g}_E^{(h)}$, and proof of items 1 and 2. Recall that $\mathfrak{g}_E^{[\eta]}(z, z') = \mathfrak{g}^{[\eta]}(z, z') - \mathfrak{g}_B^{[\eta]}(z, z')$, with $\mathfrak{g}^{[\eta]}$ as in (2.2.7) and $\mathfrak{g}_B^{[\eta]}$ as in (2.2.10). We focus on the case that $z_1 - z'_1 \neq \pm L/2$ (recall that in our conventions $z_1, z'_1 \in \{1, \dots, L\}$), in which the function s_L in (2.2.10) is equal to ± 1 ; the complementary case, $z_1 - z'_1 = \pm L/2$, in which $\mathfrak{g}_B^{[\eta]}(z, z') = 0$, can be treated in a way analogous to the discussion below, and is left to the reader. If $z_1 - z'_1 \neq \pm L/2$, by using the anti-periodicity of the propagator in the horizontal direction, we can reduce without loss of generality to the case $z_1 - z'_1 = \operatorname{per}_L(z_1 - z'_1)$ (i.e., $-L/2 < z_1 - z'_1 < L/2$), in which $\mathfrak{g}_B^{[\eta]}(z, z') = \mathfrak{g}_\infty^{[\eta]}(z - z')$, and we shall do so in the following. Therefore, in this case,

$$\mathfrak{g}_E^{[\eta]}(z, z') = \sum_{k_1 \in \mathcal{D}_L} \sum_{k_2 \in \mathcal{Q}_M(k_1)} \frac{1}{2LN_M(k_1, k_2)} \left[G_+^{[\eta]}(k_1, k_2; z, z') - G_-^{[\eta]}(k_1, k_2; z, z') \right]$$

$$-\mathbf{g}_\infty^{[\eta]}(z - z'), \tag{B.11}$$

where, recalling the definitions of f_η in (2.2.1) and of $\hat{\mathbf{g}}(k_1, k_2)$ and $\hat{g}_{\omega\omega'}(k_1, k_2)$ in (2.1.20),

$$\begin{aligned} G_+^{[\eta]}(k_1, k_2; z, z') &:= e^{-ik_1(z_1-z'_1)} e^{-ik_2(z_2-z'_2)} f_\eta(k_1, k_2) \hat{\mathbf{g}}(k_1, k_2), \\ G_-^{[\eta]}(k_1, k_2; z, z') &:= e^{-ik_1(z_1-z'_1)} e^{-ik_2(z_2+z'_2)} f_\eta(k_1, k_2) \\ &\quad \begin{bmatrix} \hat{g}_{++}(k_1, k_2) & \hat{g}_{+-}(k_1, -k_2) \\ \hat{g}_{-+}(k_1, k_2) & e^{2ik_2(M+1)} \hat{g}_{--}(k_1, k_2) \end{bmatrix}, \end{aligned} \tag{B.12}$$

which are entire functions of k_1, k_2 , and 2π -periodic in both variables. We intend to prove that, for all $z, z' \in \Lambda$,

$$\|\partial^r \mathbf{g}_E^{[\eta]}(z, z')\| \leq C^{1+|r|_1} \times \begin{cases} \mathbf{r}! \eta^{-\frac{3+|r|_1}{2}} e^{-c\eta^{-1/2} d_E(z, z')} & \text{if } 1 \leq \eta \leq 2^{-2h^*}, \\ e^{-c d_E(z, z')} & \text{if } 0 \leq \eta \leq 1, \end{cases} \tag{B.13}$$

where we recall that $h^* = -\lfloor \log_2 \min\{L, M\} \rfloor$. Recalling also the relationship between $\mathbf{g}_E^{(h)}$ and $\mathbf{g}_E^{[\eta]}$, this implies

$$\begin{aligned} \|\partial^r \mathbf{g}_E^{(h)}(z, z')\| &\leq C^{1+|r|_1} \mathbf{r}! e^{-c2^h d_E(z, z')} \int_{2^{-2h-2}}^{2^{-2h}} \eta^{-(3+|r|_1)/2} d\eta \\ &\leq (C')^{1+|r|_1} \mathbf{r}! 2^{(1+|r|_1)h} e^{-c2^h d_E(z, z')}, \end{aligned} \tag{B.14}$$

for $h^* < h < 0$, and $\|\partial^r \mathbf{g}_E^{(0)}(z, z')\| \leq C^{1+|r|_1} e^{-c d_E(z, z')}$ for $h = 0$.

Recalling that $\mathbf{g}^{(h)} = \mathbf{g}_\infty^{(h)} + \mathbf{g}_E^{(h)}$ and noting that $d_E(z, z') \geq \|z - z'\|_1$, inequalities (B.3) and (B.14) also imply that $\mathbf{g}^{(h)}$ satisfies a bound of the form (2.2.15).

In order to prove (B.13), we start from (B.11). Recalling that $\mathcal{Q}_M(k_1)$ is the set of roots of $B(k_1) \sin k_2 M - \sin k_2(M + 1)$, with $k_2 \in (-\pi, \pi]$, that $N_M(k_1, k_2)$ is given by (A.14), and that $G_\#^{[\eta]}(k_1, k_2; z, z')$ are entire and 2π -periodic, for $\# \in \{\pm\}$, we can rewrite

$$\begin{aligned} &\sum_{k_2 \in \mathcal{Q}_M(k_1)} \frac{1}{2N_M(k_1, k_2)} G_\#^{[\eta]}(k_1, k_2; z, z') \\ &= \frac{1}{2} \oint_{\mathcal{C}} \frac{B(k_1) \cos k_2 M - \cos k_2(M + 1)}{B(k_1) \sin k_2 M - \sin k_2(M + 1)} G_\#^{[\eta]}(k_1, k_2; z, z') \frac{dk_2}{2\pi i}, \end{aligned} \tag{B.15}$$

where \mathcal{C} is the boundary of the rectangle in the complex plane of vertices $-\pi - ib, \pi - ib, \pi + ib, -\pi + ib$, $b > 0$, traversed counterclockwise. We rewrite

$$\frac{B(k_1) \cos k_2 M - \cos k_2(M + 1)}{B(k_1) \sin k_2 M - \sin k_2(M + 1)} = \mp i \left[1 + 2 \frac{R_\pm(k_1, k_2) e^{\pm 2ik_2(M+1)}}{1 - R_\pm(k_1, k_2) e^{\pm 2ik_2(M+1)}} \right], \tag{B.16}$$

where

$$R_{\pm}(k_1, k_2) := \frac{1 - B(k_1)e^{\mp ik_2}}{1 - B(k_1)e^{\pm ik_2}}, \tag{B.17}$$

and using this and noting that the contributions of the left and right sides of \mathcal{C} in the integral in Eq. (B.15) cancel by the periodicity of the integrand, we then have

$$\begin{aligned} & \sum_{k_2 \in \mathcal{Q}_M(k_1)} \frac{1}{2N_M(k_1, k_2)} G_{\#}^{[\eta]}(k_1, k_2; z, z') \\ &= \sum_{\sigma=0, \pm 1} \int_{-\pi+i\sigma b}^{\pi+i\sigma b} A_{\sigma}(k_1, k_2) G_{\#}^{[\eta]}(k_1, k_2; z, z') \frac{dk_2}{2\pi} \end{aligned} \tag{B.18}$$

where $A_0(k_1, k_2) \equiv 1$ and, if $\sigma = \pm 1$,

$$A_{\sigma}(k_1, k_2) := \frac{R_{\sigma}(k_1, k_2)e^{2i\sigma k_2(M+1)}}{1 - R_{\sigma}(k_1, k_2)e^{2i\sigma k_2(M+1)}}. \tag{B.19}$$

We now need to sum (B.18) over $k_1 \in \mathcal{D}_L$. Notice that the right side of (B.18) is analytic in k_1 in a sufficiently small strip around the real axis (quantitative bounds on the width of the analyticity strip will follow) and is 2π -periodic in k_1 . Given any function $F(k_1)$ that is 2π -periodic and analytic in a strip of width $2b > 0$ around the real axis, we have

$$\frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} F(k_1) = \oint_{\mathcal{C}} \frac{dk_1}{2\pi} \frac{F(k_1)}{1 + e^{-ik_1 L}} = \sum_{\sigma'=0, \pm 1} \int_{-\pi+i\sigma' b}^{\pi+i\sigma' b} A'_{\sigma'}(k_1) F(k_1) \frac{dk_1}{2\pi}, \tag{B.20}$$

where $A'_0(k_1) \equiv 1$ and, if $\sigma' = \pm 1$, $A'_{\sigma'}(k_1) = -e^{i\sigma' k_1 L} / (1 + e^{i\sigma' k_1 L})$; moreover \mathcal{C} is the same contour defined after (B.15). Using (B.11), (B.18) and (B.20), we obtain

$$\begin{aligned} \partial^r \mathfrak{g}_E^{[\eta]}(z, z') &= - \iint_{[-\pi, \pi]^2} \frac{dk_1 dk_2}{(2\pi)^2} \left[\partial^r G_{-}^{[\eta]}(k_1, k_2; z, z') \right. \\ &+ \sum_{\# = \pm} \sum_{\sigma, \sigma' = 0, \pm}^* A'_{\sigma'}(k_1 + i\sigma' b) A_{\sigma}(k_1 + i\sigma' b, k_2 + i\sigma b) \partial^r G_{\#}^{[\eta]} \\ &\left. (k_1 + i\sigma' b, k_2 + i\sigma b; z, z') \right], \end{aligned} \tag{B.21}$$

where the $*$ on the sum indicates the constraint that $(\sigma, \sigma') \neq (0, 0)$, and

$$\begin{aligned} \partial^r G_{\#}^{[\eta]}(k_1, k_2; z, z') &= (e^{-ik_1} - 1)^{r_{1,1}} (e^{ik_1} - 1)^{r_{2,1}} (e^{-ik_2} - 1)^{r_{1,2}} \\ &(e^{ik_2} - 1)^{r_{2,2}} G_{\#}^{[\eta]}(k_1, k_2; z, z'). \end{aligned} \tag{B.22}$$

Now, by using the definition of $G_{-}^{[\eta]}$ and by proceeding as in the proof of (B.1), we see that the first term in the right side of (B.21) admits the same bound as $\mathfrak{g}_{\infty}^{[\eta]}(z - z')$, see (B.1), with the only difference that $|z - z'|_1$ should be replaced

by $|z_1 - z'_1| + \min\{z_2 + z'_2, 2(M + 1) - z_2 - z'_2\} \leq d_E(z, z')$. Therefore, the first term in the right side of (B.21) satisfies (B.13) as desired.

Let us now prove that the contribution to $[\partial^r \mathfrak{g}_E^{[\eta]}(z, z')]_{\omega\omega'}$ from the second line of (B.21) satisfy (B.13). For this purpose, if $\sigma' = 0$, we shift k_1 in the complex plane as $k_1 \rightarrow k_1 - i b \operatorname{sign}(z_1 - z'_1)$. If $\sigma = 0$, we shift k_2 in the complex plane as $k_2 \rightarrow k_2 - i\tau b$, with $\tau = \pm$, its specific valued depending on \sharp and on the matrix element (ω, ω') we are looking at; more precisely, $\tau = \tau_{\sharp, (\omega\omega')}$, with

$$\tau_{\sharp, (\omega\omega')} := \begin{cases} -\operatorname{sign}(z_2 - z'_2) & \text{if } \sharp = + \\ -1 & \text{if } \sharp = - \text{ and } (\omega, \omega') \neq (-, -) \\ +1 & \text{if } \sharp = - \text{ and } (\omega, \omega') = (-, -). \end{cases} \quad (\text{B.23})$$

Once these complex shifts are performed, we bound the contribution to $[\partial^r \mathfrak{g}_E^{[\eta]}(z, z')]_{\omega\omega'}$ from the second line of (B.21) by the sum over \sharp and over σ, σ' (with $(\sigma, \sigma') \neq (0, 0)$) of:

$$\iint_{[-\pi, \pi]^2} \frac{dk_1 dk_2}{(2\pi)^2} |A'_{\sigma'}(k_1 + i\tilde{\sigma}'b)| |A_{\sigma}(k_1 + i\tilde{\sigma}'b, k_2 + i\tilde{\sigma}b)| \left| [\partial^r G_{\sharp}^{[\eta]}(k_1 + i\tilde{\sigma}'b, k_2 + i\tilde{\sigma}b; z, z')]_{\omega\omega'} \right|, \quad (\text{B.24})$$

where $\tilde{\sigma} = \begin{cases} \sigma & \text{if } \sigma \neq 0 \\ \tau_{\sharp, (\omega\omega')} & \text{if } \sigma = 0 \end{cases}$, and $\tilde{\sigma}' = \begin{cases} \sigma' & \text{if } \sigma' \neq 0 \\ -\operatorname{sign}(z_1 - z'_1) & \text{if } \sigma' = 0 \end{cases}$. Note that, if $\sigma' \neq 0$, then

$$|A'_{\sigma'}(k_1 + i\tilde{\sigma}'b)| \leq \frac{e^{-bL}}{1 - e^{-bL}}. \quad (\text{B.25})$$

If $\sigma \neq 0$, we recall that $A_{\sigma}(k_1, k_2)$ is given by (B.19), with $R_{\sigma}(k_1, k_2)$ as in (B.17). We claim that, if $b \leq c_0$, with c_0 sufficiently small, and k_1, k_2 real, then

$$|R_{\sigma}(k_1 + i\tilde{\sigma}'b, k_2 + i\tilde{\sigma}b)| \leq e^{Cb}. \quad (\text{B.26})$$

for some $C > 0$; this will be proved momentarily, after (B.29). We now pick $b = c_0 \min\{1, \eta^{-1/2}\}$, so that, using (B.26), for $\sigma \neq 0$, M sufficiently large, and k_1, k_2 real,

$$|A_{\sigma}(k_1 + i\tilde{\sigma}'b, k_2 + i\tilde{\sigma}b)| \leq \frac{e^{Cb}}{1 - e^{-bM}} e^{-2b(M+1)}. \quad (\text{B.27})$$

If we now use (B.25) and (B.27) in the second line of (B.21) and we estimate the integral of $|\partial^r G_{\sharp}^{[\eta]}(k_1 + i\tilde{\sigma}'b, k_2 + i\tilde{\sigma}b; z, z')|$ via the same strategy used in the proof of (B.1), see Eqs. (B.7) to (B.10), we find that, for $\eta \leq 1$,

$$(\text{B.24}) \leq C^{1+|r|_1} e^{-c_0 d_E(z, z')}, \quad (\text{B.28})$$

where d_E was defined after (2.2.16), while, if $1 \leq \eta \leq 2^{-2h^*}$,

$$\begin{aligned} (\text{B.24}) &\leq C^{1+|r|_1} e^{-c_0 \eta^{-1/2} d_E(z, z')} \int_{\mathbb{R}^2} dk_1 dk_2 e^{-c\eta(k_1^2 + k_2^2)} (|k_1| + |k_2| + \eta^{-1/2})^{1+|r|_1} \\ &\leq (C')^{1+|r|_1} e^{-c_0 \eta^{-1/2} d_E(z, z')} \mathbf{r}! \eta^{-\frac{3+|r|_1}{2}}. \end{aligned} \quad (\text{B.29})$$

This completes the proof that the terms in the second line of (B.21) satisfy (B.13), with $c = c_0$, provided that the bound (B.26) holds.

Proof of (B.26). We will prove the following version of (B.26): if $\sigma \neq 0$, $b \leq c_0$ with c_0 sufficiently small and $a := \tilde{\sigma}'b$ with $\tilde{\sigma}' = \pm 1$, then

$$|R_\sigma(k_1 + ia, k_2 + i\sigma b)|^2 \leq e^{2b}(1 + C_0b), \tag{B.30}$$

where C_0 can be chosen

$$C_0 = \max \{8\kappa(1 + \pi^2), 128\pi^4\kappa(1 - 2\kappa)^{-2}\}, \tag{B.31}$$

and κ is the same as in (A.10). By using the definition of R_σ , (B.17), one sees that (B.30) is equivalent to

$$e^{-2b} + \rho^2 - 2\rho e^{-b} \cos(\sigma k_2 - \beta) \leq (1 + C_0b)[1 + \rho^2 e^{-2b} - 2\rho e^{-b} \cos(\sigma k_2 + \beta)], \tag{B.32}$$

where $\rho = \rho(k_1, a) := |B(k_1 + ia)|$ and $\beta = \beta(k_1, a) := \text{Arg}(B(k_1 + ia))$. As shown below, if $-\pi \leq k_1 \leq \pi$ and $|a| = b \leq c_0$ with c_0 sufficiently small, then

$$\rho \leq 1 + \kappa[(1 + \pi^2)b^2 - k_1^2/\pi^2], \quad |\beta| \leq \frac{4\kappa}{1 - 2\kappa}|k_1| \cdot b. \tag{B.33}$$

By rearranging the terms in the two sides, one sees that (B.32) is equivalent to

$$(\rho^2 - 1)(1 - e^{-2b}) \leq 4\rho e^{-b} \sin(\sigma k_2) \sin \beta + C_0b[(1 - \rho e^{-b})^2 + 2\rho e^{-b}(1 - \cos(\sigma k_2 + \beta))]. \tag{B.34}$$

By using the first bound in (B.33), the fact that $|a| = b \leq c_0$ with c_0 sufficiently small, and the bound $1 - \cos(\sigma k_2 + \beta) \geq (\sigma k_2 + \beta)^2/\pi^2$ valid for β sufficiently small, it is straightforward to check that the left side of (B.34) is smaller or equal than $4\kappa b[(1 + \pi^2)b^2 - k_1^2/\pi^2]$, while the right side is greater or equal than

$$-4\rho e^{-b}|k_2| \cdot |\beta| + C_0b \left[\frac{b^2}{2} + \frac{2}{\pi^2} \rho e^{-b} (|k_2| - |\beta|)^2 \right].$$

Therefore, (B.34) is a consequence of

$$4\kappa(1 + \pi^2)b^3 + 4\rho e^{-b}|k_2| \cdot |\beta| \leq 4\kappa b k_1^2/\pi^2 + C_0 \frac{b^3}{2} + C_0 b \frac{2}{\pi^2} \rho e^{-b} (|k_2| - |\beta|)^2. \tag{B.35}$$

Now, the first term in the left side of (B.35) is smaller than the third term in the right side, $C_0 b^3/2$, because $C_0 \geq 8\kappa(1 + \pi^2)$, see (B.31). By using the second bound in (B.33) and the fact that $\rho e^{-b} \leq 1$ for b small enough (thanks to the first bound in (B.33)), we see that, if $|k_2| \leq \frac{1-2\kappa}{4\pi^2}|k_1|$, then the second term in the left side of (B.35) is smaller than the first term in the right side. In the complementary case, $|k_2| \geq \frac{1-2\kappa}{4\pi^2}|k_1|$ (which implies, in particular, that $|k_2| \geq 2|\beta|$, thanks to the second bound in (B.33)), then the second term in

the left side of (B.35) is bounded from above by $\frac{16\kappa}{1-2\kappa}\rho e^{-b}|k_2| \cdot |k_1|$, while the last term in the right side is bounded from below by $\frac{C_0 b}{2\pi^2}\rho e^{-b}|k_2|^2$; now,

$$\frac{16\kappa}{1-2\kappa}\rho e^{-b}|k_2| \cdot |k_1| \leq \frac{C_0 b}{2\pi^2}\rho e^{-b}|k_2|^2 \iff |k_2| \geq \frac{32\pi^2 \kappa}{C_0(1-2\kappa)}|k_1|,$$

which is verified for $|k_2| \geq \frac{1-2\kappa}{4\pi^2}|k_1|$, because $C_0 \geq 128\pi^4\kappa(1-2\kappa)^{-2}$, see (B.31). In conclusion, (B.35) is always verified and, as a consequence, (B.34) (and, therefore, (B.30)) is, as desired.

We are left with proving the validity of (B.33) for $|a|=b$ small enough. By definition

$$\rho e^{i\beta} = 1 - \kappa + \kappa \cos k_1 \cosh a - i\kappa \sin k_1 \sinh a, \tag{B.36}$$

so that, using $1 - \cos k_1 \geq 2k_1^2/\pi^2$ and the fact that, for $|a|=b$ small, $\cosh a \leq 1 + b^2$ and $|\sinh a| \leq 2b$,

$$\rho \leq 1 + \kappa(b^2 - 2k_1^2/\pi^2 + 2b|k_1|).$$

Using $2b|k_1| \leq \pi^2 b^2 + k_1^2/\pi^2$, we get the first of (B.33). Finally, from (B.36), we find

$$|\beta| \leq \frac{\kappa|\sin k_1 \sinh a|}{1 - \kappa + \kappa \cos k_1 \cosh a} \leq \frac{2\kappa b|k_1|}{1 - \kappa(2 + b^2)}.$$

Now, picking b^2 smaller than $(1 - 2\kappa)/2\kappa$, we find that β satisfies the second of (B.33). □

Gram representation: proof of items 3 and 4. Recall that

$$\begin{aligned} & \partial^{(s,s')} g_{\omega\omega'}^{(h)}(z, z') \\ &= \int_{I_h} d\eta \sum_{k_1 \in \mathcal{D}_L} \sum_{k_2 \in \mathcal{Q}_M(k)} \frac{1}{2LN_M(k_1, k_2)} \left[\partial^{(s,s')} G_{+,\omega\omega'}^{[\eta]}(k_1, k_2; z, z') \right. \\ & \quad \left. - \partial^{(s,s')} G_{-,\omega\omega'}^{[\eta]}(k_1, k_2; z, z') \right], \tag{B.37} \end{aligned}$$

where $I_0 = [0, 1)$, and $I_h = [2^{-2h-2}, 2^{-2h})$ for all $h^* \leq h < 0$. We recall that $\partial^{(s,s')} G_{\#}^{[\eta]}$ is given by (B.22), with s playing the role of $(r_{1,1}, r_{1,2})$ and s' playing the role of $(r_{2,1}, r_{2,2})$. In the following, we will exhibit a Gram decomposition separately for the two terms in (B.37) corresponding to $G_{+}^{[\eta]}$ and $G_{-}^{[\eta]}$, which will immediately imply a Gram decomposition for the combination of the two.

We rewrite

$$\mathbb{1}_{I_h}(\eta) \partial^{(s,s')} G_{\#,\omega\omega'}^{[\eta]}(k_1, k_2; z, z') = \sum_{\sigma=1}^4 [\tilde{\gamma}_{\#, \omega, s, z}^{(h)}(k_1, k_2, \eta)]_{\sigma}^* [\gamma_{\#, \omega', s', z'}^{(h)}(k_1, k_2, \eta)]_{\sigma}, \tag{B.38}$$

where $\mathbb{1}_{I_h}$ is the characteristic function of the interval I_h , and $[\gamma_{\#, \omega, s, z}^{(h)}(k_1, k_2, \eta)]_{\sigma}$, $[\tilde{\gamma}_{\#, \omega, s, z}^{(h)}(k_1, k_2, \eta)]_{\sigma}$ are the components of the following

4-vectors:

$$\tilde{\gamma}_{\sharp,+,s,z}^{(h)}(k_1, k_2, \eta) = \mathbb{1}_{I_h}(\eta) e^{ik_1 z_1 + ik_2 z_2} (e^{ik_1} - 1)^{s_1} (e^{ik_2} - 1)^{s_2} \sqrt{f_\eta(k_1, k_2)} \begin{bmatrix} \left[\sqrt{\hat{g}_{\sharp,++}(k_1, k_2)} \right]^* \\ \left[\sqrt{\hat{g}_{\sharp,+-(k_1, k_2)} \right]^* \\ 0 \\ 0 \end{bmatrix},$$

$$\tilde{\gamma}_{\sharp,-,s,z}^{(h)}(k_1, k_2, \eta) = \mathbb{1}_{I_h}(\eta) e^{ik_1 z_1 + ik_2 z_2} (e^{ik_1} - 1)^{s_1} (e^{ik_2} - 1)^{s_2} \sqrt{f_\eta(k_1, k_2)} \begin{bmatrix} 0 \\ 0 \\ \left[\sqrt{\hat{g}_{\sharp,-+}(k_1, k_2)} \right]^* \\ \left[\sqrt{\hat{g}_{\sharp,--}(k_1, k_2)} \right]^* \end{bmatrix},$$

$$\gamma_{\sharp,+,s,z}^{(h)}(k_1, k_2, \eta) = \mathbb{1}_{I_h}(\eta) e^{ik_1 z_1 + \sharp ik_2 z_2} (e^{ik_1} - 1)^{s_1} (e^{\sharp ik_2} - 1)^{s_2} \sqrt{f_\eta(k_1, k_2)} \begin{bmatrix} \sqrt{\hat{g}_{\sharp,++}(k_1, k_2)} \\ 0 \\ \sqrt{\hat{g}_{\sharp,-+}(k, q)} \\ 0 \end{bmatrix},$$

$$\gamma_{\sharp,-,s,z}^{(h)}(k_1, k_2, \eta) = \mathbb{1}_{I_h}(\eta) e^{ik_1 z_1 + \sharp ik_2 z_2} (e^{ik_1} - 1)^{s_1} (e^{\sharp ik_2} - 1)^{s_2} \sqrt{f_\eta(k_1, k_2)} \begin{bmatrix} 0 \\ \sqrt{\hat{g}_{\sharp,+-(k_1, k_2)} \right]^* \\ 0 \\ \sqrt{\hat{g}_{\sharp,--}(k_1, k_2)} \end{bmatrix},$$

where if $\sharp = +$, then $\hat{g}_{+, \omega \omega'}(k_1, k_2)$, with $\omega, \omega' \in \{\pm\}$, are the components of the 2×2 matrix $\hat{\mathbf{g}}_+(k_1, k_2) \equiv \hat{\mathbf{g}}(k_1, k_2)$, see (2.1.20); if $\sharp = -$, then $\hat{g}_{\sharp, \omega \omega'}(k_1, k_2)$, with $\omega, \omega' \in \{\pm\}$, are the components of

$$\hat{\mathbf{g}}_-(k_1, k_2) \equiv \begin{bmatrix} \hat{g}_{++}(k_1, k_2) & \hat{g}_{+-}(k_1, -k_2) \\ \hat{g}_{-+}(k_1, k_2) & e^{2ik_2(M+1)} \hat{g}_{--}(k_1, k_2) \end{bmatrix},$$

cf. (C.6). The square roots $\sqrt{\mathbf{g}_{\sharp, \omega \omega'}(k_1, k_2)}$ of the complex numbers $\mathbf{g}_{\sharp, \omega \omega'}(k_1, k_2)$ are all defined by the same (arbitrarily chosen) branch.

In conclusion, in light of (B.38), (B.37) can be rewritten as

$$\begin{aligned} & \partial^{(s,s')} \mathbf{g}_{\omega \omega'}^{(h)}(z, z') \\ &= \int_0^\infty d\eta \sum_{k_1 \in \mathcal{D}_L} \sum_{k_2 \in \mathcal{Q}_M(k_1)} \frac{1}{2LN_M(k_1, k_2)} \sum_{\sigma=1}^4 \sum_{\sharp=\pm} \sharp [\tilde{\gamma}_{\sharp, \omega, s, z}^{(h)}(k_1, k_2, \eta)]_\sigma^* \\ & \quad [\gamma_{\sharp, \omega', s', z'}^{(h)}(k_1, k_2, \eta)]_\sigma \\ &= \left(\tilde{\gamma}_{+, \omega, s, z}^{(h)} \otimes \hat{e}_1 + \tilde{\gamma}_{-, \omega, s, z}^{(h)} \otimes \hat{e}_2, \gamma_{+, \omega', s', z'}^{(h)} \otimes \hat{e}_1 - \gamma_{-, \omega', s', z'}^{(h)} \otimes \hat{e}_2 \right) \end{aligned}$$

$$\equiv \left(\tilde{\gamma}_{\omega, s, z}^{(h)}, \gamma_{\omega', s', z'}^{(h)} \right), \tag{B.39}$$

where in the last line \hat{e}_1, \hat{e}_2 are the elements of the standard Euclidean basis of \mathbb{R}^2 . We can adapt all of the preceding discussion to $\mathfrak{g}^{(\leq h)}$ simply by replacing I_h with $[2^{-2h-2}, \infty)$; this concludes the proof of item 3.

In order to prove the bounds in item 4, we first note that the definitions given above for $\tilde{\gamma}_{\omega, s, z}^{(h)}, \gamma_{\omega, s, z}^{(h)}$ immediately imply

$$\begin{aligned} \left| \tilde{\gamma}_{\omega, s, z}^{(h)} \right|^2, \left| \gamma_{\omega, s, z}^{(h)} \right|^2 &\leq \int_{I_h} d\eta \sum_{k \in \mathcal{D}_L} \sum_{k_2 \in \mathcal{Q}_M(k)} \frac{1}{LN_M(k_1, k_2)} \\ \left| e^{ik_1} - 1 \right|^{2s_1} \left| e^{ik_2} - 1 \right|^{2s_2} \cdot f_\eta(k_1, k_2) &\sum_{\omega, \omega' = \pm} |g_{\omega\omega'}(k_1, k_2)|. \end{aligned} \tag{B.40}$$

Now, recall that the set \mathcal{D}_L consists of points in $[-\pi, \pi]$ that are equi-spaced at a mutual distance $2\pi/L$, and that the set $\mathcal{Q}_M(k)$ consists of points in $[-\pi, \pi]$ that are *almost* equi-spaced at a mutual distance $\pi/(M+1)$ (more precisely, recall that there is exactly one point of $\mathcal{Q}_M(k)$ in every interval $\frac{\pi}{M+1}(n + \frac{1}{2}, n + 1)$, $n = 0, \dots, M - 1$, and exactly one point in every interval $\frac{\pi}{M+1}(-n - 1, -n - \frac{1}{2})$, $n = 0, \dots, M - 1$). Note also that $N_M(k_1, k_2) \geq M^{11}$, and that the summand in (B.40) is continuous, so we can bound (B.40) by a Riemann sum and obtain

$$\begin{aligned} \left| \tilde{\gamma}_{\omega, s, z}^{(h)} \right|^2, \left| \gamma_{\omega, s, z}^{(h)} \right|^2 &\leq C \int_{I_h} d\eta \iint_{[-\pi, \pi]^2} dk_1 dk_2 \left| e^{ik_1} \right. \\ &\quad \left. - 1 \right|^{2s_1} \left| e^{ik_2} - 1 \right|^{2s_2} f_\eta(k_1, k_2) \sum_{\omega, \omega' = \pm} |g_{\omega\omega'}(k_1, k_2)| \\ &\leq (C')^{1+s_1+s_2} \int_{I_h} d\eta \iint_{[-\pi, \pi]^2} dk_1 dk_2 \\ &\quad \min(1, |k_1|^{2s_1} |k_2|^{2s_2} e^{-c\eta(k_1^2+k_2^2)} (|k_1| + |k_2|)) \\ &\leq (C'')^{1+s_1+s_2} s_1! s_2! \int_{I_h} \min(1, \eta^{-s_1-s_2-3/2}) d\eta \\ &\leq (C''')^{1+2s_1+2s_2} s_1! s_2! 2^{h(1+2s_1+2s_2)}. \end{aligned} \tag{B.41}$$

Similarly

$$\begin{aligned} \left| \tilde{\gamma}_{\omega, s, z}^{(\leq h)} \right|^2, \left| \gamma_{\omega, s, z}^{(\leq h)} \right|^2 &\leq (C'')^{1+s_1+s_2} s_1! s_2! \int_{2^{-2h-2}}^\infty \eta^{-s_1-s_2-3/2} d\eta \\ &\leq (C''')^{1+2s_1+2s_2} s_1! s_2! 2^{h(1+2s_1+2s_2)}, \end{aligned} \tag{B.42}$$

and these bounds constitute item 4.

¹¹In fact, by Eq. (A.13) and the definition of $\mathcal{Q}_M(k_1)$, one has $N_M(k_1, k_2) - M = (1 - B(k_1) \cos k_2)/(B^2(k_1) - 2B(k_1) \cos k_2 + 1) \geq 0$.

C. Proof of Proposition 2.9

Recall that Λ is the discrete cylinder of sides $L = 2\lfloor a^{-1}\ell_1/2 \rfloor$ and $M + 1 = \lfloor a^{-1}\ell_2 \rfloor + 1$ and $\mathfrak{g}_{\text{scal}}(z, z')$ the scaling limit propagator (2.3.3) in the continuum cylinder Λ_{ℓ_1, ℓ_2} of sides ℓ_1, ℓ_2 . In order to emphasize its dependence upon the sides of the cylinder, let us denote the scaling limit propagator in Λ_{ℓ_1, ℓ_2} by $\mathfrak{g}_{\text{scal}}(\ell_1, \ell_2; z, z')$. Note that, upon rescaling by $\xi > 0$, this propagator satisfies:

$$\xi \mathfrak{g}_{\text{scal}}(\ell_1, \ell_2; \xi z, \xi z') = \mathfrak{g}_{\text{scal}}(\xi^{-1}\ell_1, \xi^{-1}\ell_2; z, z'). \tag{C.1}$$

We will prove that, for any $z, z' \in \Lambda$ such that $z \neq z'$, and any $w, w' \in \Lambda_{a^{-1}\ell_1, a^{-1}\ell_2}$ such that $w \neq w'$, $\|w - z\| \leq \sqrt{2}$ and $\|w' - z'\| \leq \sqrt{2}$,

$$\|\mathfrak{g}_{\text{c}}(z, z') - a\mathfrak{g}_{\text{scal}}(\ell_1, \ell_2; aw, aw')\| \leq C(\min\{L, M, \|z - z'\|\})^{-2}, \tag{C.2}$$

provided that $\min\{L, M, \|z - z'\|\}$ is sufficiently large. Proposition 2.9 readily follows from (C.2), simply by rescaling by a^{-1} . Note that, thanks to (C.1), $a\mathfrak{g}_{\text{scal}}(\ell_1, \ell_2; aw, aw') = \mathfrak{g}_{\text{scal}}(a^{-1}\ell_1, a^{-1}\ell_2; w, w')$; note also that $|a^{-1}\ell_1 - L| \leq 2\sqrt{2}$ and $|a^{-1}\ell_2 - M| \leq \sqrt{2}$. By using the explicit expression of the scaling limit propagator (2.3.3) and the fact that $\|w - z\| \leq \sqrt{2}$ and $\|w' - z'\| \leq \sqrt{2}$, we find that $\|\mathfrak{g}_{\text{scal}}(a^{-1}\ell_1, a^{-1}\ell_2; w, w') - \mathfrak{g}_{\text{scal}}(L, M + 1; z, z')\| \leq C(\min\{L, M, \|z - z'\|\})^{-2}$. Therefore, in order to prove (C.2), it is enough to show that, for $\min\{L, M, \|z - z'\|\}$ large,

$$\|\mathfrak{g}_{\text{c}}(z, z') - \mathfrak{g}_{\text{scal}}(L, M + 1; z, z')\| \leq C(\min\{L, M, \|z - z'\|\})^{-2}, \tag{C.3}$$

which is what we will prove in the rest of this appendix.

Recall that $\mathfrak{g}_{\text{c}}(z, z') = \int_0^\infty \mathfrak{g}^{[\eta]}(z, z') d\eta$, with

$$\mathfrak{g}^{[\eta]}(z, z') = \sum_{\#=\pm} \# \sum_{k_1 \in \mathcal{D}_L} \sum_{k_2 \in \mathcal{Q}_M(k_1)} \frac{1}{2LN_M(k_1, k_2)} G_{\#}^{[\eta]}(k_1, k_2; z, z'), \tag{C.4}$$

where $G_{\#}^{[\eta]}$ were defined in (B.12). Similarly, $\mathfrak{g}_{\text{scal}}(L, M + 1; z, z') = \int_0^\infty \mathfrak{g}_{\text{scal}}^{[\eta]}(L, M + 1; z, z') d\eta$, with

$$\begin{aligned} &\mathfrak{g}_{\text{scal}}^{[\eta]}(L, M + 1; z, z') \\ &:= \sum_{\#=\pm} \# \sum_{k_1 \in \frac{\pi}{L}(2\mathbb{Z}+1)} \sum_{k_2 \in \frac{\pi}{2(M+1)}(2\mathbb{Z}+1)} \frac{1}{2L(M+1)} G_{\text{scal};\#}^{[\eta]}(k_1, k_2; z, z') \end{aligned} \tag{C.5}$$

where $d(k_1, k_2) := (1 - t_2)^2 k_1^2 + (1 - t_1)^2 k_2^2$ and

$$\begin{aligned} G_{\text{scal};+}^{[\eta]}(k_1, k_2; z, z') &:= e^{-ik_1(z_1 - z'_1) - ik_2(z_2 - z'_2) - \eta d(k_1, k_2)} \\ &\quad \begin{bmatrix} -2it_1 k_1 & -(1 - t_1^2)ik_2 \\ -(1 - t_1^2)ik_2 & 2it_1 k_1 \end{bmatrix}, \\ G_{\text{scal};-}^{[\eta]}(k_1, k_2; z, z') &:= e^{-ik_1(z_1 - z'_1) - ik_2(z_2 + z'_2) - \eta d(k_1, k_2)} \\ &\quad \begin{bmatrix} -2it_1 k_1 & (1 - t_1^2)ik_2 \\ -(1 - t_1^2)ik_2 & e^{2ik_2(M+1)} 2it_1 k_1 \end{bmatrix}. \end{aligned} \tag{C.6}$$

We rewrite

$$\mathfrak{g}_{\text{scal}}^{[\eta]}(L, M + 1; z, z') - \mathfrak{g}^{[\eta]}(z, z') = R_1^{[\eta]}(z, z') + R_2^{[\eta]}(z, z'), \tag{C.7}$$

where

$$R_1^{[\eta]}(z, z') := \sum_{\#=\pm} \# \frac{1}{2L(M+1)} \left[\sum_{(k_1, k_2) \in \mathcal{B}_{L,M}} G_{\text{scal};\#}(k_1, k_2; z, z') - \sum_{(k_1, k_2) \in \mathcal{D}_{L,M}} G_{\#}^{[\eta]}(k_1, k_2; z, z') \right], \tag{C.8}$$

with $\mathcal{B}_{L,M} := \frac{\pi}{L}(2\mathbb{Z} + 1) \times \frac{\pi}{2(M+1)}(2\mathbb{Z} + 1)$, and $\mathcal{D}_{L,M} := \mathcal{D}_L \times \mathcal{D}_{2(M+1)}$. Moreover,

$$R_2^{[\eta]}(z, z') := \sum_{\#=\pm} \# \frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} \left[\sum_{k_2 \in \mathcal{D}_{2(M+1)}} \frac{1}{2(M+1)} - \sum_{k_2 \in \mathcal{Q}_M(k_1, k_2)} \frac{1}{2N_M(k_1, k_2)} \right] G_{\#}^{[\eta]}(k_1, k_2; z, z'). \tag{C.9}$$

The first remainder term. We consider the contribution to $\mathfrak{g}_{\text{scal}}(L, M + 1; z, z') - \mathfrak{g}_c(z, z')$ from $R_1^{[\eta]}$ first. Examining the definitions (B.12) and (C.6), we see that, if $(k_1, k_2) \in \mathcal{D}_{L,M}$, each matrix element of $G_{\text{scal};\#}^{[\eta]}(k_1, k_2; z, z') - G_{\#}^{[\eta]}(k_1, k_2; z, z')$ is bounded in absolute value by $C|k|^2 e^{-c\eta|k|^2}$, with $|k|^2 = k_1^2 + k_2^2$, for some $C, c > 0$; if $(k_1, k_2) \in \mathcal{B}_{L,M} \setminus \mathcal{D}_{L,M}$, each matrix element of $G_{\text{scal};\#}^{[\eta]}(k_1, k_2; z, z')$ is bounded in absolute value by $C|k|e^{-c\eta|k|^2}$. These bounds are sufficient for performing the integral of $\|R_1^{[\eta]}(z, z')\|$ over $\eta \geq (\min\{L, M\})^2$. In fact, for such values of η ,

$$\frac{1}{2L(M+1)} \sum_{(k_1, k_2) \in \mathcal{B}_{L,M}} e^{-c\eta|k|^2} \cdot \begin{cases} |k|^2 & \text{if } \max\{|k_1|, |k_2|\} < \pi \\ |k| & \text{otherwise} \end{cases} \leq C\eta^{-2}, \tag{C.10}$$

for some $C > 0$, so that

$$\int_{(\min\{L, M\})^2}^{\infty} \|R_1^{[\eta]}(z, z')\| d\eta \leq C \int_{(\min\{L, M\})^2}^{\infty} \eta^{-2} d\eta \leq C'(\min\{L, M\})^{-2}. \tag{C.11}$$

In order to bound the contribution from the integral of $\|R_1^{[\eta]}(z, z')\|$ over $\eta \leq (\min\{L, M\})^2$, we need to rewrite $R_1^{[\eta]}(z, z')$ as a suitable integral in the complex plane, in analogy with what we did in Appendix B. More precisely, by using (B.20) and its analogue for the sums over k_2 , we find that the matrix elements of $R_1^{[\eta]}(z, z')$ can be rewritten as:

$$[R_1^{[\eta]}(z, z')]_{\omega\omega'} = \sum_{\#=\pm} \# \sum_{\sigma_1, \sigma_2=0, \pm} \cdot \left\{ \int_{-\infty+i\sigma_1 b}^{\infty+i\sigma_1 b} \frac{dk_1}{2\pi} \int_{-\infty+i\sigma_2 b}^{\infty+i\sigma_2 b} \frac{dk_2}{2\pi} A'_{\sigma_1}(k_1) A''_{\sigma_2}(k_2) [G_{\text{scal};\#}(k_1, k_2; z, z')]_{\omega\omega'} \right.$$

$$- \left. \int_{-\pi+i\sigma_1 b}^{\pi+i\sigma_1 b} \frac{dk_1}{2\pi} \int_{-\pi+i\sigma_2 b}^{\pi+i\sigma_2 b} \frac{dk_2}{2\pi} A'_{\sigma_1}(k_1) A''_{\sigma_2}(k_2) [G_{\#}^{[\eta]}(k_1, k_2; z, z')]_{\omega\omega'} \right\}, \tag{C.12}$$

where b will be conveniently fixed below, $A'_\sigma(k)$ was defined after (B.20) and A''_σ is defined analogously: $A''_0(k) \equiv 1$ and, if $\sigma = \pm$, $A''_\sigma(k) = -e^{2i\sigma k(M+1)}/(1 + e^{2i\sigma k(M+1)})$. We now proceed as described before and after (B.23): if $\sigma_1 = 0$, we shift k_1 in the complex plane as $k_1 \rightarrow k_1 - ib \operatorname{sign}(z_1 - z'_1)$; if $\sigma_2 = 0$, depending on the values of $\#, \omega, \omega'$, we shift $k_2 \rightarrow k_2 - i\tau_{\#, (\omega\omega')} b$, with $\tau_{\#, (\omega\omega')}$ as in (B.23). Next, we combine the third line of (C.12) with the contribution to the integral in the second line from the region $\max\{|\Re k_1|, |\Re k_2|\} \leq \pi$. After these manipulations, we find that $[R_1^{[\eta]}(z, z')]_{\omega\omega'}$ can be further rewritten as

$$[R_1^{[\eta]}(z, z')]_{\omega\omega'} = \sum_{\# = \pm} \sum_{\sigma_1, \sigma_2 = 0, \pm} \int_{-\infty+i\tilde{\sigma}_1 b}^{\infty+i\tilde{\sigma}_1 b} \frac{dk_1}{2\pi} \int_{-\infty+i\tilde{\sigma}_2 b}^{\infty+i\tilde{\sigma}_2 b} \frac{dk_2}{2\pi} A'_{\sigma_1}(k_1) A''_{\sigma_2}(k_2) \cdot \begin{cases} [G_{\text{scal}; \#}(k_1, k_2; z, z')]_{\omega\omega'} - [G_{\#}^{[\eta]}(k_1, k_2; z, z')]_{\omega\omega'} & \text{if } \max\{|\Re k_1|, |\Re k_2|\} \leq \pi \\ [G_{\text{scal}; \#}(k_1, k_2; z, z')]_{\omega\omega'} & \text{if } \max\{|\Re k_1|, |\Re k_2|\} > \pi \end{cases} \tag{C.13}$$

where $\tilde{\sigma}_1 = -\operatorname{sign}(z_1 - z'_1)$, if $\sigma_1 = 0$, and $\tilde{\sigma}_1 = \sigma_1$, otherwise; and $\tilde{\sigma}_2 = \tau_{\#, (\omega\omega')}$, if $\sigma_2 = 0$, and $\tilde{\sigma}_2 = \sigma_2$, otherwise. We now pick $b = \eta^{-1/2}$ and notice that, if $\Im k_1 = \tilde{\sigma}_1 \eta^{-1/2}$ and $\Im k_2 = \tilde{\sigma}_2 \eta^{-1/2}$, with $\eta \leq (\min\{L, M\})^2$, the integrand in the right side of (C.12) is bounded in absolute value by

$$C e^{\eta^{-1/2}} e^{-\eta^{-1/2} \|z-z'\| - c\eta |k|^2} \cdot \begin{cases} |k|^2 + \eta^{-1} & \text{if } \max\{|\Re k_1|, |\Re k_2|\} \leq \pi \\ |k| + \eta^{-1/2} & \text{if } \max\{|\Re k_1|, |\Re k_2|\} > \pi \end{cases}$$

for some $C, c > 0$. Therefore, recalling that $\|z - z'\| \gg 1$,

$$\|R_1^{[\eta]}(z, z')\| \leq C e^{-\frac{1}{2}\eta^{-1/2} \|z-z'\|} \left[\int_{[-\pi, \pi]^2} dk e^{-c\eta |k|^2} (|k|^2 + \eta^{-1}) + \int_{\mathbb{R}^2 \setminus [-\pi, \pi]^2} dk e^{-c\eta |k|^2} (|k| + \eta^{-1/2}) \right] \leq C' e^{-\frac{1}{2}\eta^{-1/2} \|z-z'\|} (\eta^{-2} + e^{-c'\eta} \eta^{-3/2}). \tag{C.14}$$

Note that $e^{-\frac{1}{4}\eta^{-1/2} \|z-z'\| - c'\eta} \leq C e^{-c'' \|z-z'\|^{2/3}}$, so that, by integrating over $\eta \leq (\min\{L, M\})^2$, we find:

$$\int_0^{(\min\{L, M\})^2} \|R_1^{[\eta]}(z, z')\| d\eta \leq C (\|z - z'\|^{-2} + \|z - z'\|^{-1} e^{-c'' \|z-z'\|^{2/3}}) \leq C' \|z - z'\|^{-2}. \tag{C.15}$$

Combining this with (C.11), we find that $\int_0^\infty \|R_1^{[\eta]}(z, z')\| d\eta \leq C(\min\{L, M, \|z - z'\|\})^{-2}$.

The second remainder term. Let us now consider the contribution to $\mathfrak{g}_{\text{scal}}(L, M + 1; z, z') - \mathfrak{g}_{\text{c}}(z, z')$ from $R_2^{[\eta]}$. By using (B.18) and the analogue of (B.20) for the sums over k_2 in $\mathcal{D}_{2(M+1)}$, we rewrite $R_2^{[\eta]}(z, z')$ as

$$R_2^{[\eta]}(z, z') := \sum_{\#=\pm} \frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} \sum_{\sigma=\pm} \int_{-\pi+i\sigma b}^{\pi+i\sigma b} \frac{dk_2}{2\pi} \times [A''_\sigma(k_2) - A_\sigma(k_1, k_2)] G_\#^{[\eta]}(k_1, k_2; z, z'), \tag{C.16}$$

where A''_σ and A_σ were defined after (C.12) and in (B.19), respectively. Note that, for $\sigma = \pm$,

$$\begin{aligned} A''_\sigma(k_2) - A_\sigma(k_1, k_2) &= -\frac{e^{2i\sigma k(M+1)}}{1 + e^{2i\sigma k(M+1)}} - \frac{R_\sigma(k_1, k_2)e^{2i\sigma k_2(M+1)}}{1 - R_\sigma(k_1, k_2)e^{2i\sigma k_2(M+1)}} \\ &= -\frac{2(1 + R_\sigma(k_1, k_2))e^{2i\sigma k_2(M+1)}}{(1 + e^{2i\sigma k_2(M+1)})(1 - R_\sigma(k_1, k_2)e^{2i\sigma k_2(M+1)})}. \end{aligned} \tag{C.17}$$

Recalling the definition (B.17) of R_σ , we have

$$|1 + R_\sigma(k_1, k_2)| = 2 \left| \frac{1 - B(k_1) \cos(k_2)}{1 - B(k_1)e^{i\sigma k_2}} \right| \leq C \frac{|k_1|^2 + |k_2|^2 + b^2}{b} \tag{C.18}$$

for $|\Im k_1| \leq \sigma \Im k_2 = b$ positive and sufficiently small; and recalling the bound (B.26) on R_σ , which remains valid with $i\tilde{\sigma}b$ replaced by $i\tilde{\sigma}b'$, $|b'| \leq b$, the denominator of (C.17) is bounded from below as

$$|1 - R_\sigma(k_1, k_2)e^{2i\sigma k_2(M+1)}| |1 + e^{2i\sigma k_2(M+1)}| \geq (1 - e^{-b(M+1)})^2 \tag{C.19}$$

for M larger than some constant and $|\Im k_1| \leq \sigma \Im k_2 = b$ positive and sufficiently small. We now proceed slightly differently, depending on whether η is larger or smaller than $(\ell(z, z'))^2$, with $\ell(z, z') := \max\{\min\{L, M\}, \|z - z'\|\}$.

The case of η smaller than $(\ell(z, z'))^2$. In this case, we rewrite (C.16) by using (B.20); if $\sigma' = 0$, we perform the complex shift $k_1 \rightarrow k_1 - ib \text{sign}(z_1 - z'_1)$, thus getting (letting $\tilde{\sigma}' = \sigma'$, if $\sigma' = \pm$, and $\tilde{\sigma}' = -\text{sign}(z_1 - z'_1)$, if $\sigma' = 0$)

$$R_2^{[\eta]}(z, z') := \sum_{\#=\pm} \sum_{\sigma'=0, \pm} \sum_{\sigma=\pm} \int_{-\pi+i\tilde{\sigma}'b}^{\pi+i\tilde{\sigma}'b} \frac{dk_1}{2\pi} \int_{-\pi+i\sigma b}^{\pi+i\sigma b} \frac{dk_2}{2\pi} A'_{\sigma'}(k_1) [A''_\sigma(k_2) - A_\sigma(k_1, k_2)] G_\#^{[\eta]}(k_1, k_2; z, z'). \tag{C.20}$$

We now bound the integrand by its absolute value, by using, in particular, (C.17)–(C.18), and by estimating the matrix elements of $G_\#^{[\eta]}$ in the same way as we did several times above and in Appendix B. We thus get, for $b =$

$c_0 \min\{1, \eta^{-1/2}\}$ with c_0 sufficiently small,

$$\begin{aligned} \left\| R_2^{[\eta]}(z, z') \right\| &\leq \frac{C e^{-b\|z-z'\|}}{(1 - e^{-bL})(1 - e^{-bM})^2} \int_{[-\pi, \pi]^2} dk \frac{(|k| + b)^3}{b} e^{-\eta|k|^2} \\ &\leq \frac{C b^4 e^{-b\|z-z'\|}}{(1 - e^{-bL})(1 - e^{-bM})^2}. \end{aligned} \tag{C.21}$$

If we now integrate this inequality with respect to η , for $0 \leq \eta \leq \ell(z, z')$, recalling that $\ell(z, z') = \max\{\min\{L, M\}, \|z - z'\|\}$, we find, for $\|z - z'\|$,

$$\begin{aligned} \int_0^{\ell(z, z')^2} \left\| R_2^{[\eta]}(z, z') \right\| d\eta &\leq C \int_0^1 e^{-c_0\|z-z'\|} d\eta \\ &\quad + C \int_1^{(\min\{L, M\})^2} \eta^{-2} e^{-c_0\eta^{-1/2}\|z-z'\|} d\eta \\ &\quad + \mathbf{1}(\|z - z'\| > \min\{L, M\}) \\ &\quad \frac{C}{LM^2} \int_{(\min\{L, M\})^2}^{\|z-z'\|^2} \eta^{-1/2} e^{-c_0\eta^{-1/2}\|z-z'\|} d\eta \\ &\leq C' \left(e^{-c_0\|z-z'\|} + \frac{1}{\|z - z'\|^2} + \frac{\|z - z'\|}{LM^2} \right) \\ &\leq \frac{C''}{(\min\{L, M, \|z - z'\|\})^2}. \end{aligned} \tag{C.22}$$

The case of η larger than $\ell(z, z')$. In this case we go back to the representation (C.16) (no rewriting of the sum over k_1 in terms of an integral in the complex plane).

We proceed slightly differently for the diagonal and off-diagonal elements of $R_2^{[\eta]}$. Let us begin with the diagonal terms. Note that the diagonal elements of $G_{\sharp}^{[\eta]}$ have the form

$$\pm 2it_1 e^{-ik_1(z_1 - z'_1)} e^{-ik_2 Z_2} e^{-\eta D(k_1, k_2)} \sin k_1$$

where Z_2 is either $z_2 - z'_2$, $z_2 + z'_2$, or $z_2 + z'_2 - 2M - 2$. We thus see that each diagonal element of $R_2^{[\eta]}$ is given by a sum of four terms (due to the sums over \sharp and σ) of the form

$$\begin{aligned} &\pm \frac{2it_1}{L} \sum_{k_1 \in \mathcal{D}_L} \int_{-\pi+i\sigma b}^{\pi+i\sigma b} [A'_\sigma(k_2) - A_\sigma(k_1, k_2)] e^{-ik_1(z_1 - z'_1)} e^{-ik_2 Z_2} e^{-\eta D(k_1, k_2)} \\ &\quad \sin k_1 \frac{dk_2}{2\pi} \\ &= \pm \frac{2t_1}{L} \sum_{k_1 \in \mathcal{D}_L} \int_{-\pi+i\sigma b}^{\pi+i\sigma b} \frac{2(1 + R_\sigma(k_1, k_2)) e^{i\sigma k_2 Z'_{2,\sigma}} \sin k_1 (z_1 - z'_1) e^{-\eta D(k_1, k_2)} \sin k_1}{(1 - R_\sigma(k_1, k_2)) e^{2i\sigma k_2(M+1)} (1 + e^{2i\sigma k_2(M+1)})} \\ &\quad \frac{dk_2}{2\pi}, \end{aligned} \tag{C.23}$$

where in passing from the first to the second line we used (C.17) and the fact that $R_\sigma(k_1, k_2)$ is even in k_1 . Moreover, in the second line, $Z'_{2,\sigma}$ is either $2(M + 1) - \sigma(z_2 - z'_2)$, $2(M + 1) - \sigma(z_2 + z'_2)$, or $2(M + 1)(1 + \sigma) - \sigma(z_2 + z'_2)$;

in any case, $Z'_{2,\sigma} \geq 2$. We can then use this, together with Inequalities (C.18) and (C.19) and the observation that $|\sin k_1(z_1 - z'_1)| \leq |k_1| \cdot \|z - z'\|$ to obtain, for $\omega = \pm$ and $b = c_0\eta^{-1/2}$, with $\eta \geq \ell(z, z')$,

$$\begin{aligned} \left| [R_2^{[\eta]}(z, z')]_{\omega\omega} \right| &\leq \frac{C\|z - z'\|}{L} \sum_{k_1 \in \mathcal{D}_L} \int_{-\pi}^{\pi} dk_2 \frac{\eta^{1/2} (k_1^2 + k_2^2 + \eta^{-1})}{(1 - e^{-c_0\eta^{-1/2}M})^2} |k_1|^2 e^{-c\eta(k_1^2 + k_2^2)} \\ &\leq \frac{C\|z - z'\|\eta^{-3/2}}{M^2}, \end{aligned} \tag{C.24}$$

and thus, recalling that $\ell(z, z') = \max\{\min\{L, M\}, \|z - z'\|\}$,

$$\left| \int_{(\ell(z, z'))^2}^{\infty} [R_2^{[\eta]}(z, z')]_{\omega\omega} d\eta \right| \leq \frac{C\|z - z'\|}{M^2 \ell(x, y)} \leq \frac{C}{M^2}, \tag{C.25}$$

which is of the desired order.

The off-diagonal elements of $G_{\sharp}^{[\eta]}$ are equal, up to a sign, to

$$2(1 - t_1^2)e^{-ik_1(z_1 - z'_1)} e^{-ik_2 Z_2} e^{-\eta D(k_1, k_2)} (1 - B(k_1)e^{\pm ik_2}),$$

where as before Z_2 is either $z_2 - z'_2$, $z_2 + z'_2$, or $z_2 + z'_2 - 2M - 2$. Noting that $R_{\pm}(k_1, k_2) = R_{\mp}(k_1, -k_2)$, we rewrite

$$\begin{aligned} &\int_{-\pi - ib}^{\pi - ib} [A''_-(k_2) - A_-(k_1, k_2)] G_{\sharp}^{[\eta]}(k_1, k_2; z, z') \frac{dk_2}{2\pi} \\ &= \int_{-\pi + ib}^{\pi + ib} [A''_+(k_2) - A_+(k_1, k_2)] G_{\sharp}^{[\eta]}(k_1, -k_2; z, z') \frac{dk_2}{2\pi} \end{aligned} \tag{C.26}$$

and so each off-diagonal element of $R_2^{[\eta]}$ can be written as a sum of two terms (due to the sum over \sharp) of the form

$$\begin{aligned} \frac{1}{L} \sum_{k_1 \in \mathcal{D}_L} \int_{-\pi + ib}^{\pi + ib} &\frac{2(1 + R_+(k_1, k_2))e^{2ik_2(M+1)}}{(1 - R_+(k_1, k_2)e^{2ik_2(M+1)})(1 + e^{2ik_2(M+1)})} e^{-ik_1(z_1 - z'_1)} e^{-\eta D(k_1, k_2)} \\ &\times [e^{-ik_2 Z_2} (1 - B(k_1)e^{\pm ik_2}) + e^{ik_2 Z_2} (1 - B(k_1)e^{\mp ik_2})] \frac{dk_2}{2\pi} \end{aligned}$$

up to uninteresting coefficients; then noting that

$$\begin{aligned} &\frac{1}{2} [e^{-ik_2 Z_2} (1 - B(k_1)e^{\pm ik_2}) + e^{ik_2 Z_2} (1 - B(k_1)e^{\mp ik_2})] \\ &= \cos k_2 Z_2 - B(k_1) \cos k_2 (Z_2 \mp 1) \\ &= (1 - \cos k_2) \cos k_2 Z_2 \mp \sin k_2 \sin k_2 Z_2 + [1 - B(k_1)] \cos k_2 (Z_2 \mp 1), \end{aligned}$$

we obtain, for $\omega = \pm$ and $b = c_0\eta^{-1/2}$ with $\eta \geq \ell(z, z')$, noting also that $|Z_2| \leq 2M$,

$$\begin{aligned} \left| [R_2^{[\eta]}(z, z')]_{\omega, -\omega} \right| &\leq \frac{C|Z_2|}{L} \sum_{k_1 \in \mathcal{D}_L} \int_{-\pi}^{\pi} dk_2 \frac{\eta^{1/2} (k_1^2 + k_2^2 + \eta^{-1})^2}{(1 - e^{-c_0\eta^{-1/2}M})^2} e^{-c\eta(k_1^2 + k_2^2)} \\ &\leq \frac{C}{M} \eta^{-3/2}, \end{aligned} \tag{C.27}$$

from which

$$\left| \int_{(\ell(z,z'))^2}^{\infty} \left[R_2^{[\eta]}(z, z') \right]_{\omega, -\omega} d\eta \right| \leq \frac{C}{M\ell(z, z')}, \tag{C.28}$$

which is again of the desired order. Combining this with (C.25) and (C.22), we find that $\int_0^\infty \|R_2^{[\eta]}(z, z')\| d\eta \leq C(\min\{L, M, \|z - z'\|\})^{-2}$. Together with the bound on $R_1^{[\eta]}$, see the line after (C.15), this concludes the proof of (C.3) and, therefore, of Proposition 2.9.

D. Non-interacting Correlation Functions in the Scaling Limit

In this appendix, we explain how to express the scaling limit of the non-interacting correlation function appearing in Theorem 1.1 in terms of the propagators studied in Appendix C, thus proving, in particular, [4, Eq.(1.12)]. For notational simplicity, in this appendix we let $t_1^*(\lambda) \equiv t_1$ and $t_2^*(\lambda) \equiv t_2 = (1 - t_1)/(1 + t_1)$. By using the Grassmann representation of Proposition 3.1 in the case $\lambda = 0$, we find that, for the lattice of unit mesh and any m -tuple of distinct edges x_1, \dots, x_m , with $m \geq 2$,

$$\begin{aligned} & \langle \epsilon_{x_1}; \dots; \epsilon_{x_m} \rangle_{0, t_1, t_2; \Lambda} \\ &= \frac{\partial^m}{\partial A_{x_1} \dots \partial A_{x_m}} \log \int \mathcal{D}\Phi e^{S_{t_1, t_2}(\Phi) + \sum_{x \in \mathfrak{B}_\Lambda} (1 - t_j^*(x)) E_x A_x} \Big|_{\mathbf{A}=\mathbf{0}} \end{aligned} \tag{D.1}$$

(note that the expectation in the left side is the *truncated* one). Introducing the rescaled energy observable $\varepsilon_\ell^a(z) := a^{-1} \sigma_z \sigma_{z+a\hat{e}_\ell}$, rescaling the lattice by a factor of a and passing over to the non-truncated expectation, we obtain

$$\begin{aligned} & \langle \varepsilon_{l_1}^a(z_1) \dots \varepsilon_{l_m}^a(z_m) \rangle_{0, t_1, t_2; \Lambda^a} \\ &= a^{-m} (1 - t_1^2)^{m_1} (1 - t_2^2)^{m_2} \langle (E_{x(z_1, l_1)} - \langle E_{x(z_1, l_1)} \rangle) \dots (E_{x(z_m, l_m)} - \langle E_{x(z_m, l_m)} \rangle) \rangle, \end{aligned} \tag{D.2}$$

where, in the right side: $E_{x(z, 1)} = \bar{H}_z H_{z+a\hat{e}_1}$ and $E_{x(z, 2)} = \bar{V}_z V_{z+a\hat{e}_2}$; the symbol $\langle (\cdot) \rangle$ indicates normalized Grassmann measure $\frac{\int \mathcal{D}\Phi e^{S_{t_1, t_2}^a(\Phi)} (\cdot)}{\int \mathcal{D}\Phi e^{S_{t_1, t_2}^a(\Phi)}}$, with \mathcal{S}_{t_1, t_2}^a the same as (2.1.2) on the rescaled lattice Λ^a .

Recall the transformation (2.1.10) relating the variables $\{\bar{H}_z, H_z, \bar{V}_z, V_z\}_{z \in \Lambda}$ to $\{\phi_{\omega, z}, \xi_{\omega, z}\}_{z \in \Lambda}^{\omega \in \{\pm\}}$, from which we see that, if x is a vertical edge of endpoints $z, z + a\hat{e}_2$, then $E_x = \phi_{+, z} \phi_{-, z+a\hat{e}_2}$, while, if x is a horizontal edge of endpoints $z, z + a\hat{e}_1$, then (with obvious notation) $E_x = [s_+ * (\phi_+ - \phi_-)(z)] [s_- * (\phi_+ + \phi_-)(z + a\hat{e}_1)]$ plus terms involving the ‘massive’ variables $\{\xi_{\omega, z}\}_{z \in \Lambda, \omega \in \{\pm\}}$.

The reader can convince herself that, for the purpose of computing the limit $a \rightarrow 0^+$ of (D.2), in the right side of (D.2) we can freely replace E_x by the following local expressions in the Grassmann ‘massless’ variables: $\phi_{+, z} \phi_{-, z}$, if x is a vertical edge of endpoints $z, z + a\hat{e}_2$ (note that $\phi_{+, z} \phi_{-, z}$ is obtained from $\phi_{+, z} \phi_{-, z+a\hat{e}_2}$ by ‘localizing’ the second field at the same position of the first

one); and $(1+t_1)^{-2}(\phi_{+,z} - \phi_{-,z})(\phi_{+,z} + \phi_{-,z}) = 2(1+t_1)^{-2}\phi_{+,z}\phi_{-,z}$, if x is a vertical edge of endpoints $z, z+a\hat{e}_1$ (note that $(1+t_1)^{-2}(\phi_{+,z} - \phi_{-,z})(\phi_{+,z} + \phi_{-,z})$ is obtained from $[s_+ * (\phi_+ - \phi_-)(z)] [s_- * (\phi_+ + \phi_-)(z+a\hat{e}_1)]$ by localizing $[s_- * (\phi_+ + \phi_-)(z+a\hat{e}_1)]$ at z , and by replacing the non-local, exponentially decaying, kernels $s_{\pm}(z_1)$ by their local counterparts, namely $c_0\delta_{z_1,0}$, with $c_0 = \lim_{L \rightarrow \infty} \sum_{y=1}^L s_{\pm}(y) = (1+t_1)^{-1}$). It is, in fact, easy to check that the difference between the exact expression of E_x and such a ‘local approximations’ is of higher order in a and its contribution to the correlation function vanishes in the limit $a \rightarrow 0$. Therefore,

$$\begin{aligned} & \lim_{a \rightarrow 0^+} \langle \varepsilon_{l_1}^a(z_1) \cdots \varepsilon_{l_m}^a(z_m) \rangle_{0,t_1,t_2;\Lambda^a} \\ &= \lim_{a \rightarrow 0^+} a^{-m} \left(\frac{2(1-t_1^2)}{(1+t_1)^2} \right)^{m_1} (1-t_2^2)^{m_2} \langle : \phi_{+,z_1} \phi_{-,z_1} : \cdots : \phi_{+,z_m} \phi_{-,z_m} : \rangle, \end{aligned} \tag{D.3}$$

where $: \phi_{+,z} \phi_{-,z} :$ denotes the difference $\phi_{+,z}\phi_{-,z} - \langle \phi_{+,z}\phi_{-,z} \rangle$. Note that $\frac{2(1-t_1^2)}{(1+t_1)^2} = 2t_2$. The Grassmann average in the right side of (D.3) can be expressed in terms of the fermionic Wick rule or, equivalently, in terms of the Pfaffian of the $2m \times 2m$ anti-symmetric matrix $\mathcal{M}^a(\mathbf{z})$, whose elements, labeled by the indices $(1,+), (1,-), \dots, (m,+), (m,-)$, are equal to

$$[\mathcal{M}^a(\mathbf{z})]_{(i,\omega)(j,\omega')} = \begin{cases} \langle \phi_{\omega,z_i} \phi_{\omega',z_j} \rangle & \text{if } i \neq j, \\ 0 & \text{otherwise.} \end{cases}$$

In view of Proposition 2.9, $\lim_{a \rightarrow 0} a^{-1} \langle \phi_{\omega,z} \phi_{\omega',z'} \rangle = [\mathfrak{g}_{\text{scal}}(z, z')]_{\omega\omega'}$ and, therefore,

$$\lim_{a \rightarrow 0^+} \langle \varepsilon_{l_1}^a(z_1) \cdots \varepsilon_{l_m}^a(z_m) \rangle_{0,t_1,t_2;\Lambda^a} = (2t_2)^{m_1} (1-t_2^2)^{m_2} \text{Pf}(\mathcal{M}(\mathbf{z})), \tag{D.4}$$

with

$$[\mathcal{M}(\mathbf{z})]_{(i,\omega)(j,\omega')} = \begin{cases} [\mathfrak{g}_{\text{scal}}(z_i, z_j)]_{\omega\omega'} & \text{if } i \neq j, \\ 0 & \text{otherwise,} \end{cases} \tag{D.5}$$

Since $\mathfrak{g}_{\text{scal}}$ is covariant under rescaling, see (C.1), the scaling limit (D.4) is, as well. Note that rescalings are, together with translations and parity, the only conformal transformations from finite cylinders to finite cylinders or, equivalently, from a finite circular annulus to a finite circular annulus: in fact, it is well known [5,28] that an annulus $\{z \in \mathbb{C} : r < |z| < R\}$ can be conformally mapped to another one only if the two annuli have the same modulus $\frac{1}{2\pi} \log(R/r)$; moreover, any automorphism of the annulus $\{z \in \mathbb{C} : r < |z| < R\}$ is either a rotation $z \rightarrow ze^{i\theta}$ or a rotation followed by an inversion $z \rightarrow Rr/z$. Equivalently, in terms of finite cylinders, this classical result of complex analysis implies that the only conformal transformations from finite cylinders to finite cylinders are uniform rescaling, translations and parity.

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Giovanni Antinucci
Section de mathématiques
Université de Genève
2-4 rue du Lièvre, 1211 Genève 4
Geneva
Switzerland

Alessandro Giuliani and Rafael L. Greenblatt
Dipartimento di Matematica e Fisica
Università degli Studi Roma Tre
L.go S. L. Murialdo 1
00146 Roma
Italy

Alessandro Giuliani
Centro Linceo Interdisciplinare Beniamino Segre
Accademia Nazionale dei Lincei, Palazzo Corsini
Via della Lungara 10
00165 Roma
Italy
e-mail: giuliani@mat.uniroma3.it

Present Address

Rafael L. Greenblatt
Scuola Internazionale Superiore di Studi Avanzati (SISSA)
Mathematics Area, Via Bonomea 265
34136 Trieste
Italy

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